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Volume 12

Issue 12

Version 1.0

ENG



GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH: F
MATHEMATICS & DECISION SCIENCES



GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH: F
MATHEMATICS & DECISION SCIENCES

VOLUME 12 ISSUE 12 (VER. 1.0)

OPEN ASSOCIATION OF RESEARCH SOCIETY

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Offset Typesetting

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

An Efficient Class of Ratio-Cum-Dual to Product Estimator of Finite Population Mean in Sample Surveys

By Sanjib Choudhury & Bhupendra Kumar Singh

North Eastern Regional Institute of Science and Technology

Abstract - We consider a class of ratio-cum-dual to product estimator for estimating a finite population mean of the study variate. The bias and mean square error of the proposed estimator have been obtained. The asymptotically optimum estimator (AOE) in this class has also been identified along with its approximate bias and mean square error. Theoretical and empirical studies have been done to demonstrate the superiority of the proposed estimator over the other estimators.

Keywords : *Finite population mean; ratio-cum-dual to product estimator; Bias; Mean square error; Efficiency.*

GJSFR-F Classification : *MSC 2010: 62D05*



AN EFFICIENT CLASS OF RATIO-CUM-DUAL TO PRODUCT ESTIMATOR OF FINITE POPULATION MEAN IN SAMPLE SURVEYS

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An Efficient Class of Ratio-Cum-Dual to Product Estimator of Finite Population Mean in Sample Surveys

Sanjib Choudhury^α & Bhupendra Kumar Singh^σ

Abstract - We consider a class of ratio-cum-dual to product estimator for estimating a finite population mean of the study variate. The bias and mean square error of the proposed estimator have been obtained. The asymptotically optimum estimator (AOE) in this class has also been identified along with its approximate bias and mean square error. Theoretical and empirical studies have been done to demonstrate the superiority of the proposed estimator over the other estimators.

Keywords : Finite population mean; ratio-cum-dual to product estimator; Bias; Mean square error; Efficiency.

I. INTRODUCTION

In sample surveys, auxiliary information is used at both selections as well as estimation stages to improve the efficiency of the estimators. The use of auxiliary information at the estimation stage appears to have started with the work of Cochran (1940). When the correlation between study variate and auxiliary variate is positive (high), the ratio method of estimation is used for estimating the population mean. The ratio method is most effective if $\rho C_y / C_x > 1/2$, where C_y , C_x and ρ are coefficient of variation of y , coefficient of variation of x and correlation coefficient between y and x respectively. On the other hand, if the correlation is negative, the product method of estimation is used and this is most effective if $\rho C_y / C_x < -1/2$, suggested by Murthy (1964). Srivenkataramana (1980) first proposed dual to ratio estimator and Bandyopadhyay (1980) proposed dual to product estimator. Singh and Tailor (2005), Singh and Espejo (2003), Tailor and Sharma (2009) worked on ratio-cum-product estimators. Sharma and Tailor (2010), Choudhury and Singh (2012) worked on ratio, dual to ratio and dual to product estimators to estimate the study variable. These motivated authors to propose a new ratio-cum-dual to product estimators for estimating the population mean.

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Consider a finite population $U = (u_1, u_2, \dots, u_N)$ of size N units. Let y and x denotes the study and auxiliary variates respectively. A sample of size n ($n < N$) is drawn using simple random sampling without replacement (SRSWOR) to estimate the population mean $\bar{Y} = (1/N) \sum_{i=1}^N y_i$ of the study variate y . Let the sample mean (\bar{x}, \bar{y}) are the unbiased estimator of (\bar{X}, \bar{Y}) based on n observations.

The usual ratio and product estimators for \bar{Y} are

$$\bar{y}_R = \bar{y} \left(\frac{\bar{X}}{\bar{x}} \right)$$

and

$$\bar{y}_P = \bar{y} \left(\bar{x} / \bar{X} \right) \quad \text{respectively,}$$

where $\bar{y} = (1/n) \sum_{i=1}^n y_i$ and $\bar{x} = (1/n) \sum_{i=1}^n x_i$.

Let $x_i^* = (1+g)\bar{X} - gx_i$, $i = 1, 2, \dots, N$, where $g = n/(N-n)$.

Then clearly $\bar{x}^* = (1+g)\bar{X} - g\bar{x}$ is also unbiased estimator for \bar{X} and $Corr(\bar{y}, \bar{x}^*) = -\rho$.

Using the transformation $x_i^* = (1+g)\bar{X} - gx_i$, Srivenkataramana (1980) obtained dual to ratio estimator as

$$\bar{y}_R^* = \bar{y} \left(\bar{x}^* / \bar{X} \right)$$

and Bandyopadhyay (1980) obtained dual to product estimator as

$$\bar{y}_P^* = \bar{y} \left(\bar{X} / \bar{x}^* \right).$$

In this paper, we have proposed a class of ratio-cum-dual to product type estimator for estimating population mean \bar{Y} . Numerical illustrations are given in the support of the present study.

II. THE PROPOSED ESTIMATOR

For estimating population mean \bar{Y} , we propose an estimator as

$$\bar{y}_{RdP} = \bar{y} \left[\alpha \left(\frac{\bar{X}}{\bar{x}} \right) + (1-\alpha) \left(\frac{\bar{X}}{\bar{x}^*} \right) \right] \quad (1)$$

where α is a suitably chosen scalar.

To obtain the bias and mean square error (MSE) of \bar{y}_{RdP} to a first degree of approximation, we write

$$e_0 = (\bar{y} - \bar{Y})/\bar{Y} \text{ and } e_1 = (\bar{x} - \bar{X})/\bar{X}$$

Such that

$$\left. \begin{aligned} E(e_0) = E(e_1) = 0, \quad E(e_0^2) &= \frac{1-f}{n} C_y^2, \\ E(e_1^2) &= \frac{1-f}{n} C_x^2, \quad E(e_0 e_1) = \frac{1-f}{n} C C_x^2, \end{aligned} \right\} \quad (2)$$

where $f = n/N$ is the sampling fraction, $C_y^2 = S_y^2/\bar{Y}^2$, $C_x^2 = S_x^2/\bar{X}^2$, $C = \rho C_y/C_x$ and defined as $\rho = S_{xy}/S_x S_y$, $S_x^2 = \frac{1}{N-1} \sum_{i=1}^N (x_i - \bar{X})^2$, $S_y^2 = \frac{1}{N-1} \sum_{i=1}^N (y_i - \bar{Y})^2$ and $S_{xy} = \frac{1}{N-1} \sum_{i=1}^N (y_i - \bar{Y})(x_i - \bar{X})$.

Expressing \bar{y}_{RdP} in terms of e 's, we obtain

$$\bar{y}_{RdP} = \bar{Y} (1 + e_0) \left\{ \alpha (1 + e_1)^{-1} + (1 - \alpha) (1 - g e_1)^{-1} \right\}.$$

We now assume that $|e_1| < 1$ and $|g e_1| < 1$, so that we may expand $(1 + e_1)^{-1}$ and $(1 - g e_1)^{-1}$ as a series in powers of e_1 and $g e_1$ respectively. Expanding, multiplying out and retaining terms of e 's to the second degree, we obtain

$$\bar{y}_{RdP} - \bar{Y} \cong \bar{Y} \left[e_0 + g (e_1 + e_0 e_1 + g e_1^2) + \alpha (1 + g) \left\{ -e_1^2 + (1 - g) e_1 - e_0 e_1 \right\} \right] \quad (3)$$

Taking the expectation of both sides in equation (3) and using the results of equation (2) we get the bias of \bar{y}_{RdP} as

$$B(\bar{y}_{RdP}) = \frac{1-f}{n} \bar{Y} C_x^2 \left[\left\{ g^2 - \alpha (g^2 - 1) \right\} + C \left\{ g - \alpha (g + 1) \right\} \right] \quad (4)$$

The bias, $B(\bar{y}_{RdP})$ in (4) is 'zero' if $\alpha = \frac{g(C+g)}{(1+g)(1-g-C)}$. Thus the estimator \bar{y}_{RdP}

with $\alpha = \frac{g(C+g)}{(1+g)(1-g-C)}$ is almost unbiased.

Squaring and taking expectations of both the sides of equation (3) and using the results of equation (2), we obtain the MSE of \bar{y}_{RdP} to the first degree of approximation as

$$M(\bar{y}_{RdP}) = \frac{1-f}{n} \bar{Y}^2 \left[C_y^2 + C_x^2 \left\{ g - \alpha (1 + g) \right\} \left\{ 2C + g - \alpha (1 + g) \right\} \right] \quad (5)$$

which is minimized when

$$\alpha = \frac{1}{1+g}(g+C) = \alpha_{opt.} \text{ (say)} \quad (6)$$

Substituting equation (6) in equation (1) yield the 'asymptotically optimum estimator' (AOE) as

$$\bar{y}_{RdP}^{opt.} = \bar{y} \left[\left(\frac{g+C}{1+g} \right) \frac{\bar{X}}{\bar{x}} + \left(\frac{1-C}{1+g} \right) \frac{\bar{X}}{\bar{x}^*} \right]$$

Thus the resulting bias and MSE of $\bar{y}_{RdP}^{opt.}$ respectively as

$$B(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y} C_x^2 (1-C)(g+C) \quad (7)$$

and

$$M(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y}^2 C_y^2 (1-\rho^2) \quad (8)$$

which is the same as the MSE of the linear regression estimator $\bar{y}_{reg.} = \bar{y} + b_{yx}(\bar{X} - \bar{x})$, where b_{yx} is the sample regression coefficient of y on x .

From equation (7), we note that the bias of AOE $\bar{y}_{RdP}^{opt.}$ is 'zero' if either $C=1$ or $C=-g$.

Remark 2.1.

To the first degree of approximation, the proposed strategy \bar{y}_{RdP} under optimality condition (6), is equal to linear regression estimator.

Remark 2.2.

For $\alpha=1$, the estimator \bar{y}_{RdP} in equation (1) boils down to the usual ratio estimator \bar{y}_R . The bias and MSE of \bar{y}_R can be obtained by putting $\alpha=1$ in equations (4) and (5) respectively as

$$B(\bar{y}_R) = \frac{1-f}{n} \bar{Y} C_x^2 (1-C)$$

$$M(\bar{y}_R) = \frac{1-f}{n} \bar{Y}^2 \{C_y^2 + C_x^2 (1-2C)\} \quad (9)$$

Remark 2.3.

For $\alpha=0$, the estimator \bar{y}_{RdP} in equation (1) boils down to the dual to product estimator \bar{y}_p^* , proposed by Bandyopadhyay (1980). The bias and MSE of \bar{y}_p^* can be obtained by putting $\alpha=0$ in equations (4) and (5) respectively as

$$B(\bar{y}_p^*) = \frac{1-f}{n} \bar{Y} C_x^2 g (g+C)$$

and

$$M(\bar{y}_p^*) = \frac{1-f}{n} \bar{Y}^2 \{C_y^2 + g C_x^2 (g+2C)\} \quad (10)$$

Thus, we see that this study provides unified treatment towards the properties of different estimators.

III. EFFICIENCY COMPARISONS

a) Comparison of \bar{y}_{RdP}

In this section, firstly, we compare MSE of conventional estimators \bar{y} , \bar{y}_R and \bar{y}_P with MSE of proposed estimator \bar{y}_{RdP} .

The MSE of sample mean \bar{y} under SRSWOR sampling scheme is given by

$$M(\bar{y}) = \frac{1-f}{n} \bar{Y}^2 C_y^2. \quad (11)$$

From equations (5) and (11), it is found that the proposed estimator \bar{y}_{RdP} is more efficient than \bar{y} if

$$\{-g + \alpha(1+g)\} \{2C + g - \alpha(1+g)\} > 0$$

This condition holds if

$$\text{either } \frac{g}{1+g} > \alpha \text{ and } \frac{1}{1+g}(2C+g) < \alpha,$$

$$\text{or } \frac{g}{1+g} < \alpha \text{ and } \frac{1}{1+g}(2C+g) > \alpha.$$

Therefore, the range of α for which the proposed estimator \bar{y}_{RdP} is more efficient than \bar{y} is

$$\left[\min \left\{ \frac{g}{1+g}, \frac{1}{1+g}(2C+g) \right\}, \max \left\{ \frac{g}{1+g}, \frac{1}{1+g}(2C+g) \right\} \right].$$

From equations (5) and (9), we note that the estimator \bar{y}_{RdP} has smaller MSE than that of the usual ratio estimator \bar{y}_R if

$$\{1+g - \alpha(1+g)\} \{1-2C - g + \alpha(1+g)\} > 0$$

This condition holds if

$$\text{either } 1 > \alpha \text{ and } \frac{1}{1+g}(2C+g-1) < \alpha,$$

$$\text{or } 1 < \alpha \text{ and } \frac{1}{1+g}(2C+g-1) > \alpha.$$

Therefore, the range of α for which the proposed estimator \bar{y}_{RdP} is better than \bar{y}_R is

$$\left[\min \left\{ 1, \frac{1}{1+g}(2C+g-1) \right\}, \max \left\{ 1, \frac{1}{1+g}(2C+g-1) \right\} \right].$$

To compare the usual product estimator \bar{y}_p , we write the bias and MSE of \bar{y}_p to the first degree of approximation respectively as

$$B(\bar{y}_p) = \frac{1-f}{n} \bar{Y} C C_x^2$$

$$M(\bar{y}_p) = \frac{1-f}{n} \bar{Y}^2 \{C_y^2 + C_x^2(1+2C)\} \quad (12)$$

We note from equations (5) and (12) that the estimator \bar{y}_{RdP} will dominate over usual product estimator \bar{y}_p if

$$\{-(g-1) + \alpha(g+1)\} \{(2C+1+g) - \alpha(g+1)\} > 0$$

This condition holds if

$$\text{either } \frac{g-1}{1+g} > \alpha \text{ and } 1 + \frac{2C}{1+g} < \alpha$$

$$\text{or } \frac{g-1}{1+g} < \alpha \text{ and } 1 + \frac{2C}{1+g} > \alpha.$$

Hence, the range of α in which the proposed estimator \bar{y}_{pdp} is better than \bar{y}_p is

$$\left\{ \min \left(\frac{g-1}{1+g}, 1 + \frac{2C}{1+g} \right), \max \left(\frac{g-1}{1+g}, 1 + \frac{2C}{1+g} \right) \right\}.$$

Secondly, comparing the MSE between the proposed estimator and dual to ratio estimator \bar{y}_R^* , proposed by Srivenkataramana (1980).

The bias and MSE of \bar{y}_R^* to the first degree of approximation respectively as

$$B(\bar{y}_R^*) = -\bar{Y} \frac{1-f}{n} g C C_x^2$$

and

$$M(\bar{y}_R^*) = \frac{1-f}{n} \bar{Y}^2 \{C_y^2 + gC_x^2(g-2C)\}. \quad (13)$$

From equations (5) and (13), it is found that the proposed estimator \bar{y}_{RdP} will dominate over Srivenkataramana (1980) estimator \bar{y}_R^* if

$$\{2g - \alpha(g+1)\}\{-2C + \alpha(g+1)\} > 0$$

This condition exist if

$$\text{either } \frac{2g}{1+g} > \alpha \text{ and } \frac{2C}{1+g} < \alpha,$$

$$\text{or } \frac{2g}{1+g} < \alpha \text{ and } \frac{2C}{1+g} > \alpha.$$

Therefore, the range of α in which the proposed estimator \bar{y}_{RdP} is more efficient than dual to ratio estimator \bar{y}_R^* is

$$\left\{ \min\left(\frac{2g}{1+g}, \frac{2C}{1+g}\right), \max\left(\frac{2g}{1+g}, \frac{2C}{1+g}\right) \right\}.$$

Lastly, we compare MSE of the proposed estimator \bar{y}_{RdP} with dual to product estimator \bar{y}_p^* .

We note from equations (5) and (10) that

$$M(\bar{y}_p^*) > M(\bar{y}_{RdP}) \text{ if}$$

$$\alpha(1+g)\{2C + 2g - \alpha(1+g)\} > 0$$

This condition exist if

$$\text{either } 0 < \alpha < \frac{2}{1+g}(C+g),$$

$$\text{or } \frac{2}{1+g}(C+g) < \alpha < 0.$$

Therefore, the range of α in which the proposed estimator \bar{y}_{RdP} is more efficient than dual to product estimator \bar{y}_p^* is

$$\left[\min\left\{\frac{2(C+g)}{1+g}, 0\right\}, \max\left\{\frac{2(C+g)}{1+g}, 0\right\} \right].$$

Thus, it seems from the above results that the proposed estimator \bar{y}_{RdP} may be made better than other estimators by making a suitable choice of the value of α within the respective ranges.

b) *Comparison of 'AOE' of $\bar{y}_{RdP}^{opt.}$*

From equations (8)-(13), it is found that the 'AOE' $\bar{y}_{RdP}^{opt.}$ is more efficient than the other existing estimators like \bar{y} , \bar{y}_R , \bar{y}_P , \bar{y}_R^* and \bar{y}_P^* . Since

$$M(\bar{y}) - M(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y}^2 \rho^2 C_y^2 > 0.$$

$$M(\bar{y}_R) - M(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y}^2 C_x^2 (1-C)^2 > 0.$$

$$M(\bar{y}_P) - M(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y}^2 C_x^2 (1+C)^2 > 0.$$

$$M(\bar{y}_R^*) - M(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y}^2 C_x^2 (C-g)^2 > 0.$$

$$M(\bar{y}_P^*) - M(\bar{y}_{RdP}^{opt.}) = \frac{1-f}{n} \bar{Y}^2 C_x^2 (C+g)^2 > 0.$$

Hence, we conclude that the proposed estimator ' \bar{y}_{RdP} ' is the best (in the sense of having optimum MSE).

IV. NUMERICAL ILLUSTRATIONS

To examine the merits of the constructed estimator over its competitors numerically, we consider eight sets of data. The source of the population, the nature of the variates y and x and the values of the various parameters are listed in Table 1.

To reflect the gain in the efficiency of the proposed estimator \bar{y}_{RdP} over the estimators \bar{y} , \bar{y}_R , \bar{y}_P , \bar{y}_R^* and \bar{y}_P^* , the effective ranges along with the optimum value of α are presented in Table 2 with respect to the population data sets.

The percent relative efficiencies (PREs) of the different estimators with respect to usual unbiased estimator \bar{y} computed by the formula

$$PRE(E, \bar{y}) = \frac{M(\bar{y})}{M(E)} \times 100$$

where

$$E = \bar{y}, \bar{y}_R, \bar{y}_P, \bar{y}_R^*, \bar{y}_P^* \text{ and } \bar{y}_{RdP} \text{ or } \bar{y}_{RdP}^{opt.}$$

and are presented in Table 3.

Table 1 : Description of the populations

Popu- lation	Source	Study variate y	Auxiliary variate x	N	n	ρ	C_y	C_x	\bar{Y}
1	Steel and Torrie (1960)	Log of leaf burn in secs	Chlorine percentage	30	6	-0.4996	0.7001	0.7493	0.6860
2	Pandey and Dubey (1988)	---	---	20	8	-0.9199	0.3552	0.3943	19.55
3	Kadilar and Cingi (2006) pp. 1054	Level of apple production	Number of apple trees	106	20	0.82	4.18	2.02	15.37
4	Sukhatme and Sukhatme (1970)	Number of villages in the circles.	A circle consisting more than five villages	89	12	0.766	0.604	2.1901	3.36
5	Maddala (1977)	Consump- tion per capita.	Deflated prices of veal	30	6	-0.6823	0.2278	0.0986	7.6375
6	Murthy (1967)	Output	Fixed capital	80	20	0.9413	0.3542	0.7507	51.8264
7	Murthy (1967)	Output	Number of workers	80	20	0.9150	0.3542	0.9484	51.8264
8	Kadilar and Cingi (2006) pp. 78	---	---	106	20	0.86	5.22	2.1	2212.59

Table 2 : Effective ranges and optimum value of α of \bar{y}_{RdP} .

Population	Ranges of α in which the proposed estimator \bar{y}_{RdP} is better than					Optimum value
	\bar{y}	\bar{y}_R	\bar{y}_P	\bar{y}_R^*	\bar{y}_P^*	$\alpha_{opt.}$
1	(-0.55, 0.20)	(-1.35, 1.00)	(-0.60, 0.25)	(-0.75, 0.40)	(-0.35, 0.00)	-0.1734
2	(-0.59, 0.40)	(-1.19, 1.00)	(-0.20, 0.01)	(-0.99, 0.80)	(-0.19, 0.00)	-0.0972
3	(0.19, 2.94)	(1.00, 2.13)	(-0.62, 3.75)	(0.38, 2.75)	(0.00, 3.13)	1.5654
4	(0.13, 0.50)	(-0.36, 1.00)	(-0.73, 1.37)	(0.27, 0.37)	(0.00, 0.64)	0.3176
5	(-2.32, 0.20)	(-3.12, 1.00)	(-1.52, -0.60)	(-2.52, 0.40)	(-2.12, 0.00)	-1.0611
6	(0.25, 0.92)	(0.17, 1.00)	(-0.50, 1.67)	(0.50, 0.67)	(0.00, 1.17)	0.5831
7	(0.25, 0.76)	(0.01, 1.00)	(-0.50, 1.51)	(0.50, 0.51)	(0.00, 1.01)	0.5063
8	(0.19, 3.66)	(1.00, 2.85)	(-0.62, 4.47)	(0.38, 3.47)	(0.00, 3.85)	1.9231

Table 3 : Percentage relative efficiency of \bar{y} , \bar{y}_R , \bar{y}_P , \bar{y}_R^* , \bar{y}_P^* and \bar{y}_{RdP} or $\bar{y}_{RdP}^{opt.}$ with respect to \bar{y} .

Population	\bar{y}	\bar{y}_R	\bar{y}_P	\bar{y}_R^*	\bar{y}_P^*	\bar{y}_{RdP} or $\bar{y}_{RdP}^{opt.}$
1	100.00	†	†	†	124.34	133.26
2	100.00	†	526.50	†	537.23	650.26
3	100.00	226.76	†	120.73	†	305.25
4	100.00	†	†	220.46	†	241.99
5	100.00	†	167.59	†	115.73	187.10
6	100.00	†	†	591.38	†	877.54
7	100.00	†	†	612.44	†	614.34
8	100.00	212.82	†	117.95	†	384.02

†Relative efficiency less than 100%.

V. CONCLUSION

Table 2 provides the wide ranges along with the optimum value of α for which the proposed estimators \bar{y}_{RdP} or $\bar{y}_{RdP}^{opt.}$ are more efficient than all other estimators considered in this paper. It is also observed from Table 2 that there is a scope for choosing α to obtain better estimators than \bar{y} , \bar{y}_R , \bar{y}_P , \bar{y}_R^* and \bar{y}_P^* .

Table 3 shows that there is a substantial gain in efficiency by using proposed estimator \bar{y}_{RdP} (or $\bar{y}_{RdP}^{opt.}$) over \bar{y} , \bar{y}_R , \bar{y}_P , \bar{y}_R^* and \bar{y}_P^* . This shows that even if the scalar α deviates from its optimum value ($\alpha_{opt.}$), the suggested estimator \bar{y}_{RdP} will yield better estimates than \bar{y} , \bar{y}_R , \bar{y}_P , \bar{y}_R^* and \bar{y}_P^* . Thus it is preferred to use the proposed estimators \bar{y}_{RdP} or $\bar{y}_{RdP}^{opt.}$ in practice.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

Proof of 'J is a Radical Class' Using Amitsur Theorem

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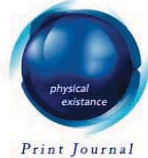
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GJSFR-F Classification : MSC 2010: 16D60, 16N80



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Abstract - The aim of this paper is to study radical class of rings, right quasi-regular rings and finally, to prove that J , the class of all right quasi-regular rings is a radical class. Amitsur gives a theorem of radical class for the sufficient condition that a class of rings would be a radical class. This paper represents, the proof of, J is a radical class using the theorem of radical class given by Amitsur.

Keywords : Ring, Ideal, Radical Class, Right quasi-regular ring.

I. INTRODUCTION

The concept of a radical was introduced by J. H. M. Wedderburn [10] in 1908, for the determination of structures of algebras and later on various radicals have been proposed by Artin [14], Baer [11], Jacobson [9], Brown-McCoy [12], Levitzki [7] etc. for the study of rings in the forties. The general theory of radicals was initiated by Kurosh [6] (1953) and Amitsur [1] in the early fifties. Andrunakievic [4], Sulinski [15], Divinsky [8] and many others have followed up the works of Kurosh and Amitsur.

Radical properties based on the notion of nilpotence do not seem to yield fruitful results for rings without chain conditions. The notion of quasi-regularity was introduced by Perlis [16]. In 1945, Jacobson [9] used it and the significant "chainless" results were obtained.

In this paper, the general ring theory covering elementary definition of rings and its ideals, homomorphism, theorem related to homomorphism and some definitions related to radical class has been discussed in preliminaries. Also, we will introduce radical class of rings and some theorems related to radical class. Amitsur gives a theorem of radical class, which works as a sufficient condition of a class of rings that would be a radical class. We will know about this theorem and also right quasi-regular ring, right quasi-regular right ideal and some lemmas related to right quasi regular rings. Finally, we will prove that J , the class of all right quasi-regular rings is a radical class. It has already been proved by using the definition of radical class. But, here we will prove this using Amitsur theorem of radical class.

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II. PRELIMINARIES

2.1. Definition

A **ring** is an ordered triple $(R, +, \cdot)$ such that R is a nonempty set and $+$ and \cdot are two binary operations on R satisfying the following axioms:

- a) R is an additive abelian group. i.e.
 - i) $a + b \in R$ for all $a, b \in R$ [closure law]
 - ii) $(a + b) + c = a + (b + c)$ for all $a, b, c \in R$. [associative law]
 - iii) there exists an element $0 \in R$ such that $a + 0 = 0 + a = a$, for all $a \in R$. [identity law]
 - iv) for every non-zero element $a \in R$ there exists an element $-a \in R$ such that $a + (-a) = (-a) + a = 0$. [inverse law]
 - v) $a + b = b + a$ for all $a, b \in R$. [commutative law]
- b) (R, \cdot) is a semi group. i.e.
 - i) $a \cdot b \in R$ for all $a, b \in R$. [closure law]
 - ii) $(a \cdot b) \cdot c = a \cdot (b \cdot c)$ for all $a, b, c \in R$ [associative law]
- c) Distributive laws are true in R . i.e. for all $a, b, c \in R$,
 - i) $a \cdot (b + c) = a \cdot b + a \cdot c$
 - ii) $(a + b) \cdot c = a \cdot c + b \cdot c$

Example

- i) $(\mathbb{Z}, +, \cdot), (\mathbb{Q}, +, \cdot), (\mathbb{R}, +, \cdot), (\mathbb{C}, +, \cdot)$ are rings.
- ii) The residue class of modulo 6, $\mathbb{Z}_6 = \{ \bar{0}, \bar{1}, \bar{2}, \bar{3}, \bar{4}, \bar{5} \}$ is a ring.
- iii) $[x]$, the set of all polynomials in x with real coefficients, is a ring.

2.2. Definition

A non-empty subset I of a ring R is called a **left (right) ideal** of R if

- i) I is an additive subgroup of R
- ii) $\forall r \in R$ and $\forall i \in I, ri \in I, (ir \in I)$.

2.3. Definition

A non-empty subset I of a ring R is called an **ideal** of R if I is both a left ideal and a right ideal of R . For a commutative ring all left and right ideals are ideals.

Example:

1. $2\mathbb{Z}$ is an ideal of \mathbb{Z} .
2. The set of integers \mathbb{Z} is only a subring but not an ideal of the ring of rational numbers \mathbb{Q} . As $3 \in \mathbb{Z}, \frac{2}{5} \in \mathbb{Q}$ but $3 \cdot \frac{2}{5} = \frac{6}{5} \notin \mathbb{Z}$.

2.4. Definition

Let R be a ring and I be an ideal of R then the **quotient ring** or **factor ring** $\frac{R}{I}$ is the set $\{ r + I : r \in R \}$, where addition and multiplication of two elements $r_1 + I, r_2 + I \in \frac{R}{I}$ are given by

- i) $(r_1 + I) + (r_2 + I) = (r_1 + r_2) + I$.
- ii) $(r_1 + I)(r_2 + I) = r_1 r_2 + I$.

Example:

$\frac{\mathbb{Z}}{2\mathbb{Z}}$ is a quotient ring.

2.5. Definition

Let $(R, +, \cdot)$ and $(R', +, \cdot)$ be two rings. A mapping $f: R \rightarrow R'$ is called a ring **homomorphism** if $\forall a, b \in R$

- i) $f(a + b) = f(a) + f(b)$ and
- ii) $f(ab) = f(a)f(b)$.

2.5.1. Theorem

Every factor ring of a ring is the homomorphic image of that ring.

Proof: Let R be a ring and I be an ideal of R , then we have to show that $\frac{R}{I}$ is a homomorphic image of R . Let us define a map $f: R \rightarrow \frac{R}{I}$ by $f(r) = r + I$ for all $r \in R$. We

need to show that f is a onto homomorphism.

Clearly f is well defined.

$$\text{Now } f(r_1 + r_2) = (r_1 + r_2) + I = (r_1 + I) + (r_2 + I) = f(r_1) + f(r_2)$$

$$\text{and } f(r_1 r_2) = r_1 r_2 + I = (r_1 + I)(r_2 + I) = f(r_1)f(r_2).$$

Thus f is a homomorphism.

Let $r + I \in \frac{R}{I}$ where $r \in R$. Then by definition, $f(r) = r + I$ i.e. $r + I = f(r)$. This implies that every element of $\frac{R}{I}$ is the image of some element in R . Thus f is onto. Hence the theorem is proved.

2.6. Definition

A ring R is said to have the **ascending chain condition (A.C.C.)** on left (right) ideals, if every ascending sequence of left (right) ideals $L_1 \subseteq L_2 \subseteq L_3 \subseteq \dots \subseteq L_n \subseteq \dots$, terminates after a finite number of steps, i.e. there exists a positive integer n such that $L_n = L_{n+1} = \dots$.

2.7. Definition

A ring R is said to have the **descending chain condition (D.C.C.)** on left (right) ideals, if every descending sequence of left (right) ideals $R \supseteq L_1 \supseteq L_2 \supseteq L_3 \supseteq \dots \supseteq L_n \supseteq \dots$, terminates after a finite number of steps, i.e. there exists a positive integer n such that $L_n = L_{n+1} = \dots$.

III. RADICAL CLASS OF RINGS

3.1. Definition

Let \mathfrak{R} be a nonempty class of rings with a certain property. A ring A is said to be an **\mathfrak{R} -ring** if $A \in \mathfrak{R}$.

Example:

Let \mathfrak{R} be the class of all nilpotent ring and A be an idempotent ring. Then A is not nilpotent ring and hence $A \notin \mathfrak{R}$. Therefore A is not an \mathfrak{R} -ring.

3.2. Definition

An ideal I of a ring A is said to be an \mathfrak{R} -ideal if I is an \mathfrak{R} -ring. i.e. $I \in \mathfrak{R}$.

Example:

Let \mathfrak{R} be the class of all nilpotent ring and I be an ideal of a nilpotent ring A . Then $I \in \mathfrak{R}$. Therefore I is an \mathfrak{R} -ideal.

3.3. Definition

A ring A is said to be \mathfrak{R} -semi-simple if A has no non-zero \mathfrak{R} -ideal.

3.4. Definition

Let \mathfrak{R} be a non-empty class of rings with a certain property. Then \mathfrak{R} is said to be a **radical property** or **radical class** if the following conditions are hold:

- A) \mathfrak{R} is homomorphically closed. i.e. every homomorphic image of an \mathfrak{R} -ring A is an \mathfrak{R} -ring. i.e. if $A \in \mathfrak{R}$ and $I \triangleleft A$, then $\frac{A}{I} \in \mathfrak{R}$.
- B) Every ring $A \in \mathfrak{R}$ contains a non-zero \mathfrak{R} -ideal $R(A)$ which contains every other \mathfrak{R} -ideals of A .
- C) $\frac{A}{R(A)}$ has no non-zero \mathfrak{R} -ideal. i.e. $\frac{A}{R(A)}$ is \mathfrak{R} -is semi-simple.

A radical class is simply called a radical.

3.5. Definition

Let \mathfrak{R} be a radical class. The \mathfrak{R} -ideal $R(A)$ of a ring A is called the \mathfrak{R} -**radical** of the ring A .

3.6. Definition

Let \mathfrak{R} -be a radical class. Then a ring A is said to be an \mathfrak{R} -**radical ring** if $R(A) = A$, where $R(A)$ is the radical of A .

3.7. Definition

Let \mathfrak{R} be a radical class. Then a ring A is said to be an \mathfrak{R} -semi-simple ring if $R(A) = 0$, where $R(A)$ is the radical of A .

0 is the only ring which is both an \mathfrak{R} -radical ring and an \mathfrak{R} -semi-simple ring.

3.7.1. Theorem [8]

Let \mathfrak{R} be a non-empty class \mathfrak{R} of rings. Then \mathfrak{R} is said to be a radical class if and only if

- A) \mathfrak{R} is homomorphically closed.
- B) If every non-zero homomorphic image of a ring A contains a non-zero \mathfrak{R} -ideal, then A is in \mathfrak{R} . i.e. $\forall I \triangleleft A$, if $\frac{A}{I} \supset \frac{B}{I} \in \mathfrak{R}$ then $A \in \mathfrak{R}$, where $B \triangleleft A$.

This theorem is known as Kurosh's Theorem.

3.7.2.1. Lemma (Zorn's Lemma)

Let A be a nonempty partially ordered set in which every totally ordered subset has an upper bound in A . Then A contains at least one maximal element.

3.7.2. Theorem (Amitsur) [3]

Let \mathfrak{R} be a nonempty class of rings. Then \mathfrak{R} is a radical class if and only if

- A') \mathfrak{R} is homomorphically closed.
- B') For any ring A and an ideal I of A if both I and $\frac{A}{I}$ is in \mathfrak{R} , then $A \in \mathfrak{R}$. i.e. \mathfrak{R} is closed under extension.
- C') If $I_1 \subseteq I_2 \subseteq I_3 \subseteq \dots$, is an ascending chain of \mathfrak{R} -ideals of a ring A , then $\bigcup_{\alpha} I_{\alpha}$ is an \mathfrak{R} -ideal.

IV. RIGHT QUASI-REGULAR RINGS

4.1. Definition

Let R be a ring and $x \in R$. Then x is called **right quasi-regular** if there exists an

Ref.

3. S. Tummurbat & R. Wisbauer, "Radicals With The α -Amitsur Property", W. Sc. P. C., Vol. 7, No. 3, 347-361, 2008.

element $y \in R$ such that $x + y + xy = 0$.

We often write $x + y + xy$ by xoy . When $xoy = 0$ then the element y is called **right quasi-inverse** of x .

4.2. Definition

A ring R is said to be **right quasi-regular** if every element in R is right quasi-regular.

4.2.1. Lemma

If R is a ring with 1, then $(1+x)$ has right inverse $(1+y)$ iff x is right quasi-regular.

Proof: Let us consider a ring with unity element 1. Let $(1+y)$ be the right inverse of $(1+x)$. Then we have,

$$(1+x)(1+y) = 1$$

$$\Rightarrow 1 + y + x + xy = 1$$

$$\Rightarrow x + y + xy = 0$$

$$\Rightarrow x \text{ is right quasi-regular.}$$

Conversely, let x be right quasi-regular. Then there is a right quasi-inverse y such that $x + y + xy = 0$

$$\Rightarrow 1 + y + x + xy = 1$$

$$\Rightarrow 1(1+y) + x(1+y) = 1$$

$$\Rightarrow (1+x)(1+y) = 1$$

i.e. $(1+y)$ is right inverse of $(1+x)$.

4.2.2. Lemma

Let R be a ring. Then for any element x in R , x is a right quasi-regular if and only if $\{r + xr\} = R, \forall r \in R$.

Proof: Let R be a ring and $x \in R$. Consider $\{r + xr\}$, the set of all elements $r + xr, \forall r \in R$. Then $\{r + xr\}$ is clearly a right ideal of R . Now suppose that $\{r + xr\} = R$. We are to show that x is right quasi-regular. Since $\{r + xr\} = R$, then $x = r + xr$ for some r in R . This implies that $x + (-r) + x(-r) = 0$. This implies that x is right quasi-regular for some $r \in R$.

Conversely, suppose that x is right quasi-regular element of R . We have to show that $R = \{r + xr\}$. Since x is right quasi-regular then \exists an element $y \in R$ such that $x + y + xy = 0 \Rightarrow x = (-y) + x(-y) \in \{r + xr\}$. Then $xr \in \{r + xr\}$ and therefore $r \in \{r + xr\}$ for every $r \in R$. Hence $\{r + xr\} = R$.

4.3. Definition

Let R be a ring and I be a right ideal of R . Then I is called a right quasi-regular right ideal if every element of I is right quasi-regular.

4.3.1. Lemma[8]

If x is right quasi-regular and if y belongs to a right quasi-regular right ideal I , then $x + y$ is right quasi-regular.

Proof: Since x is right quasi-regular, then there exists an element x' such that, $x + x' + xx' = 0$. Now, consider the element $y + yx'$. Then $y + yx'$ is in I and thus is right quasi-regular. Let z be the right quasi-inverse of $y + yx'$ then, $(y + yx') + z + (y + yx')z = 0$.

Now we will show that $x' + z + x'z$ is a right quasi-inverse of $x + y$.

$$\begin{aligned} & \text{Therefore, } (x + y) + (x' + z + x'z) + (x + y)(x' + z + x'z) \\ &= x + y + x' + z + x'z + xx' + xz + xx'z + yx' + yz + yx'z \\ &= (x + x' + xx') + (y + yx') + z + (y + yx')z + (x + x' + xx')z \\ &= 0 \end{aligned}$$

Hence $x + y$ is right quasi-regular.

4.3.2. Lemma

The sum of two right quasi-regular right ideals of a ring is also a right quasi-regular right ideal.

Proof: Let I_1 and I_2 be two right quasi-regular right ideals of a ring R . We have to show that $I_1 + I_2$ is also a right quasi-regular right ideal of R . Let $p \in I_1 + I_2$ then $p = x + y$ for some $x \in I_1$ and $y \in I_2$. Since x is right quasi-regular and $y \in I_2$ then we have $x + y$ is also right quasi-regular (by Lemma 4.3.1). i.e. p is right quasi-regular.

Hence every element of $I_1 + I_2$ is right quasi-regular.

Hence $I_1 + I_2$ is right quasi-regular right ideal of R .

4.3.3. Lemma

The sum of any finite number of right quasi-regular right ideals of a ring is again a right quasi-regular right ideal.

Proof: Let I_1, I_2, \dots, I_n are right quasi-regular right ideals of a ring R . We have to show that $I_1 + I_2 + \dots + I_n$ is right quasi-regular right ideal. We shall prove this by the method of induction on n .

If $n = 1$ then the proof is obvious. Now suppose $n = 2$, then, $I_1 + I_2$ is right quasi-regular right ideal (by Lemma 4.3.1).

Now, let $I = I_1 + I_2 + \dots + I_{n-1}$ a right quasi-regular right ideal of R . We show that $I + I_n$ is right quasi-regular right ideal of R .

Let $p \in I + I_n$ then $p = x' + y'$ for some $x' \in I$ and $y' \in I_n$. Then x' is right quasi-regular and y' belongs to a right quasi-regular right ideal I_n . Therefore $x' + y'$ is right quasi-regular (by Lemma 4.3.1). Hence $I + I_n$ is right quasi-regular right ideal. i.e. $I_1 + I_2 + \dots + I_{n-1} + I_n$ is a right quasi-regular right ideal of R .

4.3.4. Lemma

Sum (Union) of all right quasi-regular right ideals of a ring R is a right quasi-regular right ideal of R .

4.3.5. Lemma[8]

$\mathcal{J}(R)$, the sum of all right quasi-regular right ideals of a ring R is a two sided ideal of R .

Proof: Let x be any element in $\mathcal{J}(R)$ and r any element of R . We have to show that rx is in $\mathcal{J}(R)$ i.e. $\mathcal{J}(R)$ is a left ideal. We know that $\mathcal{J}(R)$ is a right quasi-regular right ideal. Hence $xr \in \mathcal{J}(R)$ is a right quasi-regular. Then there exists an element w such that

$$\begin{aligned} & xr + w + xrw = 0. \text{ Then} \\ & rx + (-rx - rwx) + rx(-rx - rwx) \\ &= rx - rx - rwx - rx \cdot rx - rx \cdot rwx \\ &= -r(w + xr + xrw)x \\ &= -r \cdot 0 \cdot x \\ &= 0 \end{aligned}$$

Ref.

8. N. Divinsky, "Rings and Radicals (University of Toronto press)", 1965.

Therefore, rx is right quasi-regular.

Next consider the right ideal generated by rx . This is the set of all $rx_i + rxs$, where i is an integer and s is in R . The element $xi + xs$ is in $J(R)$ and, as above, $r(xi + xs)$ is right quasi-regular. Therefore, $\{rx_i + rxs\}$ is a right quasi-regular right ideal. It is thus in $J(R)$ and then, in particular, rx is in $J(R)$. Therefore $J(R)$ is a two-sided ideal of R .

4.3.6. Lemma

Every homomorphic image of a right quasi-regular ring R is right quasi-regular.

Proof: Let R be a right quasi-regular ring and I be any ideal of R , then we have to show that $\frac{R}{I}$ is right quasi-regular. Let $x \in \frac{R}{I}$ then $x = r + I$ for some $r \in R$.

Since R is right quasi-regular, then r is right quasi-regular. Then there exists an element $r' \in R$ such that $r + r' + rr' = 0$.

$$\begin{aligned} \text{Now } (r + I) + (r' + I) + (r + I)(r' + I) &= r + r' + I + rr' + I \\ &= r + r' + rr' + I \\ &= 0 + I \\ &= I \end{aligned}$$

But I is the zero element of $\frac{R}{I}$. Therefore $r' + I$ is right quasi-inverse of $r + I$. Hence $r + I$ is right quasi-regular i.e. x is right quasi-regular. Therefore $\frac{R}{I}$ is right quasi-regular.

4.3.7. Lemma

Let R be a ring and I be an ideal of R . If I and $\frac{R}{I}$ are right quasi-regular then R is right quasi-regular.

Proof: Since $\frac{R}{I}$ is right quasi-regular, then for every $x \in \frac{R}{I}$, there exists $y \in \frac{R}{I}$ such that

$$\begin{aligned} (x + I) + (y + I) + (x + I)(y + I) &= I \\ \Rightarrow x + I + y + I + xy + I &= I \\ \Rightarrow x + y + xy + I &= I \\ \Rightarrow x + y + xy \in I \end{aligned}$$

Since I is right quasi-regular then there exists $w \in I$ such that

$$\begin{aligned} x + y + xy + w + (x + y + xy)w &= 0 \\ \Rightarrow x + y + xy + w + xw + yw + xyw &= 0 \\ \Rightarrow x + (y + w + yw) + x(y + w + yw) &= 0 \end{aligned}$$

This implies that $y + w + yw$ is a right quasi-inverse of x and thus x is a right quasi-regular. Hence R is right quasi-regular.

V. CONCLUSIONS

From the above discussions, we can prove the following theorem.

5.1. Theorem

The class of all right quasi-regular rings is a radical class.

Proof: Let J be the class of all right quasi-regular rings. We shall prove this using Amitsur theorem.

By Lemma 4.3.6, A' holds. i.e. J is homomorphically closed.

By Lemma 4.3.7, B') holds. i.e. J is closed under extension.

To prove C'), let $I_1 \subseteq I_2 \subseteq I_3 \subseteq \dots$ be an ascending chain of right quasi-regular right ideals of a ring R . We have to show that $\bigcup_{\alpha} I_{\alpha}$ is right quasi-regular. Let $x \in \bigcup_{\alpha} I_{\alpha}$ then $x \in I_{\alpha}$ for some α . Since each I_{α} is right quasi-regular right ideal then \exists an element x' such that $x + x' + xx' = 0$. i.e. x is right quasi-regular. Hence every element of $\bigcup_{\alpha} I_{\alpha}$ is right quasi-regular. i.e. $\bigcup_{\alpha} I_{\alpha}$ is right quasi-regular. Hence J is a radical class.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

Positive Solutions for Systems of Three-Point Nonlinear Boundary Value Problems on Time Scales

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Abstract - Values of λ are determined for which there exist positive solutions of the system of dynamic equations, $u^{\Delta\Delta}(t) + \lambda p(t)f(v(\sigma(t))) = 0$, $v^{\Delta\Delta}(t) + \lambda q(t)g(u(\sigma(t))) = 0$, for $t \in [a, b]_{\mathbb{T}}$ Satisfying the three - point boundary conditions, $\alpha u(a) - \beta u^{\Delta}(a) = 0$, $u(\sigma^2(b)) - \delta u(\eta) = 0$, $\alpha v(a) - \beta v^{\Delta}(a) = 0$, $v(\sigma^2(b)) - \delta v(\eta) = 0$, where \mathbb{T} is a time scales. A Guo-Krasnosel'skii fixed point theorem is applied.

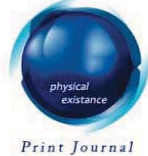
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GJSFR-F Classification : MSC 2010: 34B15, 39B10, 34B18



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Positive Solutions for Systems of Three-Point Nonlinear Boundary Value Problems on Time Scales

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Abstract - Values of λ are determined for which there exist positive solutions of the system of dynamic equations, $u^{\Delta\Delta}(t) + \lambda p(t)f(v(\sigma(t))) = 0$, $v^{\Delta\Delta}(t) + \lambda q(t)g(u(\sigma(t))) = 0$, for $t \in [a, b]_{\mathbb{T}}$. Satisfying the three-point boundary conditions, $\alpha u(a) - \beta u^{\Delta}(a) = 0$, $u(\sigma^2(b)) - \delta u(\eta) = 0$, $\alpha v(a) - \beta v^{\Delta}(a) = 0$, $v(\sigma^2(b)) - \delta v(\eta) = 0$, where \mathbb{T} is a time scales. A Guo-Krasnosel'skii fixed point theorem is applied.

Keywords : Time scales, three-point boundary value problems, dynamic equations, system of equations, positive solution, eigenvalue problem.

I. INTRODUCTION

Let \mathbb{T} be a time scale with $a, \sigma^2(b) \in \mathbb{T}$. Given an interval J of \mathbb{R} , we will use the interval notation

$$J_{\mathbb{T}} = J \cap \mathbb{T}.$$

We are concerned with determining intervals of the parameter λ (eigenvalues) for which there exist positive solutions for the system of dynamic equations,

$$\begin{aligned} u^{\Delta\Delta}(t) + \lambda p(t)f(v(\sigma(t))) &= 0, \quad t \in [a, b]_{\mathbb{T}}, \\ v^{\Delta\Delta}(t) + \lambda q(t)g(u(\sigma(t))) &= 0, \quad t \in [a, b]_{\mathbb{T}}, \end{aligned} \tag{1.1}$$

satisfying the boundary conditions

$$\begin{aligned} \alpha u(a) - \beta u^{\Delta}(a) &= 0, & u(\sigma^2(b)) - \delta u(\eta) &= 0, \\ \alpha v(a) - \beta v^{\Delta}(a) &= 0, & v(\sigma^2(b)) - \delta v(\eta) &= 0, \end{aligned} \tag{1.2}$$

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where $\alpha, \beta \geq 0, \alpha + \beta > 0, \lambda > 0, 0 < \delta < 1, \eta \in [a, \sigma^2(b)]$, and

(A1) $f, g \in C([0, \infty), [0, \infty))$,

(A2) $p, q \in C([a, \sigma(b)]_{\mathbb{T}}, [0, \infty))$, and each does not vanish identically on any closed subinterval of $[a, \sigma(b)]_{\mathbb{T}}$,

(A3) All of

$$f_0 := \lim_{x \rightarrow 0^+} \frac{f(x)}{x}, \quad g_0 := \lim_{x \rightarrow 0^+} \frac{g(x)}{x},$$

$$f_\infty := \lim_{x \rightarrow \infty} \frac{f(x)}{x}, \quad g_\infty := \lim_{x \rightarrow \infty} \frac{g(x)}{x}$$

exist as positive real numbers.

On a larger scale, there has been a great deal of activity in studying positive solutions of boundary value problems for ordinary differential equations. Interest in such solutions is high from a theoretical sense [9, 10, 12, 15] and as applications for which only positive solutions are meaningful. These considerations are cast primarily for scalar problems, but good attention has been given to boundary value problems for systems of differential equations [13, 14, 19, 21, 22].

The main tool in this paper is an application of the Guo-Krasnoselskii fixed point theorem for operators leaving a Banach space cone invariant [10]. A Green's function plays a fundamental role in defining an appropriate operator on a suitable cone.

II. SOME PRELIMINARIES

In this section, we state some preliminary lemmas and the well-known Guo-Krasnosel'skii fixed point theorem.

Let $G(t, s)$ be the Green's function for the boundary value problem

$$-y^{\Delta\Delta}(t) = 0, \tag{2.1}$$

$$\alpha y(a) - \beta y^\Delta(a) = 0, \quad y(\sigma^2(b)) - \delta y(\eta) = 0, \tag{2.2}$$

which is given by

$$G(t, s) = \frac{1}{d} \begin{cases} G_1(t, s) : a \leq s \leq \eta \\ G_2(t, s) : \eta \leq \sigma(s) \leq \sigma^2(b) \end{cases}$$

where

$$G_1(t, s) = \begin{cases} [\beta + \alpha(\sigma(s) - a)] [\sigma^2(b) - \delta\eta - t(1 - \delta)], & \sigma(s) \leq t \\ [\beta + \alpha(t - a)] [\sigma^2(b) - \delta\eta - \sigma(s)(1 - \delta)], & t \leq s \end{cases}$$

R_{ef.}

[9] L. H. Erbe and H. Wang, On the existence of positive solutions of ordinary differential equations, *Proc. Amer. Math. Soc.*, **120**(1994), No. 3, 743-748.



$$G_2(t, s) = \begin{cases} [\beta + \alpha(\sigma(s) - a)](\sigma^2(b) - t) + (t - \sigma(s))(\eta + \beta - \alpha a)\delta, & \sigma(s) \leq t \\ [\beta + \alpha(t - a)](\sigma^2(b) - \sigma(s)), & t \leq s \end{cases}$$

and

$$d := \beta(1 - \delta) + \alpha(\sigma^2(b) - a - \delta(\eta - a)).$$

Lemma 2.1 For $h(t) \in C[a, \sigma^2(b)]_{\mathbb{T}}$, the BVP

$$-y^{\Delta\Delta}(t) = h(t), \quad t \in [a, b]_{\mathbb{T}}, \quad (2.3)$$

$$\alpha y(a) - \beta y^{\Delta}(a) = 0, \quad y(\sigma^2(b)) - \delta y(\eta) = 0, \quad (2.4)$$

has a unique solution

$$\begin{aligned} y(t) = & \frac{\beta + \alpha(t - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))h(s)\Delta s \\ & - \frac{\delta(\beta + \alpha(t - a))}{d} \int_a^{\eta} (\eta - \sigma(s))h(s)\Delta s - \int_a^t (t - \sigma(s))h(s)\Delta s. \end{aligned} \quad (2.5)$$

From (2.5) obviously we have that

$$y(t) \leq \frac{\beta + \alpha(t - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))h(s)\Delta s, \quad (2.6)$$

and

$$y(\eta) \geq \frac{\beta + \alpha(\eta - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))h(s)\Delta s. \quad (2.7)$$

Lemma 2.2 Let $0 < \delta < 1$. If $h(t) \in C[a, \sigma^2(b)]_{\mathbb{T}}$, and $h \geq 0$, then the unique solution y of the problem (2.3), (2.4) satisfies

$$y(t) \geq 0, \quad t \in (a, \sigma^2(b))_{\mathbb{T}}.$$

Proof: From the fact that $y^{\Delta\Delta}(t) = -h(t) \leq 0$, we know that the graph of $y(t)$ is concave down on $[a, \sigma^2(b)]_{\mathbb{T}}$ and $y^{\Delta}(t)$ is monotone decreasing. Thus $y^{\Delta}(t) \leq y^{\Delta}(a) = \frac{\alpha}{\beta}y(a)$, where $\beta \neq 0$.

Case 1. If $y(a) < 0$, then $y^{\Delta}(t) < 0$ for $[a, \sigma^2(b)]_{\mathbb{T}}$. Thus y is a monotone decreasing function, that is $y(t) \geq y(\sigma^2(b))$.

1. If $y(\sigma^2(b)) \geq 0$, then $y(t) > 0$. So this contradicts the assertion $y(t)$ is a monotone decreasing function.

2. If $y(\sigma^2(b)) < 0$, then we have that

$$y(\eta) = \frac{1}{\delta}y(\sigma^2(b)) < 0,$$

$$y(\sigma^2(b)) = \delta y(\eta) \geq y(\eta),$$

which contradicts the assertion that $y(t)$ is monotone decreasing.

Case 2. If $y(a) \geq 0$, then $y^\Delta(a) \geq 0$. So $y(t)$ is a monotone increasing on $[a, a + \epsilon]$.

1. If $y(\sigma^2(b)) \geq 0$, then $y(t) \geq 0$ on $[a, \sigma^2(b)]_{\mathbb{T}}$.
2. If $y(\sigma^2(b)) < 0$, then we have that

$$y(\eta) = \frac{1}{\delta}y(\sigma^2(b)) < 0,$$

$$y(\sigma^2(b)) = \delta y(\eta) \geq y(\eta),$$

which contradicts the assertion that the graph of $y(t)$ is concave down on $(a, \sigma^2(b))_{\mathbb{T}}$.

If $\beta = 0$, from the boundary conditions we obtain $y(a) = 0$.

1. If $y(\sigma^2(b)) \geq 0$, then the concavity of y implies that $y(t) \geq 0$ for $t \in [a, \sigma^2(b)]_{\mathbb{T}}$.
2. If $y(\sigma^2(b)) < 0$, then

$$y(\eta) = \frac{1}{\delta}y(\sigma^2(b)) < 0,$$

$$y(\sigma^2(b)) = \delta y(\eta) \geq y(\eta).$$

This contradicts with the concavity of y .

Lemma 2.3 *If $y^{\Delta\Delta}(t) \leq 0$, then $\frac{y(\sigma^2(b))}{\sigma^2(b)} \leq \frac{y(t)}{t} \leq \frac{y(\eta)}{\eta}$ for all $t \in [\eta, \sigma^2(b)]_{\mathbb{T}}$.*

Proof: Let $h(t) := y(t) - \frac{t}{\sigma^2(b)-a}y(\sigma^2(b))$. Thus, we have $h(\eta) > 0$ and $h(\sigma^2(b)) = 0$. Since $h^{\Delta\Delta}(t) \leq 0$ then $h(t) \geq 0$ on $[\eta, \sigma^2(b)]_{\mathbb{T}}$. So $\frac{y(\sigma^2(b))}{\sigma^2(b)} \leq \frac{y(t)}{t}$. For the function $h(t)$, since $h(\eta) > 0$, $h(\sigma^2(b)) = 0$ and $h^{\Delta\Delta}(t) \leq 0$ then the function $h(t)$ is decreasing on $[\eta, \sigma^2(b)]_{\mathbb{T}}$. So $\frac{y(t)}{t} \leq \frac{y(\eta)}{\eta}$ for all $t \in [\eta, \sigma^2(b)]_{\mathbb{T}}$.

Lemma 2.4 *Let $0 < \delta < 1$. If $h(t) \in C[a, \sigma^2(b)]_{\mathbb{T}}$, and $h \geq 0$, then the unique solution y of the problem (2.3), (2.4) satisfies*

$$\inf_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} y(t) \geq \gamma \|y\|,$$

where

$$\gamma := \min \left\{ \frac{\delta(\sigma^2(b) - \eta)}{\sigma^2(b) - \delta\eta - a(1 - \delta)}, \frac{\delta\eta}{\sigma^2(b)} \right\}.$$

Proof: By the second boundary condition we know that $y(\eta) \geq y(\sigma^2(b))$. Pick $t_0 \in (a, \sigma^2(b))_{\mathbb{T}}$ such that $y(t_0) = \|y\|$. If $t_0 < \eta < \sigma^2(b)$, then

$$\min_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} y(t) = y(\sigma^2(b)),$$

and

$$\frac{y(\sigma^2(b)) - y(\eta)}{\sigma^2(b) - \eta} \leq \frac{y(\eta) - y(t_0)}{\eta - t_0}.$$

Therefore

$$\min_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} y(t) \geq \frac{\delta(\sigma^2(b) - \eta)}{\sigma^2(b) - \delta\eta - a(1 - \delta)} \|y\|.$$

If $\eta \leq t_0 < \sigma^2(b)$, again we have $y(\sigma^2(b)) = \min_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} y(t)$. From Lemma 2.3, we have $\frac{y(\eta)}{\eta} \geq \frac{y(t_0)}{t_0}$. Combining with the boundary condition $\delta y(\eta) = y(\sigma^2(b))$, we conclude that

$$\frac{y(\sigma^2(b))}{\delta\eta} \geq \frac{y(t_0)}{t_0} \geq \frac{y(t_0)}{\sigma^2(b)} = \frac{\|y\|}{\sigma^2(b)}.$$

This is

$$\min_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} y(t) \geq \frac{\delta\eta}{\sigma^2(b)} \|y\|.$$

We note that a pair $(u(t), v(t))$ is a solution of the eigenvalue problem (1.1), (1.2) if, and only if,

$$u(t) = \lambda \int_a^{\sigma(b)} G(t, s)p(s) f \left(\lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r \right) \Delta s, \quad a \leq t \leq \sigma^2(b),$$

where

$$v(t) = \lambda \int_a^{\sigma(b)} G(t, s)q(s)g(u(\sigma(s)))\Delta s, \quad a \leq t \leq \sigma^2(b).$$

Values of λ for which there are positive solutions (positive with respect to a cone) of (1.1), (1.2) will be determined via applications of the following fixed point-theorem [17].

Theorem 2.5 (Krasnosel'skii) *Let \mathcal{B} be a Banach space, and let $\mathcal{P} \subset \mathcal{B}$ be a cone in \mathcal{B} . Assume Ω_1 and Ω_2 are open subsets of \mathcal{B} with $0 \in \Omega_1 \subset \bar{\Omega}_1 \subset \Omega_2$, and let*

$$T : \mathcal{P} \cap (\bar{\Omega}_2 \setminus \Omega_1) \rightarrow \mathcal{P}$$

be a completely continuous operator such that, either

- (i) $\|Tu\| \leq \|u\|$, $u \in \mathcal{P} \cap \partial\Omega_1$, and $\|Tu\| \geq \|u\|$, $u \in \mathcal{P} \cap \partial\Omega_2$, or

Ref.

[17] M. A. Krasnosel'skii, *Positive Solutions of Operator Equations*, P. Noordhoff Ltd, Groningen, The Netherlands (1964).

(ii) $\|Tu\| \geq \|u\|$, $u \in \mathcal{P} \cap \partial\Omega_1$, and $\|Tu\| \leq \|u\|$, $u \in \mathcal{P} \cap \partial\Omega_2$.
 Then, T has a fixed point in $\mathcal{P} \cap (\overline{\Omega_2} \setminus \Omega_1)$.

III. POSITIVE SOLUTIONS IN A CONE

In this section, we apply Theorem 2.5 to obtain solutions in a cone (i.e., positive solutions) of (1.1), (1.2).

For our construction, let $\mathcal{B} = \{x : [a, \sigma^2(b)]_{\mathbb{T}} \rightarrow \mathbb{R}\}$ with supremum norm $\|x\| = \sup\{|x(t)| : t \in [a, \sigma^2(b)]_{\mathbb{T}}\}$ and define a cone $\mathcal{P} \subset \mathcal{B}$ by

$$\mathcal{P} = \left\{ x \in \mathcal{B} \mid x(t) \geq 0 \text{ on } [a, \sigma^2(b)]_{\mathbb{T}}, \text{ and } \min_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} x(t) \geq \gamma \|x\| \right\}.$$

For our first result, define positive numbers L_1 and L_2 by

$$L_1 := \max \left\{ \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) \Delta s f_{\infty} \right]^{-1}, \right. \\ \left. \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) q(s) \Delta s g_{\infty} \right]^{-1} \right\},$$

and

$$L_2 := \min \left\{ \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) \Delta s f_0 \right]^{-1}, \right. \\ \left. \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s)) q(s) \Delta s g_0 \right]^{-1} \right\}.$$

Theorem 3.1 *Assume that conditions (A1) – (A3) are satisfied. Then, for each λ satisfying*

$$L_1 < \lambda < L_2, \tag{3.1}$$

there exists a pair (u, v) satisfying (1.1), (1.2) such that $u(x) > 0$ and $v(x) > 0$ on $(a, \sigma^2(b))_{\mathbb{T}}$.

Proof: Let λ be as in (3.1). And let $\epsilon > 0$ be chosen such that

$$\max \left\{ \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) \Delta s (f_{\infty} - \epsilon) \right]^{-1}, \right.$$



$$\left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))q(s)\Delta s (g_{\infty} - \epsilon) \right]^{-1} \} \leq \lambda,$$

and

$$\lambda \leq \min \left\{ \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)\Delta s (f_0 + \epsilon) \right]^{-1}, \right. \\ \left. \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))q(s)\Delta s (g_0 + \epsilon) \right]^{-1} \right\},$$

Define an integral operator $T : \mathcal{P} \rightarrow \mathcal{B}$ by

$$Tu(t) := \lambda \int_a^{\sigma(b)} G(t, s)p(s)f \left(\lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r \right) \Delta s, \quad u \in \mathcal{P}. \quad (3.2)$$

We seek suitable fixed points of T in the cone \mathcal{P} .

By Lemma 2.4, $T\mathcal{P} \subset \mathcal{P}$. In addition, standard arguments show that T is completely continuous.

Now, from the definitions of f_0 and g_0 , there exists $H_1 > 0$ such that

$$f(x) \leq (f_0 + \epsilon)x \quad \text{and} \quad g(x) \leq (g_0 + \epsilon)x, \quad 0 < x \leq H_1.$$

Let $u \in \mathcal{P}$ with $\|u\| = H_1$. We first have from (2.6) and choice of ϵ ,

$$\begin{aligned} & \lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r \\ & \leq \lambda \frac{\beta + \alpha(t - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)g(u(\sigma(r)))\Delta r \\ & \leq \lambda \frac{\beta + \alpha(t - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)(g_0 + \epsilon)u(r)\Delta r \\ & \leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)\Delta r (g_0 + \epsilon)\|u\| \\ & \leq \|u\| \\ & = H_1. \end{aligned}$$

As a consequence, we next have from (2.6) and choice of ϵ ,

$$Tu(t) = \lambda \int_a^{\sigma(b)} G(t, s)p(s)f \left(\lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r \right) \Delta s$$

$$\begin{aligned}
 &\leq \lambda \frac{\beta + \alpha(t-a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\
 &\quad f\left(\lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r\right)\Delta s \\
 &\leq \lambda \frac{\beta + \alpha(t-a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\
 &\quad (f_0 + \epsilon)\lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r\Delta s \\
 &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)(f_0 + \epsilon)H_1\Delta s \\
 &\leq H_1 \\
 &= \|u\|.
 \end{aligned}$$

So, $\|Tu\| \leq \|u\|$. If we set

$$\Omega_1 = \{x \in \mathcal{B} : \|x\| < H_1\},$$

then

$$\|Tu\| \leq \|u\|, \quad \text{for } u \in \mathcal{P} \cap \partial\Omega_1. \quad (3.3)$$

Next, from the definitions of f_∞ and g_∞ , there exists $\bar{H}_2 > 0$ such that

$$f(x) \geq (f_\infty - \epsilon)x \text{ and } g(x) \geq (g_\infty - \epsilon)x, \quad x \geq \bar{H}_2.$$

Let

$$H_2 = \max\left\{2H_1, \frac{\bar{H}_2}{\gamma}\right\}.$$

Let $u \in \mathcal{P}$ and $\|u\| = H_2$. Then,

$$\min_{t \in [\eta, \sigma^2(b)]_{\mathbb{T}}} u(t) \geq \gamma\|u\| \geq \bar{H}_2.$$

Consequently, from (2.7) and choice of ϵ ,

$$\begin{aligned}
 &\lambda \int_a^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r \\
 &\geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_\eta^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)g(u(\sigma(r)))\Delta r \\
 &\geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_\eta^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)(g_\infty - \epsilon)u(r)\Delta r
 \end{aligned}$$

$$\begin{aligned}
 &\geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)(g_{\infty} - \epsilon)\Delta r \gamma \|u\| \\
 &\geq \|u\| \\
 &= H_2.
 \end{aligned}$$

And so, we have from (2.7) and choice of ϵ ,

$$\begin{aligned}
 Tu(\eta) &\geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\
 &\quad f\left(\lambda \int_{\eta}^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r\right)\Delta s \\
 &\geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\
 &\quad (f_{\infty} - \epsilon)\lambda \int_{\eta}^{\sigma(b)} G(\sigma(s), r)q(r)g(u(\sigma(r)))\Delta r \Delta s \\
 &\geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)(f_{\infty} - \epsilon)H_2 \Delta s \\
 &\geq \lambda \gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)(f_{\infty} - \epsilon)H_2 \Delta s \\
 &\geq H_2 \\
 &= \|u\|.
 \end{aligned}$$

Hence, $\|Tu\| \geq \|u\|$. So, if we set

$$\Omega_2 = \{x \in \mathcal{B} : \|x\| < H_2\},$$

then

$$\|Tu\| \geq \|u\|, \quad \text{for } u \in \mathcal{P} \cap \partial\Omega_2. \quad (3.4)$$

Applying Theorem 2.5 to (3.3) and (3.4), we obtain that T has a fixed point $u \in \mathcal{P} \cap (\overline{\Omega_2} \setminus \Omega_1)$. As such, and with v being defined by

$$v(t) = \lambda \int_a^{\sigma(b)} G(t, s)q(s)g(u(\sigma(s)))\Delta s,$$

the pair (u, v) is a desired solution of (1.1), (1.2) for the given λ . The proof is complete.

Prior to our next result, we define positive numbers L_3 and L_4 by

$$L_3 := \max \left\{ \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)\Delta s f_0 \right]^{-1}, \right. \\
 \left. \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s))q(s)\Delta s g_0 \right]^{-1} \right\},$$

and

$$L_4 := \min \left\{ \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)\Delta s f_\infty \right]^{-1}, \right. \\ \left. \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))q(s)\Delta s g_\infty \right]^{-1} \right\}.$$

Theorem 3.2 Assume that conditions (A1) – (A4) are satisfied. Then, for each λ satisfying

$$L_3 < \lambda < L_4, \tag{3.5}$$

there exists a pair (u, v) satisfying (1.1), (1.2) such that $u(x) > 0$ and $v(x) > 0$ on $(a, \sigma^2(b))_{\mathbb{T}}$.

Proof: Let λ be as in (3.5). And let $\epsilon > 0$ be chosen such that

$$\max \left\{ \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_\eta^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)\Delta s (f_0 - \epsilon) \right]^{-1}, \right. \\ \left. \left[\gamma \frac{\beta + \alpha(\eta - a)}{d} \int_\eta^{\sigma(b)} (\sigma^2(b) - \sigma(s))q(s)\Delta s (g_0 - \epsilon) \right]^{-1} \right\} \leq \lambda,$$

and

$$\lambda \leq \min \left\{ \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)\Delta s (f_\infty + \epsilon) \right]^{-1}, \right. \\ \left. \left[\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))q(s)\Delta s (g_\infty + \epsilon) \right]^{-1} \right\}.$$

Let T be the cone preserving, completely continuous operator that was defined by (3.2).

From the definitions of f_0 and g_0 , there exists $H_3 > 0$ such that

$$f(x) \geq (f_0 - \epsilon)x \text{ and } g(x) \geq (g_0 - \epsilon)x, \quad 0 < x \leq H_3.$$

Also, from the definition of g_0 it follows that $g(0) = 0$ and so there exists $0 < H_3 < \overline{H}_3$ such that

$$\lambda g(x) \leq \frac{\overline{H}_3}{\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)\Delta r}, \quad 0 \leq x \leq H_3.$$



Choose $u \in \mathcal{P}$ with $\|u\| = H_3$. Then

$$\begin{aligned}
 & \lambda \int_a^{\sigma(b)} G(\sigma(s), r) q(r) g(u(\sigma(r))) \Delta r \\
 & \leq \lambda \frac{\beta + \alpha(t - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r)) q(r) g(u(\sigma(r))) \Delta r \\
 & \leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r)) q(r) g(u(\sigma(r))) \Delta r \\
 & \leq \frac{\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r)) q(r) \overline{H}_3 \Delta r}{\frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s)) q(s) \Delta s} \\
 & \leq \overline{H}_3.
 \end{aligned}$$

Then, by (2.7)

$$\begin{aligned}
 Tu(\eta) & \geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) \\
 & \quad f \left(\lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(r)) q(r) g(u(\sigma(r))) \Delta r \right) \Delta s \\
 & \geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) \\
 & \quad (f_0 - \epsilon) \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(r)) q(r) g(u(\sigma(r))) \Delta r \Delta s \\
 & \geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) \\
 & \quad (f_0 - \epsilon) \lambda \gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(r)) q(r) (g_0 - \epsilon) \|u\| \Delta r \Delta s \\
 & \geq \lambda \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) (f_0 - \epsilon) \|u\| \Delta s \\
 & \geq \lambda \gamma \frac{\beta + \alpha(\eta - a)}{d} \int_{\eta}^{\sigma(b)} (\sigma^2(b) - \sigma(s)) p(s) (f_0 - \epsilon) \|u\| \Delta s \\
 & \geq \|u\|.
 \end{aligned}$$

So, $\|Tu\| \geq \|u\|$. If we put

$$\Omega_3 = \{x \in \mathcal{B} : \|x\| < H_3\},$$

then

$$\|Tu\| \geq \|u\|, \quad \text{for } u \in \mathcal{P} \cap \partial\Omega_3. \quad (3.6)$$

Next, by definitions of f_∞ and g_∞ , there exists \bar{H}_4 such that

$$f(x) \leq (f_\infty + \epsilon)x \quad \text{and} \quad g(x) \leq (g_\infty + \epsilon)x, \quad x \geq \bar{H}_4$$

Clearly, since g_∞ is assumed to be a positive real number, it follows that g is unbounded at ∞ , and so, there exists $\tilde{H}_4 > \max\{2H_3, \bar{H}_4\}$ such that $g(x) \leq g(\tilde{H}_4)$, for $0 < x \leq \tilde{H}_4$.

Set

$$f^*(t) = \sup_{a \leq s \leq t} f(s), \quad g^*(t) = \sup_{a \leq s \leq t} g(s), \quad \text{for } t \geq 0.$$

Clearly f^* and g^* are nondecreasing real valued functions for which it holds

$$\lim_{x \rightarrow \infty} \frac{f_i^*(x)}{x} = f_\infty, \quad \lim_{x \rightarrow \infty} \frac{g_i^*(x)}{x} = g_\infty.$$

Hence, there exists H_4 such that $f^*(x) \leq f^*(H_4)$, $g^*(x) \leq g^*(H_4)$, for $0 < x \leq H_4$.

Choosing $u \in \mathcal{P}$ with $\|u\| = H_4$, we have

$$\begin{aligned} Tu(t) &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\ &\quad f\left(\lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)g(u(\sigma(r)))\Delta r\right)\Delta s \\ &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\ &\quad f^*\left(\lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)g(u(\sigma(r)))\Delta r\right)\Delta s \\ &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\ &\quad f^*\left(\lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)g^*(u(\sigma(r)))\Delta r\right)\Delta s \\ &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\ &\quad f^*\left(\lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)g^*(H_4)\Delta r\right)\Delta s \\ &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s) \\ &\quad f^*\left(\lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(r))q(r)(g_\infty + \epsilon)H_4\Delta r\right)\Delta s \\ &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^{\sigma(b)} (\sigma^2(b) - \sigma(s))p(s)f^*(H_4)\Delta s \end{aligned}$$

$$\begin{aligned} &\leq \lambda \frac{\beta + \alpha(\sigma^2(b) - a)}{d} \int_a^\sigma (\sigma^2(b) - \sigma(s))p(s)(f_\infty + \epsilon)H_4\Delta s \\ &\leq H_4 \\ &= \|u\|, \end{aligned}$$

and so $\|Tu\| \leq \|u\|$. For this case, if we let

$$\Omega_4 = \{x \in \mathcal{B} : \|x\| < H_4\},$$

then

$$\|Tu\| \leq \|u\|, \quad \text{for } u \in \mathcal{P} \cap \partial\Omega_4. \quad (3.7)$$

Application of part (ii) of Theorem 2.5 yields a fixed point u of T belonging to $\mathcal{P} \cap (\overline{\Omega}_4 \setminus \Omega_3)$, which in turn yields a pair (u, v) satisfying (1.1), (1.2) for the chosen value of λ . The proof is complete.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

Application of Laplace Transform

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Abstract - The present discounted value equation in finance has a broad range of uses and may be applied to various areas of finance including corporate finance, banking finance and investment finance etc. The basic premise of present discounted value is the time value money. Not many analytic solutions exist for present discounted value problems but by using Laplace transform we can deduce some of the closed form solutions quite easily. In this note we show how present discounted value in finance related to Laplace transforms. Also we discuss on the present value rules for the elementary functions and the general properties of the Laplace transform. And we will focus on the application of time derivative property using Laplace transforms to each present value rule.

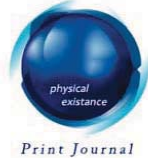
Keywords : Present discounted value, cash flow, perpetuity, Time derivative, Laplace transform.

GJSFR-F Classification : MSC 2010: 44A10



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Application of Laplace Transform

Dr. N. A. Patil^α & Vijaya N. Patil^σ

Abstract - The present discounted value equation in finance has a broad range of uses and may be applied to various areas of finance including corporate finance, banking finance and investment finance etc. The basic premise of present discounted value is the time value of money. Not many analytic solutions exist for present discounted value problems but by using Laplace transform we can deduce some of the closed form solutions quite easily. In this note we show how present discounted value in finance related to Laplace transforms. Also we discuss on the present value rules for the elementary functions and the general properties of the Laplace transform. And we will focus on the application of time derivative property using Laplace transforms to each present value rule.

Keywords : Present discounted value, cash flow, perpetuity, Time derivative, Laplace transform.

I. INTRODUCTION

During the past few decades, methods based on integral transforms, in particular, the Laplace transforms, are being increasingly employed in mathematics, physics, mechanics and other engineering sciences. Laplace transforms have a wide variety of applications in the solution of differential, integral and difference equations. It is much less used in financial engineering. One reason is technical: not many people know that all that they need to do is to make simple calculations in the Laplace domain.

The outline of this note is as follows –

In section 1 we show the relation between present discounted value and Laplace transforms.

In section 2 we identified the present value rules for each of the cash flow.

In section 3 we discussed the general properties of Laplace transforms with present value rules.

In section 4 we show the application of time derivative property to each of the present value rules.

II. RELATION BETWEEN PRESENT VALUE AND LAPLACE TRANSFORM

The Present value of a series of payments given by,

$$(PV)_t = \sum_{t=1}^T \frac{C(t)}{(1+r)^t} \quad (1)$$

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Where, $(PV)_t$ = Present discounted value at time t

$C(t)$ = Cash flow

r = rate of discount

t = Period

Here we assume the Present value with continuous compounding. It is the current value of a stream of cash flows. In other words, it is the amount that we would be willing to pay today in order to receive a cash flow or a series of them in the future.

Now by using an exponential series we can write equation (1) as,

$$(PV)_t = \sum_{t=1}^T e^{-rt} C(t) \quad (2)$$

In the limiting case replacing summation to an integral, equation (2) can be written as

$$(PV)_r = \int_0^T e^{-rt} C(t) dt \quad (3)$$

Again here T is some finite quantity. So if we consider as $T \rightarrow \infty$, equation (3) will become

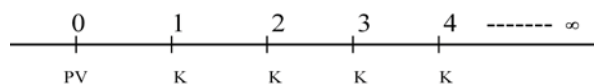
$$(PV)_r = \int_0^{\infty} e^{-rt} C(t) dt \quad (4)$$

This equation is the exact replica of Laplace transform in mathematics.

Therefore,
$$(PV)_r = L[C(t)] \quad (5)$$

III. LAPLACE TRANSFORMS AND PRESENT VALUE RULES FOR SOME CASH FLOWS

Using Present value equation: Consider the case of constant cash payment K made at the end of each year at interest rate r as shown in the following time line,



Here the cash flow is continuous forever. Therefore the Present value is given by an infinite geometric series:

$$PV = \frac{K}{1+r} + \frac{K}{(1+r)^2} + \frac{K}{(1+r)^3} + \dots \quad (6)$$

Dividing both sides by $(1+r)$ we get,

$$\frac{PV}{1+r} = \frac{K}{(1+r)^2} + \frac{K}{(1+r)^3} + \frac{K}{(1+r)^4} + \dots \quad (7)$$

Subtracting equation (7) from (6) we get,

$$PV - \frac{PV}{1+r} = \frac{K}{1+r}$$

On solving we get the Present value of perpetuity.

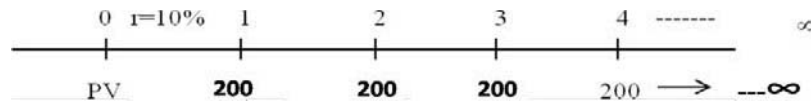
Using Laplace transform equation: If the cash flow is constant say K then the Present discounted value of a stream at interest rate r is given by,

$$L[K] = K \int_0^{\infty} e^{-rt} dt = \frac{K}{r}$$

This is the same formula as above.

For example: An insurance company has just launched a security that will pay Rs.200 indefinitely, starting the first payment next year. How much should this security be worth today if the appropriate return is 10% ?

We solve this example by using the time line,



$$PV = \frac{K}{r} = \frac{200}{0.10} = Rs. 2000$$

Using Present value equation: Consider the payments in perpetuity increases at a certain growth rate g as shown on the time line:

The Present value of a growing perpetuity can be written as the following infinite series-

$$PV = \frac{K}{1+r} + \frac{K(1+g)}{(1+r)^2} + \frac{K(1+g)^2}{(1+r)^3} + \dots \quad (8)$$

Multiplying both sides by $\frac{(1+g)}{(1+r)}$. Hence we get,

$$PV \frac{(1+g)}{(1+r)} = \frac{K(1+g)}{(1+r)^2} + \frac{K(1+g)^2}{(1+r)^3} + \frac{K(1+g)^3}{(1+r)^4} + \dots \quad (9)$$

Subtracting equation (9) from (8) we get,

$$PV - \frac{PV(1+g)}{1+r} = \frac{K}{1+r}$$

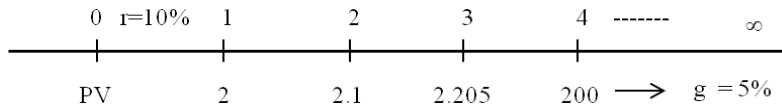
On solving we get the Present value of a growing perpetuity.

Using Laplace transform equation: For an exponential or geometric cash flow the Present discounted value of a stream growing at rate g, is given by:

$$L[C(t)] = K L[e^{gt}] = \frac{K}{r-g} \text{ if } r > g \quad (10)$$

This is the geometric growth stream or Present value of growing perpetuity having cash flow after the first period divided by the difference between the discount rate and the growth rate and the growth rate must be less than the interest rate.

For example: A company is expected to pay Rs.2 of dividend per share that will increase 5% forever. If investors require 10% return on the company's stocks, how much should investors pay for the stocks? The cash flows are as follows:



$$PVG = \frac{K}{r - g} = \frac{2}{0.10 - 0.05} = Rs. 40 \quad (11)$$

For an arithmetic cash flow the Present discounted value of a stream at rate r , is given by:

$$L[C(t)] = K L[t] = \frac{K}{r^2} \quad (12)$$

This shows that an arithmetic growth stream is equivalent to receiving one consol per period in perpetuity. This rule is widely used in finance for solving Present value.

The above rules are commonly used transforms but more useful are the general properties of the Laplace transforms in an algebraic fashion. Let us look at some of the main properties.

Property 1: Linearity: The Laplace Transform is a linear operator. Hence if the Laplace Transforms of the cash flows $f(t)$ and $g(t)$ both exist then we have for any arbitrary constants (a, b) that:

$$L\{a f(t) + b g(t)\} = a L\{f(t)\} + b L\{g(t)\} \quad (13)$$

This property allows us to deduce more complex transforms to simple transforms.

Property 2: Geometric scaling:

$$L\{e^{\alpha t} c(t)\} = V(r - \alpha) \quad \text{for } \alpha < r \quad (14)$$

This property shows that the scaling a cash flow by geometric term is equivalent to corresponding reduction in the rate of discount.

Property 3: Multiplication by t :

$$L\{t C(t)\} = -V'(r) \quad (15)$$

We can confirm this property by using the derivative of exponential function.

Property 4: Time shifting:

$$\left. \begin{array}{l} L\{C(a+bt)\} \text{ for } t \geq a/b \\ 0 \quad \quad \quad \text{for } t < a/b \end{array} \right\} = e^{ra/b} \left(\frac{1}{b}\right) P(r/b) \quad (16)$$

This property applies the change of variable theorem of integral calculus and helpful for finding cash flows with altered time schedules.

Property 5: Time derivative:

$$L\{C'(t)\} = r L\{C(t)\} - C(0) \quad (17)$$

This property identifies a fundamental linear relationship between Laplace transform for cash flows and their time derivatives. For the confirmation we use integration by parts:

$$\int u dv = uv - \int v du$$

When we evaluate the integral over relevant range 0 to ∞ for the Laplace transform and impose a standard assumption in present value problems that the marginal present value of the cash flow vanishes as t gets large.

All present value rules of section second can be derived from this time derivative property of Laplace transform and hence having particular significance in finance. For the confirmation we rewrite the time-derivative property by using notation

$$(PV)_r = L[C(t)], \text{ as:}$$

$$L[C(t)] = \frac{C(0)}{r} + \frac{L[C'(t)]}{r} \quad (18)$$

Apply the property, for some cash flows.

Ex: For $S(t) = K \Rightarrow S(0) = K$ & $S'(t) = 0$, we get the consol rule by using property 5.

Alternatively, each asset is valued as if its cash flow were projected at a constant level equal to the current rate plus the present value of the time derivative of the cash flow.

For geometric cash flow (2.I)

$$S(t) = e^{\alpha t} \Rightarrow S(0) = 1 \text{ \& } S'(t) = \alpha e^{\alpha t}, \text{ we get}$$

$$L[e^{\alpha t}] = \frac{1}{\theta} + \frac{\alpha L[e^{\alpha t}]}{\theta} \quad (19)$$

Alternatively, we could combine the consol rule with property 2 to establish the rule for geometric growth.

Similarly we can derive the rule for arithmetic growth by using equation (19) or combining the consol rule and property 3.

IV. CONCLUSION

In this article we have presented the close relationship between present discounted value in finance and Laplace transform. We can solve the present discounted value examples within a few minutes by using Laplace equation method. The result seems to be new & to have a potential to increase the practical utility of Laplace transform especially in finance. However it is important to notice that frequency domain is possible appreciate also in the real world & applied in the areas like economics or finance. But the Laplace transform is the big source for present discounted value function to illustrate the enhanced problems.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

Formation of a Summation Formula Enmeshed with Hypergeometric Function

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Keywords : *Gaussian Hypergeometric function , Contiguous function, Recurrence relation, Bailey summation theorem and Legendre duplication formula.*

GJSFR-F Classification : *MSC 2010: 65B10, 33D60*



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I. INTRODUCTION

Generalized Gaussian hypergeometric function of one variable is defined by

$${}_A F_B \left[\begin{matrix} a_1, a_2, \dots, a_A ; \\ b_1, b_2, \dots, b_B ; \end{matrix} z \right] = \sum_{k=0}^{\infty} \frac{(a_1)_k (a_2)_k \dots (a_A)_k z^k}{(b_1)_k (b_2)_k \dots (b_B)_k k!}$$

or

$${}_A F_B \left[\begin{matrix} (a_A) ; \\ (b_B) ; \end{matrix} z \right] \equiv {}_A F_B \left[\begin{matrix} (a_j)_{j=1}^A ; \\ (b_j)_{j=1}^B ; \end{matrix} z \right] = \sum_{k=0}^{\infty} \frac{((a_A))_k z^k}{((b_B))_k k!} \quad (1)$$

where the parameters b_1, b_2, \dots, b_B are neither zero nor negative integers and A, B are non-negative integers.

Contiguous Relation[E. D. p.51(10), Andrews p.363(9.16)] is defined as follows

$$(a-b) {}_2 F_1 \left[\begin{matrix} a, b ; \\ c ; \end{matrix} z \right] = a {}_2 F_1 \left[\begin{matrix} a+1, b ; \\ c ; \end{matrix} z \right] - b {}_2 F_1 \left[\begin{matrix} a, b+1 ; \\ c ; \end{matrix} z \right] \quad (2)$$

Recurrence relation of gamma function is defined as follows

$$\Gamma(z+1) = z \Gamma(z) \quad (3)$$

Legendre duplication formula[Bells & Wong p.26(2.3.1)] is defined as follows

$$\sqrt{\pi} \Gamma(2z) = 2^{(2z-1)} \Gamma(z) \Gamma\left(z + \frac{1}{2}\right) \quad (4)$$

$$\Gamma\left(\frac{1}{2}\right) = \sqrt{\pi} = \frac{2^{(b-1)} \Gamma\left(\frac{b}{2}\right) \Gamma\left(\frac{b+1}{2}\right)}{\Gamma(b)} \tag{5}$$

$$= \frac{2^{(a-1)} \Gamma\left(\frac{a}{2}\right) \Gamma\left(\frac{a+1}{2}\right)}{\Gamma(a)} \tag{6}$$

Bailey summation theorem [Prud, p.491(7.3.7.8)] is defined as follows

$${}_2F_1 \left[\begin{matrix} a, 1-a & ; & 1 \\ c & ; & 2 \end{matrix} \right] = \frac{\Gamma\left(\frac{c}{2}\right) \Gamma\left(\frac{c+1}{2}\right)}{\Gamma\left(\frac{c+a}{2}\right) \Gamma\left(\frac{c+1-a}{2}\right)} = \frac{\sqrt{\pi} \Gamma(c)}{2^{c-1} \Gamma\left(\frac{c+a}{2}\right) \Gamma\left(\frac{c+1-a}{2}\right)} \tag{7}$$

II. MAIN RESULT OF SUMMATION FORMULA

$$\begin{aligned} & {}_2F_1 \left[\begin{matrix} a & , & -a-48 & ; & 1 \\ & c & & ; & 2 \end{matrix} \right] \\ &= \frac{\sqrt{\pi} \Gamma(c)}{2^{c+48}} \left[\frac{-3638347904750543085030062414561280000a}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \right. \\ & \quad + \frac{4230534459144635193233439012777984000a^2}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{-1773530111269024379361184888978329600a^3}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{-26072127788223783027931526767525632a^5 - 140082855115327355120120725082688a^6}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{113210267032593419892743400611520a^7 - 1591308493265001080090252839184a^8}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{-249488180140450905945277350672a^9 + 1896611354738752958221477012a^{10}}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{340576268631980322449458500a^{11} + 2232393636721350709468081a^{12}}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{-221961189928545392626392a^{13} - 4488938418788950118258a^{14}}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{26237742565110055980a^{15} + 1726308193141046911a^{16} + 16724033867495328a^{17}}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \\ & \quad + \frac{-67838656895768a^{18} - 2424088437540a^{19} - 16239228929a^{20} - 12952632a^{21} + 295702a^{22}}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} \\ & \quad + \frac{1140a^{23} + a^{24} + 3638347904750568937046801299537920000c}{\Gamma\left(\frac{c-a+1}{2}\right) \Gamma\left(\frac{c+a+48}{2}\right)} + \end{aligned}$$

$$\begin{aligned}
 & + \frac{-11023860306689859165632917200175104000ac}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{8296047725075578715435877350611353600a^2c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2589060527415141315999913883524792320a^3c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{366873883396175498327904190658199552a^4c - 17957990686264683551051799087968256a^5c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-783414259355080511741474195684352a^6c + 84340873453502002421398843682048a^7c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{973370298518103355293671126656a^8c - 160687148751021926409412982272a^9c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2248983658767185268213909888a^{10}c + 149172577626311143061653392a^{11}c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{3615991694923403891801160a^{12}c - 40237584058803766916008a^{13}c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2145579120612795343784a^{14}c - 15847495431380520424a^{15}c + 294522342618394704a^{16}c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{5916283797667056a^{17}c + 27173155233840a^{18}c - 187374214720a^{19}c - 2438354072a^{20}c}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-8088520a^{21}c - 1416a^{22}c + 24a^{23}c + 6793325847545320511366130680463360000c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-12135827068916077894200206118617088000ac^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{6753675210879510270039714594306785280a^2c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-1624869911805656070962191215597797376a^3c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{171251591774737199593026229784064000a^4c^2 - 4093063242641347774786447363015680a^5c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-511149924510502042072739692658944a^6c^2 + 22504067285994043924692302863232a^7c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{913186425471360033494807245312a^8c^2 - 33804718262755533682100484480a^9c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-1295637035432751803115296208a^{10}c^2 + 15890879633892055504532184a^{11}c^2}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} +
 \end{aligned}$$

$$\begin{aligned}
& + \frac{1022450363916086376744200a^{12}c^2 + 5608077520737975363080a^{13}c^2}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-278653737771626178104a^{14}c^2 - 4752455855653094160a^{15}c^2 - 5301794887227696a^{16}c^2}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{495348526392720a^{17}c^2 + 4527490455488a^{18}c^2 + 8270782520a^{19}c^2 - 69085016a^{20}c^2}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-326040a^{21}c^2 - 312a^{22}c^2 + 5613309455109683097975789723844608000c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-7268957877902293679747953623910318080ac^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{3148476881911767042344731733158526976a^2c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-591776472395510749153521515998937088a^3c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{45128923170400422228420079851380736a^4c^3 + 5555930291075194182761162530816a^5c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-150716557255326505370710216430080a^6c^3 + 2217081280693047664145316025088a^7c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{254900430630476339868277406976a^8c^3 - 1908319886666337016042402944a^9c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-255978165939209274666102816a^{10}c^3 - 1370080885460819936032720a^{11}c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{113902943230614008414944a^{12}c^3 + 1870509542899797944144a^{13}c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-7980655855696546880a^{14}c^3 - 425296980132412960a^{15}c^3 - 3111050730435648a^{16}c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{7889860890912a^{17}c^3 + 200089649760a^{18}c^3 + 804243440a^{19}c^3 + 352352a^{20}c^3 - 2288a^{21}c^3}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{2768586081553372140200792041508044800c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-2788045206948865839414630121447882752ac^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{962315653870080299754170370182283264a^2c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-141439234805306232175165147182489600a^3c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} +
\end{aligned}$$

$$\begin{aligned}
& + \frac{7350569555757964653087742802558976a^4c^4 + 199068791705835251665671359802880a^5c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-25294347834251193603125177586944a^6c^4 - 147437993181774554303313826560a^7c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{36251936778359362132532742528a^8c^4 + 346029886905457750751639520a^9c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-24523514138602855095025488a^{10}c^4 - 447590500623950105662240a^{11}c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{4712185069329296679120a^{12}c^4 + 177404400410212015040a^{13}c^4 + 932298241605633760a^{14}c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-14323139972581440a^{15}c^4 - 193197305325024a^{16}c^4 - 553302990240a^{17}c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{2501298800a^{18}c^4 + 15215200a^{19}c^4 + 16016a^{20}c^4}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{923636481234908158042872857449463808c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-747023445605858669273498410297589760ac^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{207668786206542828581701432891146240a^2c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-23663020067855212035887457634222080a^3c^5 + 732904567621987230351419111325696a^4c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{49903353805450335424368547782656a^5c^5 - 2578893746999530831934304224256a^6c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-66577688249629584056310230016a^7c^5 + 2874347583304554423474949632a^8c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{74463995548879100021824512a^9c^5 - 1062528271193599593294912a^{10}c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-42807675912180201294272a^{11}c^5 - 124023155193475329792a^{12}c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{7667675875623731968a^{13}c^5 + 86927587193935488a^{14}c^5 + 5208901746048a^{15}c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-4594408955136a^{16}c^5 - 22868541696a^{17}c^5 - 15951936a^{18}c^5 + 64064a^{19}c^5}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{223330745897982709518529548993429504c^6 - 147472358833668973831737490003722240ac^6}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} +
\end{aligned}$$



$$\begin{aligned}
 & + \frac{33122202957018388352610598824574976a^2c^6 - 2871630447882289896332967060062208a^3c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{34441427528122369258470645260288a^4c^6 + 6831394798548698327794234414080a^5c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-150155381817771045217714537472a^6c^6 - 8692060148779265084777793024a^7c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{108392221996820991579377664a^8c^6 + 6575336744109767454300480a^9c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{8164199124092113502400a^{10}c^6 - 2080877086126562051328a^{11}c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-21558516560807978752a^{12}c^6 + 121341101729527680a^{13}c^6 + 2983367975812224a^{14}c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{12147729834240a^{15}c^6 - 32015343360a^{16}c^6 - 273873600a^{17}c^6 - 320320a^{18}c^6}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{40877151497014378839771354551549952c^7 - 22219502427899702338520014030635008ac^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{4022878448009248326347020243828736a^2c^7 - 256858653782602724496822439018496a^3c^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-1641647966593763191402106634240a^4c^7 + 613388087046982794870551930880a^5c^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2367132688592308795171872768a^6c^7 - 654998067407893710964829184a^7c^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-1523650746127540570315776a^8c^7 + 329539524916344426313344a^9c^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{3425646083934616768512a^{10}c^7 - 50038536300856097280a^{11}c^7 - 974138332318255104a^{12}c^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2184451291486464a^{13}c^7 + 45506452992000a^{14}c^7 + 291087413760a^{15}c^7 + 283345920a^{16}c^7}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-823680a^{17}c^7 + 5832594412243055928973325422821376c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2618050661653803280443469114376192ac^8 + 379493114458898267502470696665088a^2c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-16864207890780906911526743900160a^3c^8 - 447720321025657093838032576512a^4c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{37986129070713604063172198400a^5c^8 + 389769280200184576457945088a^6c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} +
 \end{aligned}$$

$$\begin{aligned}
& + \frac{-31372314252822465838755840a^7c^8 - 411109178500644611480064a^8c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{9363608406942489722880a^9c^8 + 187215352358190114816a^{10}c^8 - 146424613495480320a^{11}c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-21752854863387648a^{12}c^8 - 123540894597120a^{13}c^8 + 173781104640a^{14}c^8}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{2503987200a^{15}c^8 + 3294720a^{16}c^8 + 662317910051181926669156716904448c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-245400115796071532958765453148160ac^9 + 28153238255209821490041472417792a^2c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-781005645276012502462741086208a^3c^9 - 42641895620379018617550274560a^4c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{1612351973725244164720885760a^5c^9 + 39662648103197518265810944a^6c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-926453550912364548243456a^7c^9 - 23555359060666652876800a^8c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{100903225406804254720a^9c^9 + 5299360980595089408a^{10}c^9 + 24692042573406208a^{11}c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-227873864458240a^{12}c^9 - 1996213739520a^{13}c^9 - 2571345920a^{14}c^9 + 5857280a^{15}c^9}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{60743744451907338804193681997824c^{10} - 18518348550706651821965710458880ac^{10}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{1653100889178452682493737828352a^2c^{10} - 21740721608271376099456843776a^3c^{10}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-2600813329461903794303729664a^4c^{10} + 42802618860435796873871360a^5c^{10}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{2038210720517206870835200a^6c^{10} - 12463849823864973467648a^7c^{10}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{-758892596094305001472a^8c^{10} - 2575932665689743360a^9c^{10} + 83354346743635968a^{10}c^{10}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{679546442358784a^{11}c^{10} - 290098192384a^{12}c^{10} - 13243310080a^{13}c^{10} - 19914752a^{14}c^{10}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{4546715789559603603424633946112c^{11} - 1133943987191935145013892087808ac^{11}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
& + \frac{76868622553074562239414927360a^2c^{11} - 6333853945783824633298944a^3c^{11}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} +
\end{aligned}$$



$$\begin{aligned}
 & + \frac{-112614782541939583740936192a^4c^{11} + 383669909202609748246528a^5c^{11}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{67555227087049795731456a^6c^{11} + 197559515798431514624a^7c^{11}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-1499874634184294400a^8c^{11} - 121790210462269440a^9c^{11} + 596723478134784a^{10}c^{11}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{8090331791360a^{11}c^{11} + 13534789632a^{12}c^{11} - 25346048a^{13}c^{11}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{279714662512116971474262163456c^{12} - 56599454327811805687654121472ac^{12}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{2812155919136601281891139584a^2c^{12} + 34239457060775837590487040a^3c^{12}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-3555265567506078642012160a^4c^{12} - 21202879517144336793600a^5c^{12}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{1500614958719249924096a^6c^{12} + 13623195491897671680a^7c^{12} - 170239161849888768a^8c^{12}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-2179202261483520a^9c^{12} - 1109168406528a^{10}c^{12} + 43341742080a^{11}c^{12} + 76038144a^{12}c^{12}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{14202828556049274882760900608c^{13} - 2306244763812047662339850240ac^{13}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{79549699725335858543329280a^2c^{13} + 1970351955192820966359040a^3c^{13}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-81197049931788972195840a^4c^{13} - 1121385597257318400000a^5c^{13}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{21460893335812374528a^6c^{13} + 318310898510266368a^7c^{13} - 703360552796160a^8c^{13}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-20253520035840a^9c^{13} - 44148916224a^{10}c^{13} + 70189056a^{11}c^{13}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{596172495237225865390587904c^{14} - 76594063955412337515560960ac^{14}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{1673197695444638048976896a^2c^{14} + 67447668828029406412800a^3c^{14}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{-1280086116620532449280a^4c^{14} - 28565196119329013760a^5c^{14}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} + \\
 & + \frac{163144044077973504a^6c^{14} + 4184336248995840a^7c^{14} + 6613981593600a^8c^{14}}{\Gamma(\frac{c-a+1}{2}) \Gamma(\frac{c+a+48}{2})} +
 \end{aligned}$$

$$\begin{aligned}
& + \frac{-90493747200a^9c^{14} - 190513152a^{10}c^{14} + 20667778097297440376881152c^{15}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-2063368899940734098997248ac^{15} + 23616227634398246207488a^2c^{15}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{1609625006757485477888a^3c^{15} - 11946606167053565952a^4c^{15} - 457413975935680512a^5c^{15}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-102213608275968a^6c^{15} + 31611720302592a^7c^{15} + 91954348032a^8c^{15} - 127008768a^9c^{15}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{589580948824343210622976c^{16} - 44694485078557324214272ac^{16}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{136597103714150121472a^2c^{16} + 27691358161274142720a^3c^{16} - 16081668010082304a^4c^{16}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-4728686302986240a^5c^{16} - 14549489418240a^6c^{16} + 120658329600a^7c^{16} + 317521920a^8c^{16}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{13745163084312984158208c^{17} - 767797631994193510400ac^{17}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-2697012025875234816a^2c^{17} + 340925659342700544a^3c^{17} + 1206869439283200a^4c^{17}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-29915346370560a^5c^{17} - 122376683520a^6c^{17} + 149422080a^7c^{17}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{259068608527627976704c^{18} - 10245061664235847680ac^{18} - 86103626336960512a^2c^{18}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{2895000352849920a^3c^{18} + 16701453762560a^4c^{18} - 99365683200a^5c^{18} - 348651520a^6c^{18}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{3884083664023191552c^{19} - 102879574253109248ac^{19} - 1204310832578560a^2c^{19}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{15688841297920a^3c^{19} + 100631838720a^4c^{19} - 110100480a^5c^{19} + 45210938613170176c^{20}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-739630552973312ac^{20} - 9921293713408a^2c^{20} + 46022000640a^3c^{20} + 242221056a^4c^{20}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{393630531452928c^{21} - 3495365181440ac^{21} - 46552580096a^2c^{21} + 46137344a^3c^{21}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{2410953048064c^{22} - 9164554240ac^{22} - 96468992a^2c^{22} + 9261023232c^{23}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \\
& + \frac{-8388608ac^{23} + 16777216c^{24}}{\Gamma\left(\frac{c-a+1}{2}\right)\Gamma\left(\frac{c+a+48}{2}\right)} + \frac{20007974164906320568399715106816000000}{\Gamma\left(\frac{c-a}{2}\right)\Gamma\left(\frac{c+a+49}{2}\right)} +
\end{aligned}$$

$$\begin{aligned}
& + \frac{-36922405754626178036890405761515520000a}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{22065278387991542066307045169926144000a^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-5809841151573544819649225637045350400a^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{679890009913045305132937518437944320a^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-18472923622042793122488427319724288a^5 - 2560574556753342210803824950095424a^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{132219440896221918244499646073152a^7 + 5904047252981777731527533921776a^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-270614639795328768343142065008a^9 - 11480344457959138778418103724a^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{191451864133347179813345532a^{11} + 13328712827913082236490801a^{12}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{73453466060412362003352a^{13} - 5838069209154766013234a^{14}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-119007946686312580908a^{15} - 64104578911758209a^{16} + 23731154067653472a^{17}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{293580605335336a^{18} + 706255392612a^{19} - 12674658689a^{20} - 115467528a^{21} - 326954a^{22}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{12a^{23} + a^{24} + 51436861851110719507236588764528640000c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-70899646160167879803660918205710336000ac}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{33418551215289286318709655103183257600a^2c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-6936716063495713404096607576151162880a^3c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{592162566124636481520902123469422592a^4c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-150763178991495629513120805163008a^5c - 2543564167189971442636762196069376a^6c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{45503143417288234509441538577152a^7c + 5720189897921291255677630092928a^8c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
\end{aligned}$$

$$\begin{aligned}
& + \frac{-57850481183031873845533379072a^9c - 8079408263510288034016072320a^{10}c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-42658602715638740094989712a^{11}c + 5457516920998845643189704a^{12}c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{105080628717885503259304a^{13}c - 718597316917131560744a^{14}c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-42146779055614771736a^{15}c - 394633823765699760a^{16}c + 1815118515463440a^{17}c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{59605259080752a^{18}c + 390797915200a^{19}c + 260642536a^{20}c - 7396664a^{21}c - 27912a^{22}c}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-24a^{23}c + 53879676136106962853910284379095040000c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-58701850922925277436346363011137536000ac^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{22272862407231657294514826679624007680a^2c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-3645281850231197173180825687195828224a^3c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{213875375121643552409476002283352064a^4c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{6491596305419248116184393185321984a^5c^2 - 964599604570031846728877390420224a^6c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-5682533971627991823205758376832a^7c^2 + 1885588931111013766114178951680a^8c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{19394898395592969380873022336a^9c^2 - 1851518450114596444564235088a^{10}c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-39083335675241940394365144a^{11}c^2 + 580438231442930787182984a^{12}c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{25546912919341268358904a^{13}c^2 + 159272816221878580936a^{14}c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-3904286889491205360a^{15}c^2 - 70854972633732144a^{16}c^2 - 286985878971792a^{17}c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{2582951379008a^{18}c^2 + 30021551560a^{19}c^2 + 93509416a^{20}c^2 - 3432a^{21}c^2 - 312a^{22}c^2}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{32211699088981901572499271667679232000c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
\end{aligned}$$



$$\begin{aligned}
 & + \frac{-28512000389819083595321676326366085120ac}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{8790460634142328234758745607006846976a^2c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-1124467001863561580537016353123794944a^3c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{39708713605091532228302178253774848a^4c^3 + 3072311886315380216747428668760064a^5c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-186888540899160276337102312129024a^6c^3 - 5460277329981356928340767544064a^7c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{292228466284140197571175445760a^8c^3 + 8719189987830229672642717824a^9c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-165175187685162073177772832a^{10}c^3 - 7770572051846792035610672a^{11}c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-22875449892544073945120a^{12}c^3 + 2404075406651070648496a^{13}c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{35321722511563533760a^{14}c^3 - 16711413272084960a^{15}c^3 - 4275303741893184a^{16}c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-34756964834592a^{17}c^3 - 42598876320a^{18}c^3 + 641040400a^{19}c^3 + 2658656a^{20}c^3}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{2288a^{21}c^3 + 12673569486529179630190013064413184000c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-9247136634374056328148193250852732928ac^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{2322904856073730886919387342994735104a^2c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-227997454418110745446600708021592064a^3c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{3169731737784104346723877333917696a^4c^4 + 717153243408099923350630112145920a^5c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-18894426862062122171165144950016a^6c^4 - 1256503503761835993387006160128a^7c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{20005312090114578497340010368a^8c^4 + 1399506341758816898106208800a^9c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{1107293475704607644242608a^{10}c^4 - 714829981963561507018976a^{11}c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
 \end{aligned}$$

$$\begin{aligned}
 & + \frac{-9234815852307223978800a^{12}c^4 + 80197368741037387840a^{13}c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{2582041372038150880a^{14}c^4 + 15258431854585920a^{15}c^4 - 78560643024864a^{16}c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-1260904925280a^{17}c^4 - 4363799440a^{18}c^4 + 160160a^{19}c^4 + 16016a^{20}c^4}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{3549767559427168161601437434501922816c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-2151099368274358759726535915983601664ac^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{439330611840452027347549954963931136a^2c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-32045308509021985102462728518369280a^3c^5 - 230117021186821381378230594748416a^4c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{103351131588547583108228728397824a^5c^5 - 526153805312068114419158172672a^6c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-156554337016017490934865859584a^7c^5 - 342370320670336052097249792a^8c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{120430698898953310510476288a^9c^5 + 1514876828200500103475136a^{10}c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-31496137302336777598528a^{11}c^5 - 792727673743617900288a^{12}c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-2214018290093978368a^{13}c^5 + 83755863780580992a^{14}c^5 + 874309907382912a^{15}c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{1577740260096a^{16}c^5 - 16154121984a^{17}c^5 - 74378304a^{18}c^5 - 64064a^{19}c^5}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{743559650440452121302608512460783616c^6 - 375380418098956909389583998592745472ac^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{61934956167026537832602378208935936a^2c^6 - 3176115614993973085877719175380992a^3c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-97570846916520135171755792490496a^4c^6 + 9891668493540563867446482152448a^5c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{116489775146410674839857159168a^6c^6 - 11974895925870999956260747776a^7c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-187271227517652896016043008a^8c^6 + 5743465474229360858844864a^9c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
 \end{aligned}$$



$$\begin{aligned}
& + \frac{144554417686389868597440a^{10}c^6 - 230607270695987657472a^{11}c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-33251222890838208256a^{12}c^6 - 277238139202856832a^{13}c^6 + 778197251255424a^{14}c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{20174474960640a^{15}c^6 + 78548870400a^{16}c^6 - 2882880a^{17}c^6 - 320320a^{18}c^6}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{120419109249818260699006806855254016c^7 - 50643920204146384742582467755507712ac^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{6678206027798100901235978413015040a^2c^7 - 216349934068120128428111178465280a^3c^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-13849802762834607412635933622272a^4c^7 + 634544063521152036454215260160a^5c^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{18576717170330425255016964096a^6c^7 - 556528190887633754610422784a^7c^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-17422709590477274298209280a^8c^7 + 107962987229429307820416a^9c^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{7038701856856382828544a^{10}c^7 + 44211950709134866944a^{11}c^7 - 681094885538715648a^{12}c^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-9700638027015936a^{13}c^7 - 23794160332800a^{14}c^7 + 184617438720a^{15}c^7 + 955468800a^{16}c^7}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{823680a^{17}c^7 + 15435801243454637955415444937506816c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-5394264531666420310950929291542528ac^8 + 559524248170904730920655318548480a^2c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-8736796136721904754238181933056a^3c^8 - 1237734142501349223302162522112a^4c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{25253673488097375487494881280a^5c^8 + 1455699176221267243359055872a^6c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-12031819362886782611902464a^7c^8 - 901850710109690716936704a^8c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-3798874296621284474880a^9c^8 + 194341463383802701824a^{10}c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{2337313594001498112a^{11}c^8 - 2395393722129408a^{12}c^8 - 161395768934400a^{13}c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-718165493760a^{14}c^8 + 26357760a^{15}c^8 + 3294720a^{16}c^8}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
\end{aligned}$$

$$\begin{aligned}
& + \frac{1592821042644515671417734442254336c^9 - 460335074446405836514058089529344ac^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{36708731576238908672976168681472a^2c^9 - 9269192159255026038805102592a^3c^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-77931011136370050568382054400a^4c^9 + 355242949210634290070487040a^5c^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{73346374205830358870228992a^6c^9 + 251928402592276494729216a^7c^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-29002037259752505794560a^8c^9 - 323984159153576796160a^9c^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{2604142834119450624a^{10}c^9 + 56843610573381632a^{11}c^9 + 183910257090560a^{12}c^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-1148723896320a^{13}c^9 - 6788587520a^{14}c^9 - 5857280a^{15}c^9}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{133944629219042741828987058651136c^{10} - 31792962283341780433921881669632ac^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{1884977432671390301494262628352a^2c^{10} + 27453830985840912961790017536a^3c^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-3577899019963626250693312512a^4c^{10} - 25831060180904558769373184a^5c^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{2500762955655770765639680a^6c^{10} + 29770829896804480360448a^7c^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-554861886520804335616a^8c^{10} - 10477914723705741312a^9c^{10} - 6504383856402432a^{10}c^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{731660679766016a^{11}c^{10} + 3798320734208a^{12}c^{10} - 139403264a^{13}c^{10} - 19914752a^{14}c^{10}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{9259084618069199752092016705536c^{11} - 1788433975536396105112398856192ac^{11}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{74879477966758136016883679232a^2c^{11} + 2270097366668305229482033152a^3c^{11}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-119779384831109120949092352a^4c^{11} - 2108720637666398843895808a^5c^{11}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{56211218206919898071040a^6c^{11} + 1154277387520437207040a^7c^{11}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-4271295545929531392a^8c^{11} - 191483317319024640a^9c^{11} - 815139396943872a^{10}c^{11}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
\end{aligned}$$

$$\begin{aligned}
& + \frac{4260696014848a^{11}c^{11} + 29350723584a^{12}c^{11} + 25346048a^{13}c^{11}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{529170073110991196497521934336c^{12} - 82199817672273746020257497088ac^{12}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{2226310998038921142800482304a^2c^{12} + 112335291284151811064856576a^3c^{12}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-2816859314269848893194240a^4c^{12} - 83269328357107000442880a^5c^{12}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{705630734345855811584a^6c^{12} + 26456679475060703232a^7c^{12} + 65094234212646912a^8c^{12}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-1995437837352960a^9c^{12} - 12430943895552a^{10}c^{12} + 456228864a^{11}c^{12} + 76038144a^{12}c^{12}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{25084567677926533835412996096c^{13} - 3087201226200895282242125824ac^{13}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{45090915657763002741948416a^2c^{13} + 3939189564189899494522880a^3c^{13}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-40575227582415871672320a^4c^{13} - 2139706573184788070400a^5c^{13}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-195459046294880256a^6c^{13} + 382648988297920512a^7c^{13} + 2188013035192320a^8c^{13}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-9832316928000a^9c^{13} - 81208737792a^{10}c^{13} - 70189056a^{11}c^{13}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{987125304445771508818640896c^{14} - 94450060802172951369613312ac^{14}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{392250491442329132466176a^2c^{14} + 102263791723363946987520a^3c^{14}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-108880942638335262720a^4c^{14} - 37648129311289835520a^5c^{14}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-187827301329862656a^6c^{14} + 3333038389985280a^7c^{14} + 25954876784640a^8c^{14}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-952565760a^9c^{14} - 190513152a^{10}c^{14} + 32197592876889535514935296c^{15}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{-2337876825862944317243392ac^{15} - 10571276177233084940288a^2c^{15}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
& + \frac{1978662384171409211392a^3c^{15} + 10187017606434127872a^4c^{15} - 447086626545205248a^5c^{15}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} +
\end{aligned}$$

$$\begin{aligned}
 & + \frac{-3611715821371392a^6c^{15} + 14233364594688a^7c^{15} + 146822135808a^8c^{15} + 127008768a^9c^{15}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{866877134589278443012096c^{16} - 46282113328159709462528ac^{16}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-539135018545206001664a^2c^{16} + 28101817396392886272a^3c^{16}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{266925771696439296a^4c^{16} - 3333037754941440a^5c^{16} - 34606714060800a^6c^{16}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{1270087680a^7c^{16} + 317521920a^8c^{16} + 19126537123824754753536c^{17}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-719582772331401969664ac^{17} - 12566499273518088192a^2c^{17} + 281603400092614656a^3c^{17}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{3578843352268800a^4c^{17} - 12558776401920a^5c^{17} - 172582502400a^6c^{17} - 149422080a^7c^{17}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{341997755091266830336c^{18} - 8538595036697198592ac^{18} - 189974093577060352a^2c^{18}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{1829902732492800a^3c^{18} + 28499821199360a^4c^{18} - 1045954560a^5c^{18} - 348651520a^6c^{18}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{4874760435045236736c^{19} - 73736935746568192ac^{19} - 1950607118172160a^2c^{19}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{6169186795520a^3c^{19} + 127055953920a^4c^{19} + 110100480a^5c^{19} + 54050689474625536c^{20}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{-423766867836928ac^{20} - 13199997927424a^2c^{20} + 484442112a^3c^{20} + 242221056a^4c^{20}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{449048494473216c^{21} - 1292583829504ac^{21} - 53196357632a^2c^{21} - 46137344a^3c^{21}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{2628587094016c^{22} - 96468992ac^{22} - 96468992a^2c^{22} + 9663676416c^{23}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} + \\
 & + \frac{8388608ac^{23} + 16777216c^{24}}{\Gamma(\frac{c-a}{2}) \Gamma(\frac{c+a+49}{2})} \Big] \tag{8}
 \end{aligned}$$

Derivation of result (8):

Substituting $b = -a - 48, z = \frac{1}{2}$ in given result (2), we get

$$(2a + 48) {}_2F_1 \left[\begin{matrix} a, & -a - 48 & ; & 1 \\ & c & & 2 \end{matrix} \right]$$

$$= a {}_2F_1 \left[\begin{matrix} a+1, & -a-48 \\ c & \end{matrix} ; \frac{1}{2} \right] + (a+48) {}_2F_1 \left[\begin{matrix} a, & -a-47 \\ c & \end{matrix} ; \frac{1}{2} \right]$$

Now applying same parallel method which is used in Ref[6], we can prove the main formula.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

Orbit - Orbit Resonance of Pluto and Neptune

By M. A. Sharaf & L. A. Alaqal

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Abstract - In the present paper, an algorithm for the planar restricted circular threebody problem in rotating symbolic system is developed to determine orbit-orbit resonance of Pluto and Neptune.

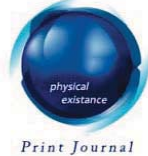
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GJSFR-F Classification : *MSC 2010: 70M20*



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Orbit - Orbit Resonance of Pluto and Neptune

M. A. Sharaf^α & L. A. Alaql^σ

Abstract - In the present paper, an algorithm for the planar restricted circular threebody problem in rotating symbolic system is developed to determine orbit-orbit resonance of Pluto and Neptune.

Keywords : Celestial Mechanics-Planetary close encounters-Pluto- Neptune- solar system dynamics

I. INTRODUCTION

Because of Pluto is locked into a 3:2 resonance with Neptune, Pluto completes 2 orbits every 3 orbits of the Sun completed by Neptune. Although, the ratio is not exactly 3:2, sometimes Pluto's period is slightly faster than its average value or slower. It is the possibility to indefinitely close approach if the perihelion and nod of Pluto are unrestricted, where, the radius of perihelion of Pluto is less than the radius of Neptune's orbit. This is the most important case of orbit - orbit resonance in the solar system. It is well known that the orbit of Pluto has a large eccentricity of $e = 0.247$, which brings the planet at a certain moment inside the orbit of Neptune. The two planets are trapped in an orbit-orbit resonance. The period of Pluto is 3:2 times the period of Neptune.

II. EQUATIONS OF MOTION

We consider the planar restricted circular three-body problem in rotating synodic system (e.g .Szebehely 1967) in which the two primaries are the Sun and Neptune while the third body is Pluto. The equations of motion to be solved are

$$\dot{x} = u, \quad (1.1)$$

$$\dot{y} = v, \quad (1.2)$$

$$\ddot{u} = -(1-\mu) \frac{x-\mu}{R_1^3} - \mu \frac{x+1-\mu}{R_2^3} + x + 2v, \quad (1.3)$$

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$$\dot{v} = -(1-\mu)\frac{y}{R_1^3} - \mu\frac{y}{R_2^3} + y + 2u, \quad (1.4)$$

with

$$R_1 = \sqrt{(x-\mu)^2 + y^2} \quad (2.1)$$

and

$$R_2 = \sqrt{(x+1-\mu)^2 + y^2}, \quad (2.2)$$

where dot denotes differentiation with respect to the time t , (x, y) are the coordinates of the third body, μ denotes the mass of the smaller primary when the total mass of the primaries has been normalized to unity.

In these equations, the unit of length is the distance between the primaries, the unit of mass is the sum of the masses of the primaries. The unit of time is $1/n$. (n is the mean motion). Normally, n is expressed in a number of radians per second, hence in $1/\text{sec}$. Its inverse $1/n$ is therefore expressed in seconds and may be interpreted as a unit of time.

III. ORBIT DETERMINATION OF PLUTO AND NEPTUNE

All the numerical values of the following are taken from the reference (Hellings, 1994)

a) *Orbital elements*

Neptune:

$$\begin{aligned} e &= 0 \\ P &= 165.62 \\ a &= 30.1584 \end{aligned}$$

Pluto:

$$\begin{aligned} e &= 0.247 \\ P &= 248.43 \\ a &= 39.5187 \end{aligned}$$

and $\mu = 0.0000525$

Amplitude of the libration is known to be 38° .

b) *Initial conditions*

The Initial conditions are

$$x_0 = -0.6073955952; y_0 = -0.7774968265; u_0 = 0.1083342234; v_0 = -0.0863997159.$$

c) *The results*

It should be noted that ,all the computations are performed using *Mathematica 7*. For clear illustrations of our analysis, the results are displayed graphically in the following manner

Fig.1 : Basic part of the orbit of Pluto relative to the Sun and Neptune

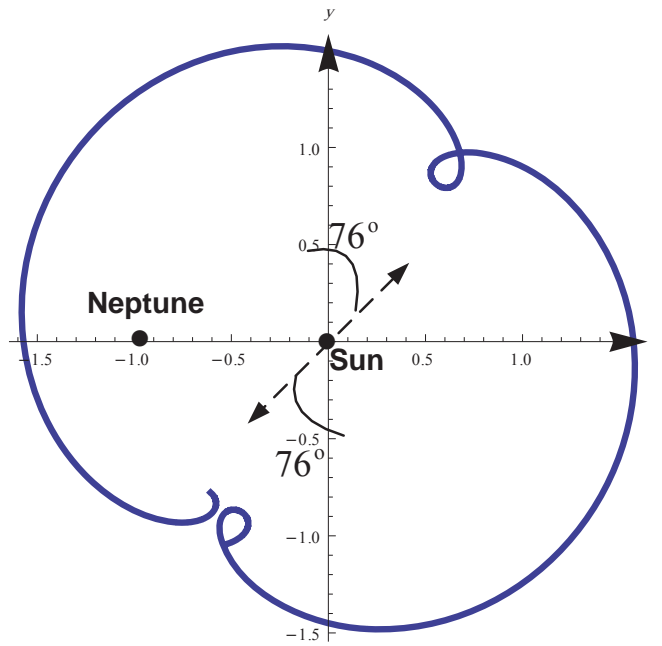


Fig.2 : Plot between u and v for the basic part of the orbit of Pluto relative to the Sun and Neptune

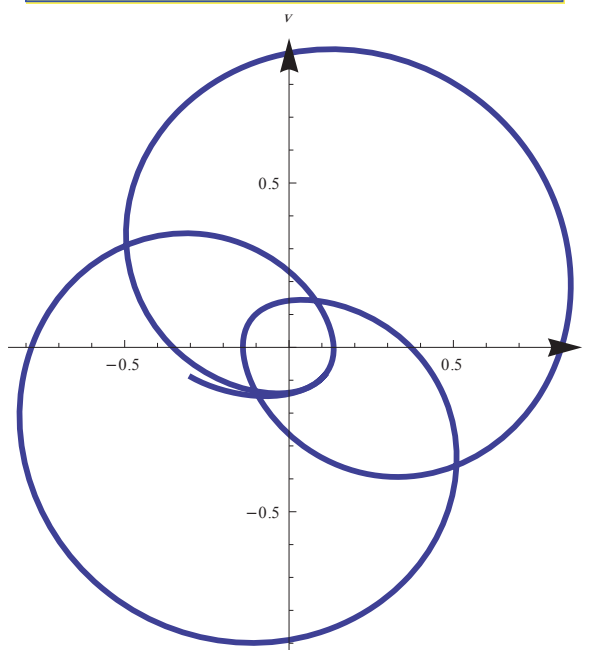


Fig.3 : Plot between x and u
for the basic part of the orbit of Pluto relative
to the Sun and Neptune

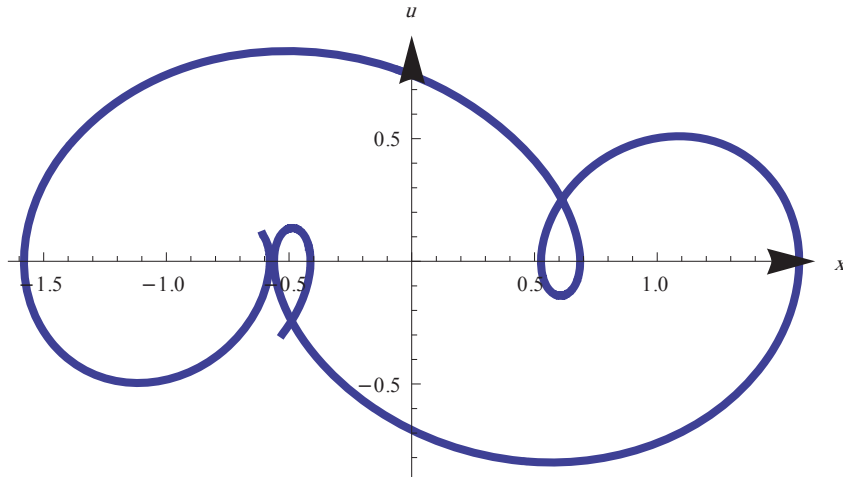


Fig.4 : Plot between y and v
for the basic part of the orbit of Pluto relative
to the Sun and Neptune

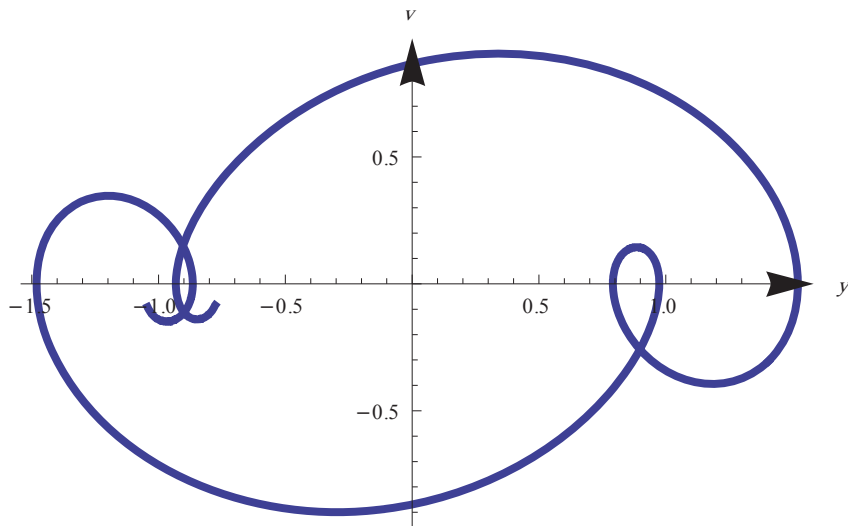


Fig.5 : The basic orbit of Fig.1 oscillates with an amplitude of 38°

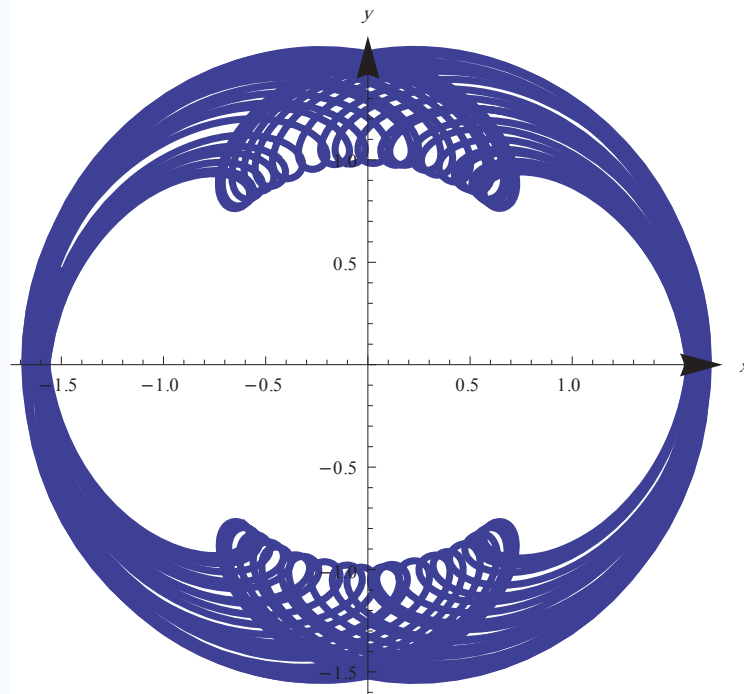


Fig.6 : The basic orbit of Fig.2 oscillates with an amplitude of 38°

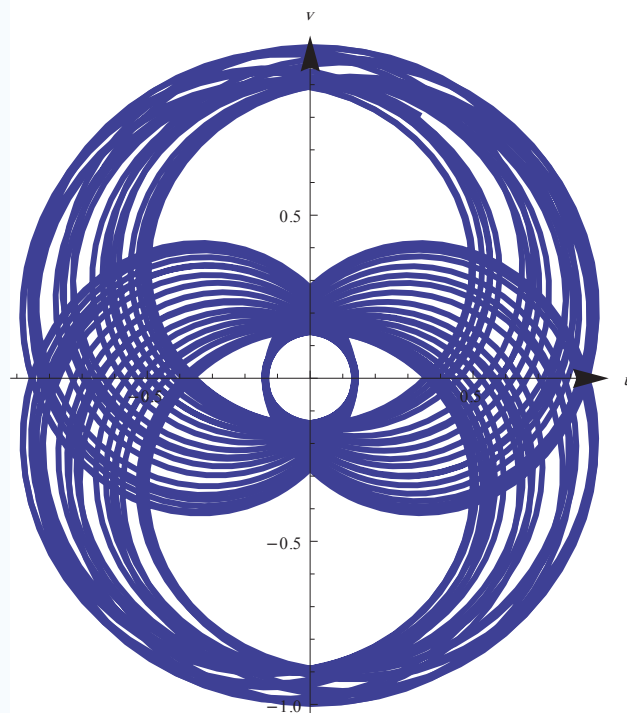


Fig.7 : The basic orbit of Fig.3 oscillates with an amplitude of 38°

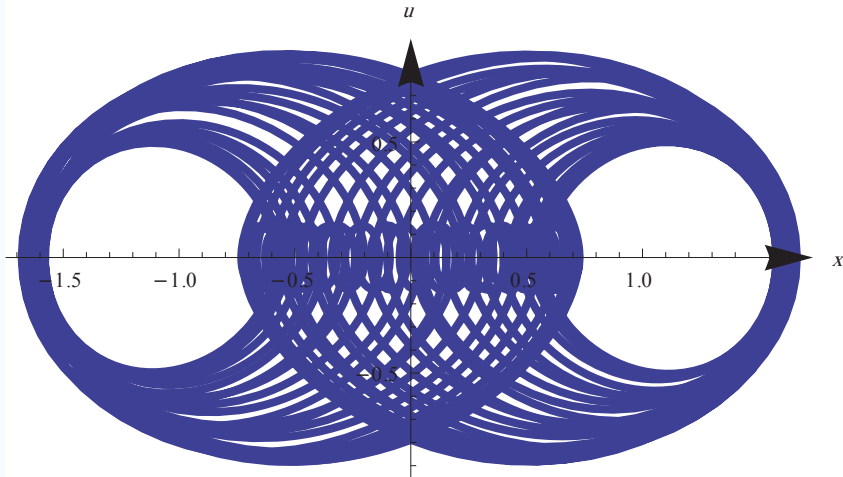
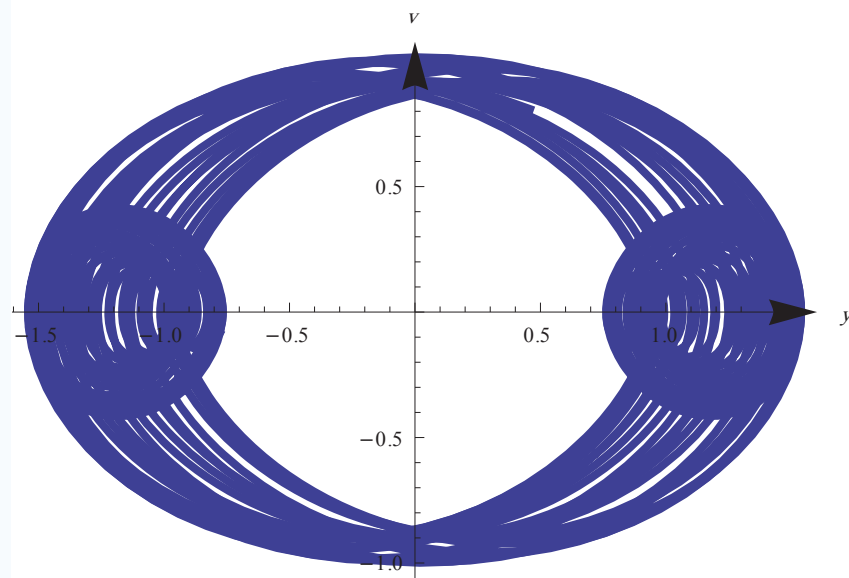


Fig.8 : The basic orbit of Fig.4 oscillates with an amplitude of 38°



IV. CONCLUSION

In this paper, general computational algorithm for the planar restricted circular three-body problem in rotating synodic system is developed in Section 2. This algorithm is applied to determinate orbit-orbit resonance of Pluto and Neptune. Finally the results are illustrate graphically in Section 3 which could be summarized as :

- 1-Figure 1 shows the basic part of the orbit of Pluto relative to the Sun and Neptune. This part represents two revolutions of Pluto around the Sun.
- 2- Pluto reaches two times a distance closer to the sun than Neptune.
- 3-The next two revolutions have the same shape, but the figure is rotated a little bit counter clockwise as the two perihelia approach the y- axis. This phenomenon increases until the whole figure is rotated over 76° .
- 4- A total libration, shown in Fig. 5, illustrates that Pluto will never collide with Neptune since its distance to Neptune is always larger than about 17 A.U.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

New Representations in Terms of q -product Identities for Ramanujan's Results IV

By M.P. Chaudhary, Upendra Kumar Pandit & Ashish Arora

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Abstract - In this paper author has established seven q -product identities, which are presumably new, and not available in the literature.

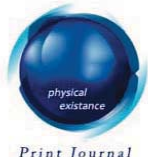
Keywords : *Theta functions, functions, triple product identities.*

GJSFR-F Classification : *AMS Subject Classifications: Primary 05A17, 05A15; Secondary 11P83*



Strictly as per the compliance and regulations of :





Ref.

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New Representations in Terms of q-product Identities for Ramanujan's Results IV

M.P. Chaudhary^α, Upendra Kumar Pandit^σ & Ashish Arora^ρ

Abstract - In this paper author has established seven q-product identities, which are presumably new, and not available in the literature.

Keywords : Theta functions, functions, triple product identities.

I. INTRODUCTION

For $|q| < 1$,

$$(a; q)_\infty = \prod_{n=0}^{\infty} (1 - aq^n) \tag{1.1}$$

$$(a; q)_\infty = \prod_{n=1}^{\infty} (1 - aq^{(n-1)}) \tag{1.2}$$

$$(a_1, a_2, a_3, \dots, a_k; q)_\infty = (a_1; q)_\infty (a_2; q)_\infty (a_3; q)_\infty \dots (a_k; q)_\infty \tag{1.3}$$

Ramanujan has defined general theta function, as

$$f(a, b) = \sum_{-\infty}^{\infty} a^{\frac{n(n+1)}{2}} b^{\frac{n(n-1)}{2}} ; |ab| < 1, \tag{1.4}$$

Jacobi's triple product identity [9,p.35] is given, as

$$f(a, b) = (-a; ab)_\infty (-b; ab)_\infty (ab; ab)_\infty \tag{1.5}$$

Special cases of Jacobi's triple products identity are given, as

$$\Phi(q) = \sum_{n=-\infty}^{\infty} q^{n^2} = (-q; q^2)_\infty^2 (q^2; q^2)_\infty \tag{1.6}$$

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$$\Psi(q) = \sum_{n=0}^{\infty} q^{\frac{n(n+1)}{2}} = \frac{(q^2; q^2)_{\infty}}{(q; q^2)_{\infty}} \tag{1.7}$$

$$f(-q) = \sum_{n=-\infty}^{\infty} (-1)^n q^{\frac{n(3n-1)}{2}} = (q; q)_{\infty} \tag{1.8}$$

Equation (1.8) is known as Euler’s pentagonal number theorem. Euler’s another well known identity is as

$$(q; q^2)_{\infty}^{-1} = (-q; q)_{\infty} \tag{1.9}$$

Roger-Ramanujan identities [6, p.578] are given as

$$G(q) = \sum_{n=0}^{\infty} \frac{q^{n^2}}{(q; q)_n} = \frac{1}{(q; q^5)_{\infty} (q^4; q^5)_{\infty}} = \frac{(q^2; q^5)_{\infty} (q^3; q^5)_{\infty} (q^5; q^5)_{\infty}}{(q; q)_{\infty}} \tag{1.10}$$

$$H(q) = \sum_{n=0}^{\infty} \frac{q^{n(n+1)}}{(q; q)_n} = \frac{1}{(q^2; q^5)_{\infty} (q^3; q^5)_{\infty}} = \frac{(q; q^5)_{\infty} (q^4; q^5)_{\infty} (q^5; q^5)_{\infty}}{(q; q)_{\infty}} \tag{1.11}$$

Roger-Ramanujan function is given by

$$R(q) = q^{\frac{1}{5}} \frac{H(q)}{G(q)} = q^{\frac{1}{5}} \frac{(q; q^5)_{\infty} (q^4; q^5)_{\infty}}{(q^2; q^5)_{\infty} (q^3; q^5)_{\infty}} \tag{1.12}$$

Throughout this paper we use the following representations

$$(q^a; q^n)_{\infty} (q^b; q^n)_{\infty} (q^c; q^n)_{\infty} \cdots (q^t; q^n)_{\infty} = (q^a, q^b, q^c \cdots q^t; q^n)_{\infty} \tag{1.13}$$

$$(q^a; q^n)_{\infty} (q^a; q^n)_{\infty} (q^c; q^n)_{\infty} \cdots (q^t; q^n)_{\infty} = (q^a, q^a, q^c \cdots q^t; q^n)_{\infty} \tag{1.14}$$

Now we can have following q-products identities, as

$$\begin{aligned} (q^2; q^2)_{\infty} &= \prod_{n=0}^{\infty} (1 - q^{2n+2}) \\ &= \prod_{n=0}^{\infty} (1 - q^{2(4n)+2}) \times \prod_{n=0}^{\infty} (1 - q^{2(4n+1)+2}) \times \prod_{n=0}^{\infty} (1 - q^{2(4n+2)+2}) \times \prod_{n=0}^{\infty} (1 - q^{2(4n+3)+2}) \\ &= \prod_{n=0}^{\infty} (1 - q^{8n+2}) \times \prod_{n=0}^{\infty} (1 - q^{8n+4}) \times \prod_{n=0}^{\infty} (1 - q^{8n+6}) \times \prod_{n=0}^{\infty} (1 - q^{8n+8}) \\ &= (q^2; q^8)_{\infty} (q^4; q^8)_{\infty} (q^6; q^8)_{\infty} (q^8; q^8)_{\infty} = (q^2, q^4, q^6, q^8; q^8)_{\infty} \end{aligned} \tag{1.15}$$

Ref.

6. G.E. Andrews, R. Askey and R. Roy; *Special Functions*, Cambridge University Press, Cambridge, 1999.

$$\begin{aligned}
(q^4; q^4)_\infty &= \prod_{n=0}^{\infty} (1 - q^{4n+4}) \\
&= \prod_{n=0}^{\infty} (1 - q^{4(3n)+4}) \times \prod_{n=0}^{\infty} (1 - q^{4(3n+1)+4}) \times \prod_{n=0}^{\infty} (1 - q^{4(3n+2)+4}) \\
&= \prod_{n=0}^{\infty} (1 - q^{12n+4}) \times \prod_{n=0}^{\infty} (1 - q^{12n+8}) \times \prod_{n=0}^{\infty} (1 - q^{12n+12}) \\
&= (q^4; q^{12})_\infty (q^8; q^{12})_\infty (q^{12}; q^{12})_\infty = (q^4, q^8, q^{12}; q^{12})_\infty \tag{1.16}
\end{aligned}$$

$$\begin{aligned}
(q^4; q^{12})_\infty &= \prod_{n=0}^{\infty} (1 - q^{12n+4}) = \prod_{n=0}^{\infty} (1 - q^{12(5n)+4}) \times \prod_{n=0}^{\infty} (1 - q^{12(5n+1)+4}) \times \\
&\quad \times \prod_{n=0}^{\infty} (1 - q^{12(5n+2)+4}) \times \prod_{n=0}^{\infty} (1 - q^{12(5n+3)+4}) \times \prod_{n=0}^{\infty} (1 - q^{12(5n+4)+4}) \\
&= \prod_{n=0}^{\infty} (1 - q^{60n+4}) \times \prod_{n=0}^{\infty} (1 - q^{60n+16}) \times \prod_{n=0}^{\infty} (1 - q^{60n+28}) \times \prod_{n=0}^{\infty} (1 - q^{60n+40}) \times \prod_{n=0}^{\infty} (1 - q^{60n+52}) \\
&= (q^4; q^{60})_\infty (q^{16}; q^{60})_\infty (q^{28}; q^{60})_\infty (q^{40}; q^{60})_\infty (q^{52}; q^{60})_\infty = (q^4, q^{16}, q^{28}, q^{40}, q^{52}; q^{60})_\infty \tag{1.17}
\end{aligned}$$

Similarly we can compute following as

$$\begin{aligned}
(q^4; q^{12})_\infty &= (q^4; q^{60})_\infty (q^{16}; q^{60})_\infty (q^{28}; q^{60})_\infty (q^{40}; q^{60})_\infty (q^{52}; q^{60})_\infty \\
&= (q^4, q^{16}, q^{28}, q^{40}, q^{52}; q^{60})_\infty \tag{1.18}
\end{aligned}$$

$$(q^6; q^6)_\infty = (q^6; q^{24})_\infty (q^{12}; q^{24})_\infty (q^{18}; q^{24})_\infty (q^{24}; q^{24})_\infty = (q^6, q^{12}, q^{18}, q^{24}; q^{24})_\infty \tag{1.19}$$

$$\begin{aligned}
(q^6; q^{12})_\infty &= (q^6; q^{60})_\infty (q^{18}; q^{60})_\infty (q^{30}; q^{60})_\infty (q^{42}; q^{60})_\infty (q^{54}; q^{60})_\infty \\
&= (q^6, q^{18}, q^{30}, q^{42}, q^{54}; q^{60})_\infty \tag{1.20}
\end{aligned}$$

$$\begin{aligned}
(q^8; q^8)_\infty &= (q^8; q^{48})_\infty (q^{16}; q^{48})_\infty (q^{24}; q^{48})_\infty (q^{32}; q^{48})_\infty (q^{40}; q^{48})_\infty (q^{48}; q^{48})_\infty \\
&= (q^8, q^{16}, q^{24}, q^{32}, q^{40}, q^{48}; q^{48})_\infty \tag{1.21}
\end{aligned}$$

$$\begin{aligned}
(q^8; q^{12})_\infty &= (q^8; q^{60})_\infty (q^{20}; q^{60})_\infty (q^{32}; q^{60})_\infty (q^{44}; q^{60})_\infty (q^{56}; q^{60})_\infty \\
&= (q^8, q^{20}, q^{32}, q^{44}, q^{56}; q^{60})_\infty \tag{1.22}
\end{aligned}$$

$$(q^8; q^{16})_\infty = (q^8; q^{48})_\infty (q^{24}; q^{48})_\infty (q^{40}; q^{48})_\infty = (q^8, q^{24}, q^{40}; q^{48})_\infty \tag{1.23}$$

$$(q^{10}; q^{20})_\infty = (q^{10}; q^{60})_\infty (q^{30}; q^{60})_\infty (q^{50}; q^{60})_\infty = (q^{10}, q^{30}, q^{50}; q^{60})_\infty \tag{1.24}$$

$$\begin{aligned}
(q^{12}; q^{12})_\infty &= (q^{12}; q^{60})_\infty (q^{24}; q^{60})_\infty (q^{36}; q^{60})_\infty (q^{48}; q^{60})_\infty (q^{60}; q^{60})_\infty \\
&= (q^{12}, q^{24}, q^{36}, q^{48}, q^{60}; q^{60})_\infty \tag{1.25}
\end{aligned}$$

$$(q^{16}; q^{16})_{\infty} = (q^{16}; q^{48})_{\infty} (q^{32}; q^{48})_{\infty} (q^{48}; q^{48})_{\infty} = (q^{16}, q^{32}, q^{48}; q^{48})_{\infty} \quad (1.26)$$

$$(q^{20}; q^{20})_{\infty} = (q^{20}; q^{60})_{\infty} (q^{40}; q^{60})_{\infty} (q^{60}; q^{60})_{\infty} = (q^{20}, q^{40}, q^{60}; q^{60})_{\infty} \quad (1.27)$$

The outline of this paper is as follows. In sections 2, we have recorded some well known results, those are useful to the rest of the paper. In section 3, we state and prove seven new q-product identities, which are not available in the literature of special functions.

II. PRELIMINARIES

Let us recall the definition of cubic theta functions $A(q), B(q)$ and $C(q)$ due to Borwein et al.[4], as

$$A(q) = \sum_{m,n=-\infty}^{\infty} q^{m^2+mn+n^2} \quad (2.1)$$

$$B(q) = \sum_{m,n=-\infty}^{\infty} \omega^{m-n} q^{m^2+mn+n^2}; \quad \omega = \exp\left(\frac{2\pi i}{3}\right) \quad (2.2)$$

$$C(q) = \sum_{m,n=-\infty}^{\infty} q^{m^2+mn+n^2+m+n} \quad (2.3)$$

Borwein et al.[4] established the following relations

$$A(q) = A(q^3) + 2qC(q^3) \quad (2.4)$$

$$B(q) = A(q^3) - qC(q^3) \quad (2.5)$$

$$C(q) = \frac{3(q^3; q^3)_{\infty}^3}{(q; q)_{\infty}} \quad (2.6)$$

$$A(q)A(q^2) = B(q)B(q^2) + qC(q)C(q^2) \quad (2.7)$$

Entry-2, in Ramanujan's first note book [8, p.230], [10, p.356] is stated as

$$\Psi(q)\Psi(q^3) - \Psi(-q)\Psi(-q^3) = 2q\Phi(q^2)\Psi(q^{12}) \quad (2.8)$$

Entry-4(iv), in the chapter 20 of Ramanujan's second note book [8], [9, p.359] is stated as

$$\Phi(q)\Phi(q^{27}) - \Phi(-q)\Phi(-q^{27}) = 4qf(-q^6)f(-q^{18}) + 4q^7\Psi(q^2)\Psi(q^{54}) \quad (2.9)$$

Entry-9(i), in the chapter 20 of Ramanujan's second note book [8], [9, p.277] is stated as

$$\Psi(q^3)\Psi(q^5) - \Psi(-q^3)\Psi(-q^5) = 2q^3\Psi(q^2)\Psi(q^{30}) \quad (2.10)$$

Ref.

4. J.M. Borwein, P.B. Borwein and F.G. Garvan; *Some cubic modular identities of Ramanujan*, trans.Amer.Math.Soc., 343(1994), 35-47.

Entry-9(iii), in the chapter 20 of Ramanujan’s second note book [8], [9, p.377] is stated as

$$\Phi(q^3)\Phi(q^5) = \Phi(-q^2)\Phi(-Q^2) + 2q^2\Psi(q)\Psi(Q); \text{ where } Q = q^{15} \tag{2.11}$$

Entry-9(iv), in the chapter 20 of Ramanujan’s second note book [8], [9, p.377] is stated as

$$\Psi(q)\Psi(q^{15}) + \Psi(-q)\Psi(-q^{15}) = 2\Psi(q^6)\Psi(q^{10}) \tag{2.12}$$

Entry-25, in Ramanujan’s note book [9, p.39] is stated as

$$\Phi(q) + \Phi(-q) = 2\Phi(q^4) \tag{2.13}$$

$$\Phi(q) - \Phi(-q) = 4q\Psi(q^8) \tag{2.14}$$

$$\Phi(q)\Phi(-q) = \Phi(-q^2) \tag{2.15}$$

III. MAIN RESULTS

We have establish following

$$(q^2, q^4, q^6; q^8)_\infty [(-q; q^2)_\infty^2 + (q; q^2)_\infty^2] = 2(-q^4; q^8)_\infty^2 \tag{3.1}$$

$$(q^2, q^4, q^6, q^8; q^8)_\infty [(-q; q^2)_\infty^2 - (q; q^2)_\infty^2] = 4q \frac{(q^{16}, q^{32}, q^{48}; q^{48})_\infty}{(q^8, q^{24}, q^{40}; q^{48})_\infty} \tag{3.2}$$

$$\frac{(-q; q^2)_\infty^2 + (q; q^2)_\infty^2}{(-q; q^2)_\infty^2 - (q; q^2)_\infty^2} = \frac{(-q^4; q^8)_\infty^2 (q^8, q^8, q^{24}, q^{24}, q^{40}, q^{40}; q^{48})_\infty}{2q} \tag{3.3}$$

$$(-q; q^2)_\infty^2 (q; q^2)_\infty^2 (q^2; q^2)_\infty^2 = (q^2, q^2, q^4; q^4)_\infty \tag{3.4}$$

$$\frac{(-q; q^2)_\infty (-q^3; q^6)_\infty - (q; q^2)_\infty (q^3; q^6)_\infty}{(-q; q^2)_\infty \times (-q^3; q^6)_\infty \times (q; q^2)_\infty \times (q^3; q^6)_\infty} = \frac{2q(-q^2; q^4)_\infty^2 (q^4, q^8, q^{16}, q^{20}, q^{24}; q^{24})_\infty}{(q^2, q^4, q^6, q^8; q^8)_\infty (q^6, q^{12}, q^{18}; q^{24})_\infty} \tag{3.5}$$

$$\begin{aligned} & \frac{(-q^3; q^6)_\infty (-q^5; q^{10})_\infty - (q^3; q^6)_\infty (q^5; q^{10})_\infty}{(-q^3; q^6)_\infty \times (-q^5; q^{10})_\infty \times (q^3; q^6)_\infty \times (q^5; q^{10})_\infty} = \frac{(q^4, q^8, q^{12}; q^{12})_\infty}{(q^6, q^{12}, q^{18}, q^{24}; q^{24})_\infty} \times \\ & \times \frac{2q^3}{(q^2, q^6, q^{10}; q^{12})_\infty (q^{10}, q^{20}, q^{30}, q^{30}, q^{40}, q^{50}; q^{60})_\infty} \end{aligned} \tag{3.6}$$

$$\begin{aligned} & \frac{[(q; q^2)_\infty (q^{15}; q^{30})_\infty] + [(-q; q^2)_\infty (-q^{15}; q^{30})_\infty]}{[(q; q^2)_\infty (q^{15}; q^{30})_\infty][(-q; q^2)_\infty (-q^{15}; q^{30})_\infty]} = \frac{(q^{12}, q^{20}, q^{24}, q^{36}, q^{40}, q^{48}, q^{60}, q^{60}; q^{60})_\infty}{(q^{10}, q^{30}, q^{30}, q^{50}, q^{60}; q^{60})_\infty} \times \\ & \times \frac{2}{(q^2, q^4, q^6, q^8, q^8; q^8)_\infty (q^6, q^{18}, q^{30}, q^{42}, q^{54}; q^{60})_\infty} \end{aligned} \tag{3.7}$$

Proof of (3.1):Employing equation (1.6) in equation (2.13), we have

$$(-q; q^2)_\infty^2 (q^2; q^2)_\infty + (q; q^2)_\infty^2 (q^2; q^2)_\infty = 2(-q^4; q^8)_\infty^2 (q^8; q^8)_\infty$$

$$\begin{aligned}
& (q^2; q^2)_\infty [(-q; q^2)_\infty^2 + (q; q^2)_\infty^2] = 2(-q^4; q^8)_\infty^2 (q^8; q^8)_\infty \\
& (q^2; q^8)_\infty (q^4; q^8)_\infty (q^6; q^8)_\infty (q^8; q^8)_\infty [(-q; q^2)_\infty^2 + (q; q^2)_\infty^2] = 2(-q^4; q^8)_\infty^2 (q^8; q^8)_\infty \\
& (q^2; q^8)_\infty (q^4; q^8)_\infty (q^6; q^8)_\infty [(-q; q^2)_\infty^2 + (q; q^2)_\infty^2] = 2(-q^4; q^8)_\infty^2 \\
& (q^2, q^4, q^6; q^8)_\infty [(-q; q^2)_\infty^2 + (q; q^2)_\infty^2] = 2(-q^4; q^8)_\infty^2
\end{aligned}$$

which establish the result (3.1).

Proof of (3.2):Employing equations (1.6) and (1.7) in equation (2.14), we have

$$(-q; q^2)_\infty^2 (q^2; q^2)_\infty - (q; q^2)_\infty^2 (q^2; q^2)_\infty = \frac{4q(q^{16}; q^{16})_\infty}{(q^8; q^{16})_\infty}$$

$$(q^2; q^2)_\infty [(-q; q^2)_\infty^2 - (q; q^2)_\infty^2] = 4q \frac{(q^{16}, q^{32}, q^{48}; q^{48})_\infty}{(q^8, q^{24}, q^{40}; q^{48})_\infty}$$

$$(q^2, q^4, q^6, q^8; q^8)_\infty [(-q; q^2)_\infty^2 - (q; q^2)_\infty^2] = 4q \frac{(q^{16}, q^{32}, q^{48}; q^{48})_\infty}{(q^8, q^{24}, q^{40}; q^{48})_\infty}$$

which establish the result (3.2).

Proof of (3.3):Dividing equation (3.1) by (3.2), we get equation (3.3).

Proof of (3.4):Employing equation (1.6) in equation (2.15), we have

$$(-q; q^2)_\infty^2 (q^2; q^2)_\infty (q; q^2)_\infty^2 (q^2; q^2)_\infty = (q^2; q^4)_\infty^2 (q^4; q^4)_\infty$$

$$(-q; q^2)_\infty^2 (q; q^2)_\infty^2 (q^2; q^2)_\infty^2 = (q^2; q^4)_\infty^2 (q^4; q^4)_\infty$$

$$(-q; q^2)_\infty^2 (q; q^2)_\infty^2 (q^2; q^2)_\infty^2 = (q^2; q^4)_\infty (q^2; q^4)_\infty (q^4; q^4)_\infty$$

$$(-q; q^2)_\infty^2 (q; q^2)_\infty^2 (q^2; q^2)_\infty^2 = (q^2, q^2, q^4; q^4)_\infty$$

which establish the result (3.4).

Proof of (3.5):Employing equations (1.6) and (1.7) in equation (2.8), we get.

$$\frac{(q^2; q^2)_\infty (q^6; q^6)_\infty}{(q; q^2)_\infty (q^3; q^6)_\infty} - \frac{(q^2; q^2)_\infty (q^6; q^6)_\infty}{(-q; q^2)_\infty (-q^3; q^6)_\infty} = \frac{2q(-q^2; q^4)_\infty^2 (q^4; q^4)_\infty (q^{24}; q^{24})_\infty}{(q^{12}; q^{24})_\infty}$$

$$\frac{(-q; q^2)_\infty (-q^3; q^6)_\infty - (q; q^2)_\infty (q^3; q^6)_\infty}{(-q; q^2)_\infty (-q^3; q^6)_\infty (q; q^2)_\infty (q^3; q^6)_\infty} = \frac{2q(-q^2; q^4)_\infty^2 (q^4; q^4)_\infty (q^{24}; q^{24})_\infty}{(q^{12}; q^{24})_\infty (q^2; q^2)_\infty (q^6; q^6)_\infty}$$

$$\frac{(-q; q^2)_\infty (-q^3; q^6)_\infty - (q; q^2)_\infty (q^3; q^6)_\infty}{(-q; q^2)_\infty (-q^3; q^6)_\infty (q; q^2)_\infty (q^3; q^6)_\infty} = \frac{2q(-q^2; q^4)_\infty^2 (q^4, q^8, q^{16}, q^{20}, q^{24}; q^{24})_\infty}{(q^2; q^2)_\infty (q^6, q^{12}, q^{18}; q^{24})_\infty}$$

$$\frac{(-q; q^2)_\infty (-q^3; q^6)_\infty - (q; q^2)_\infty (q^3; q^6)_\infty}{(-q; q^2)_\infty (-q^3; q^6)_\infty (q; q^2)_\infty (q^3; q^6)_\infty} = \frac{2q(-q^2; q^4)_\infty^2 (q^4, q^8, q^{16}, q^{20}, q^{24}; q^{24})_\infty}{(q^2, q^4, q^6, q^8; q^8)_\infty (q^6, q^{12}, q^{18}; q^{24})_\infty}$$

which establish the result (3.5).

Proof of (3.6):Employing equation (1.7) in equation (2.10), we get.

$$\frac{(q^6; q^6)_\infty (q^{10}; q^{10})_\infty}{(q^3; q^6)_\infty (q^5; q^{10})_\infty} - \frac{(q^6; q^6)_\infty (q^{10}; q^{10})_\infty}{(-q^3; q^6)_\infty (-q^5; q^{10})_\infty} = \frac{2q^3(q^4; q^4)_\infty (q^{60}; q^{60})_\infty}{(q^2; q^4)_\infty (q^{30}; q^{60})_\infty}$$

$$\frac{(-q^3; q^6)_\infty (-q^5; q^{10})_\infty - (q^3; q^6)_\infty (q^5; q^{10})_\infty}{(-q^3; q^6)_\infty (-q^5; q^{10})_\infty (q^3; q^6)_\infty (q^5; q^{10})_\infty} = \frac{2q^3 (q^4; q^4)_\infty (q^{60}; q^{60})_\infty}{(q^2; q^4)_\infty (q^6; q^6)_\infty (q^{10}; q^{10})_\infty (q^{30}; q^{60})_\infty}$$

$$\frac{(-q^3; q^6)_\infty (-q^5; q^{10})_\infty - (q^3; q^6)_\infty (q^5; q^{10})_\infty}{(-q^3; q^6)_\infty (-q^5; q^{10})_\infty (q^3; q^6)_\infty (q^5; q^{10})_\infty} = \frac{(q^4, q^8, q^{12}; q^{12})_\infty}{(q^6, q^{12}, q^{18}, q^{24}; q^{24})_\infty} \times$$

$$\times \frac{2q^3}{(q^2, q^6, q^{10}; q^{12})_\infty (q^{10}, q^{20}, q^{30}, q^{30}, q^{40}, q^{50}; q^{60})_\infty}$$

which establish the result (3.6).

Proof of (3.7): Employing equation (1.7) in equation (2.12), we get.

$$\frac{(q^2; q^2)_\infty (q^{30}; q^{30})_\infty}{(q; q^2)_\infty (q^{15}; q^{30})_\infty} + \frac{(q^2; q^2)_\infty (q^{30}; q^{30})_\infty}{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty} = \frac{2(q^{12}; q^{12})_\infty (q^{20}; q^{20})_\infty}{(q^6; q^{12})_\infty (q^{10}; q^{20})_\infty}$$

$$\frac{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty + (q; q^2)_\infty (q^{15}; q^{30})_\infty}{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty (q; q^2)_\infty (q^{15}; q^{30})_\infty} = \frac{2(q^{12}; q^{12})_\infty (q^{20}; q^{20})_\infty}{(q^2; q^2)_\infty (q^6; q^{12})_\infty (q^{10}; q^{20})_\infty (q^{30}; q^{30})_\infty}$$

$$\frac{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty + (q; q^2)_\infty (q^{15}; q^{30})_\infty}{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty (q; q^2)_\infty (q^{15}; q^{30})_\infty} = \frac{2(q^{12}, q^{20}, q^{24}, q^{36}, q^{40}, q^{48}, q^{60}, q^{60}; q^{60})_\infty}{(q^2; q^2)_\infty (q^6; q^{12})_\infty (q^{10}; q^{20})_\infty (q^{30}; q^{30})_\infty}$$

$$\frac{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty + (q; q^2)_\infty (q^{15}; q^{30})_\infty}{(-q; q^2)_\infty (-q^{15}; q^{30})_\infty (q; q^2)_\infty (q^{15}; q^{30})_\infty} = \frac{(q^{12}, q^{20}, q^{24}, q^{36}, q^{40}, q^{48}, q^{60}, q^{60}; q^{60})_\infty}{(q^{10}, q^{30}, q^{30}, q^{50}, q^{60}; q^{60})_\infty} \times$$

$$\times \frac{2}{(q^2, q^4, q^6, q^8, q^8; q^8)_\infty (q^6, q^{18}, q^{30}, q^{42}, q^{54}; q^{60})_\infty}$$

which establish the result (3.7).

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Note Uncertain Field of Fractions

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Abstract - The set of some real rhotrices of the same dimension D^* was defined in [2] to be an integral domain. An example of a finite field $M[R_3]$ was given in [4] based on this definition also and on the construction of finite fields presented in [3]. It was discovered that the finite sub collection of the elements of $M[R_3]$ as contained in D^* is not closed under rhotrix addition and hence not an integral domain. More generally, D^* is not an integral domain as it is not closed under rhotrix addition. This problem affects the field of fractions constructed in [8]. A solution to this problem is provided in this article and the construction method of such fields is reviewed. This reviewed version gives the generalization of such construction as the n-dimensional rhotrices are used.

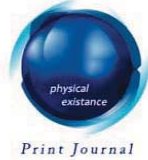
Keywords : *n-dimensional rhotrix; Quotient rhotrix; Integral domain; Field of fraction.*

GJSFR-F Classification : *MSC 2010: 83A05*



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Note Uncertain Field of Fractions

S. Usaini^α & S. M. Tudunkaya^σ

Abstract -The set of some real rhotrices of the same dimension D^* was defined in [2] to be an integral domain. An example of a finite field $M[R_3]$ was given in [4] based on this definition also and on the construction of finite fields presented in [3]. It was discovered that the finite sub collection of the elements of $M[R_3]$ as contained in D^* is not closed under rhotrix addition and hence not an integral domain. More generally, D^* is not an integral domain as it is not closed under rhotrix addition. This problem affects the field of fractions constructed in [8].A solution to this problem is provided in this article and the construction method of such fields is reviewed. This reviewed version gives the generalization of such construction as the n-dimensional rhotrices are used.

Keywords : n-dimensional rhotrix; Quotient rhotrix; Integral domain; Field of fraction

1. INTRODUCTION

The idea of classifying the set of all rhotrices of dimension 3 as abstract structures was presented in [1] and [2].In [2] the set of some 3-dimensional real rhotrices

$$D^* = \langle (R - ZD), +, \circ \rangle \tag{1.1}$$

was defined to be an integral domain under rhotrix addition and multiplication, where R is the set of all real rhotrices of dimension 3 as defined in [6] by

$$R = \left\{ \left\langle \begin{matrix} a \\ b & c & d \\ e \end{matrix} \right\rangle : a, b, c, d, e \in \mathfrak{R} \right\},$$

$$ZD = \left\{ \left\langle \begin{matrix} a \\ b & 0 & d \\ e \end{matrix} \right\rangle : a, b, d, e, 0 \in \mathfrak{R} \text{ and at least one of } a, b, d, e \neq 0 \right\}.$$

Recall that an integral domain is a commutative ring with out zero divisors. However, D^* is not even a ring because the additive closure is not there. This can be seen as follows: Let $R, Q \in D^*$ such that $h(R) = c$ and $h(Q) = -c$. If $R + Q = S$ then $h(S) = 0$ and at least one of $a, b, d, e \neq 0$ which implies that $S \notin D^*$.

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II. A PARTICULAR FIELD OF FRACTION OF AN INTEGRAL DOMAIN

Theorem 2.1

Let H_n be the set of all n-dimensional integer hearty rhotrices. If $H_n^* = H_n \cup \{O_n\}$ then H_n^* is an integral domain, where O_n is an n-dimensional zero rhotrix.

Proof

It suffices to show that H_n^* is isomorphic to an integral domain \mathbb{Z} . That is $H_n^* \cong \mathbb{Z}$.

Define a mapping $\tau : \mathbb{Z} \rightarrow H_n^*$ by $\tau(c) = C_n$.

For homomorphism, let $c, d \in \mathbb{Z}$, then

$$(i) \tau(c + d) = C_n + D_n = \tau(c) + \tau(d) \quad (ii) \tau(cd) = C_n \circ D_n = \tau(c)\tau(d)$$

Therefore τ is a homomorphism

Since $\forall \tau(c) \in H_n^*$ there exists $c \in \mathbb{Z}$ such that $\tau(c) = C_n$ then τ is onto.

Now let $\tau(c), \tau(d) \in H_n^*$ such that $\tau(c) = \tau(d)$.

$$\tau(c) = \tau(d) \Rightarrow C_n = D_n \Rightarrow c = d.$$

Thus τ is one to one.

Hence $H_n^* \cong \mathbb{Z}$.

Definition 2.1

Let H_n^* and H_n be as in Theorem 2.1 above. Then a relation \sim on $H_n^* \times H_n$ defined by cross multiplication as $(C_{n1}, D_{n1}) \sim (C_{n2}, D_{n2})$ if $C_{n1} \circ D_{n2} = C_{n2} \circ D_{n1}$, $\exists C_{n1}, C_{n2} \in H_n^*$; $D_{n1}, D_{n2} \in H_n$.

Proposition 2.1

The relation \sim as defined in Definition (2.1) is an equivalence relation.

Proof

Reflexivity and Symmetry of the relation are obvious.

For transitivity, let $C_{n1} \circ D_{n2} = C_{n2} \circ D_{n1}$ and $C_{n2} \circ D_{n3} = C_{n3} \circ D_{n2}$. To show that $C_{n1} \circ D_{n3} = C_{n3} \circ D_{n1}$ we have

$$\begin{aligned} (C_{n1} \circ D_{n3}) \circ D_{n2} &= (C_{n1} \circ D_{n2}) \circ D_{n3} \\ &= (C_{n2} \circ D_{n1}) \circ D_{n3} \\ &= (C_{n2} \circ D_{n3}) \circ D_{n1} \\ &= (C_{n3} \circ D_{n2}) \circ D_{n1} \\ &= (C_{n3} \circ D_{n1}) \circ D_{n2} \end{aligned}$$

$\therefore (C_{n1} \circ D_{n3}) = (C_{n3} \circ D_{n1})$ by cancellation law.

We denote by $\frac{C_n}{D_n}$ the equivalence class of (C_n, D_n) in $H_n^* \times H_n$ and define $H_n^*[H_n^{-1}]$ to

be the set of all the equivalence classes $\frac{C_n}{D_n}$, where $C_n \in H_n^*$ and $D_n \in H_n$.

For all $\frac{C_{n1}}{D_{n1}}, \frac{C_{n2}}{D_{n2}} \in H_n^*[H_n^{-1}]$ we define addition and multiplication on $H_n^*[H_n^{-1}]$ as

$$\text{follows: } \frac{C_{n1}}{D_{n1}} +' \frac{C_{n2}}{D_{n2}} = \frac{C_{n1} \circ D_{n2} + C_{n2} \circ D_{n1}}{D_{n1} \circ D_{n2}} \text{ and } \frac{C_{n1}}{D_{n1}} \circ' \frac{C_{n2}}{D_{n2}} = \frac{C_{n1} \circ C_{n2}}{D_{n1} \circ D_{n2}}.$$

Proposition 2.2

The operations $(+')$, (\circ') as defined above are well-defined.

Proof

Suppose $\frac{C'_{n1}}{D'_{n1}} = \frac{C_{n1}}{D_{n1}}$ and $\frac{C'_{n2}}{D'_{n2}} = \frac{C_{n2}}{D_{n2}}$; then $C'_{n1} \circ D_{n1} = D'_{n1} \circ C_{n1}$ and $C'_{n2} \circ D_{n2} = D'_{n2} \circ C_{n2}$, so

$$\text{that } (C'_{n1} \circ D'_{n2} + C'_{n2} \circ D'_{n1})D_{n1}D_{n2} = C'_{n1} \circ D_{n1} \circ D'_{n2} \circ D_{n2} + C'_{n2} \circ D_{n2} \circ D'_{n1} \circ D_{n1}$$

$$= C_{n1} \circ D'_{n1} \circ D'_{n2} \circ D_{n2} + C_{n2} \circ D'_{n2} \circ D'_{n1} \circ D_{n1}$$

$$= (C_{n1} \circ D_{n2} + C_{n2} \circ D_{n1})D'_{n1} \circ D'_{n2}$$

$$\text{implying that } \frac{C'_{n1}}{D'_{n1}} +' \frac{C'_{n2}}{D'_{n2}} = \frac{C_{n1}}{D_{n1}} +' \frac{C_{n2}}{D_{n2}}.$$

$$\text{Similarly } (C'_{n1} \circ C'_{n2})D_{n1} \circ D_{n2} = (C_{n1} \circ C_{n2})D'_{n1} \circ D'_{n2} \text{ implies that } \frac{C'_{n1}}{D'_{n1}} \circ' \frac{C'_{n2}}{D'_{n2}} = \frac{C_{n1} \circ C_{n2}}{D_{n1} \circ D_{n2}}.$$

By definition (1.3) the equivalence class $\frac{C_n}{D_n} = C_n \circ D_n^{-1}$ since $D_n \neq 0_n \in H_n$. Therefore, for

all $D_n \in H_n$, $0_n \in H_n^*$, $\frac{0_n}{D_n} = 0_n \circ D_n^{-1} = 0_n = 0_n \circ I_n = \frac{0_n}{I_n}$. Thus $\frac{0_n}{I_n} = \frac{0_n}{D_n}$ is the additive

identity and $-\frac{C_n}{D_n} = \frac{-C_n}{D_n}$ is the additive inverse. Similarly, $\frac{I_n}{I_n} = \frac{D_n}{D_n}$ is the multiplicative identity.

Theorem 2.2

With the above definitions and the definitions of the operations $(+')$ and (\circ') , the set of the equivalence classes $H_n^*[H_n^{-1}]$ is a commutative ring.

Proof

One should check that the properties of a ring are fulfilled. But the proof follows from the fact that addition and multiplication are the regular addition and multiplication of fractions.

Proposition 2.3

The function $\psi : H_n^* \rightarrow H_n^*[H_n^{-1}]$ defined by $\psi(C_n) = \frac{C_n}{I_n}$ is a ring homomorphism whose kernel is $\{C_n \in H_n^* : C_n \circ' D_n = 0 \text{ for some } D_n \in H_n^*[H_n^{-1}]\}$.

Proof

Let $C_{n1}, C_{n2} \in H_n^*$, then

$$\psi(C_{n1} + C_{n2}) = \frac{C_{n1} + C_{n2}}{I_n} = (C_{n1} + C_{n2}) \circ I_n = C_{n1} \circ I_n + C_{n2} \circ I_n = \frac{C_{n1}}{I_n} + \frac{C_{n2}}{I_n} = \psi(C_{n1}) + \psi(C_{n2})$$

$$\psi(C_{n1} \circ' C_{n2}) = \frac{C_{n1} \circ' C_{n2}}{I_n} = (C_{n1} \circ' C_{n2}) \circ I_n = C_{n1} \circ I_n \circ' C_{n2} \circ I_n = \frac{C_{n1}}{I_n} \circ' \frac{C_{n2}}{I_n} = \psi(C_{n1}) \circ' \psi(C_{n2})$$

$$\psi(I_n) = \frac{I_n}{I_n}.$$

Now $C_n \in \ker \psi$ if and only if $\frac{C_n}{I_n} = \frac{0_n}{I_n}$, if and only if $C_n \circ I_n = 0_n \circ I_n = 0_n$, which imply that $\ker \psi = \{0_n\}$.

Recall from [1] that, the set $M = \{nI : n \in \mathbb{Z}\}$ where I is the unity element of the commutative ring of 3-dimensional rhotrices R is a subring and submonoid of R under multiplication (\circ). Thus the set $M_n = \{nI_n : n \in \mathbb{Z}\}$ is a subring and submonoid of the commutative ring R_n^* of n-dimensional rhotrices. Therefore any submonoid, H_n of R_n^* with property that for all $Q_n \neq 0 \in R_n^*$ and $S_n \in H_n$, $Q_n \circ S_n \neq 0$ can serve in the above construction for the generalization of proposition 2.3 as stated in the following proposition.

Proposition 2.4

$R_n^*[H_n^{-1}]$ as constructed above is a ring, and there is a homomorphism $\psi : R_n^* \rightarrow R_n^*[H_n^{-1}]$ given by $\psi(Q_n) = \frac{Q_n}{I_n}$.

Proof

The proof follows from propositions 2.2 and 2.3.

As defined in [1], a diagonal rhotrix of dimension 3 is a rhotrix whose two non-diagonal entries are all zero. Let D be the set of all n-dimensional diagonal rhotrices then it is easy for someone to verify that $W = D \cup \{O_n, I_n\}$, where O_n is the n-dimensional additive identity; I_n is the n-dimensional multiplicative identity is a group and is normal in R_n^* under multiplication.

Proposition 2.5

Let R_n^* be a commutative ring of n-dimensional rhotrices, and let H_n be a submonoid of R_n^* such that $Q_n \circ S_n \neq 0_n$ for every $Q_n \neq 0_n \in R_n^*$ and $S_n \in H_n$. Then every ideal of $R_n^*[H_n^{-1}]$ has the form $W[H_n^{-1}]$, for suitable W normal in R_n^* .

Proof

Since all the elements of $W[H_n^{-1}]$ are also elements of $R_n^*[H_n^{-1}]$ and $I_n \in H_n$ then obviously $W[H_n^{-1}]$ is an additive subgroup of $R_n^*[H_n^{-1}]$.

Ref.

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For all $\frac{C_n}{D_n} \in R_n^*[H_n^{-1}] ; \frac{W_n}{D'_n} \in W[H_n^{-1}] ; \frac{C_n}{D_n} \circ' \frac{W_n}{D'_n} = \frac{C_n \circ W_n}{D_n \circ D'_n} \in W[H_n^{-1}]$ since

$$C_n \circ W_n \in W, D_n \circ D'_n \in H_n.$$

W is normal in R_n^* implies that $\frac{W_n}{D'_n} \circ \frac{C_n}{D_n} \in W[H_n^{-1}]$.

Proposition 2.6

$H_n^*[H_n^{-1}]$ is an integral domain.

Proof

Suppose $\frac{C_{n1}}{D_{n1}} \circ' \frac{C_{n2}}{D_{n2}} = 0_n \in H_n^*[H_n^{-1}]$, that is $\frac{C_{n1} \circ C_{n2}}{D_{n1} \circ D_{n2}} = \frac{0_n}{I_n}$

$\Rightarrow (C_{n1} \circ C_{n2}, D_{n1} \circ D_{n2}) \sim (0_n, I_n)$ and $C_{n1} \circ C_{n2} \circ D_n = 0$ for some $D_n \in H_n$.

$C_{n1} \circ C_{n2} \circ D_n = 0_n \in H_n^*$, which is an integral domain, and $D_n \neq 0_n$, thus $C_{n1} \circ C_{n2} = 0_n$.

So either C_{n1} or C_{n2} is 0_n and consequently either $\frac{C_{n1}}{D_{n1}}$ or $\frac{C_{n2}}{D_{n2}}$ is 0_n .

Theorem 2.3

The set $H_n^*[H_n^{-1}]$ of all equivalence classes $\frac{C_n}{D_n}$ is a field.

Proof

From Theorem 2.2, $H_n^*[H_n^{-1}]$ is a commutative ring with unity $\frac{I_n}{I_n}$. So we just need to show that every non zero element of $H_n^*[H_n^{-1}]$ has multiplicative inverse.

Suppose $\frac{C_n}{D_n} \neq \frac{0_n}{I_n}$, then $C_n \neq 0_n$, so $C_n \in H_n$ which implies that $\frac{C_n}{D_n} \in H_n^*[H_n^{-1}]$.

Clearly, $\frac{C_n}{D_n} \circ' \frac{D_n}{C_n} = \frac{C_n \circ D_n}{D_n \circ C_n} = \frac{I_n}{I_n}$. Thus $\frac{D_n}{C_n}$ is the multiplicative inverse of $\frac{C_n}{D_n}$.

III. CONCLUSION

In this short note, amendment concerning some definitions in [2] and [8] with their generalizations were provided. The steps observed in the construction of field of fractions illustrated in [8] were also amended respectively.

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GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH
MATHEMATICS AND DECISION SCIENCES
Volume 12 Issue 12 Version 1.0 Year 2012
Type : Double Blind Peer Reviewed International Research Journal
Publisher: Global Journals Inc. (USA)
Online ISSN: 2249-4626 & Print ISSN: 0975-5896

Bianchi Type- VI_0 Dark Energy Cosmological Models in General Relativity

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Abstract - Bianchi type- VI_0 cosmological models of the universe filled with dark energy with constant and time-dependent equation of state parameters are investigated in general relativity. We obtain exact solutions of Einstein's field equations using the condition that the shear scalar is proportional to the expansion scalar, which represent singular and non-singular cosmological models of the universe. The physical behavior of the models are discussed. We conclude that the universe models do not approach isotropy through the evolution of the universe.

Keywords : *Bianchi type-VI. Dark energy. Cosmological models.*

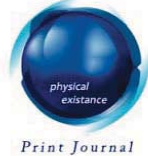
GJSFR-F Classification : *MSC 2010: 83A05*



BIANCHI TYPE-VI₀ DARK ENERGY COSMOLOGICAL MODELS IN GENERAL RELATIVITY

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I. INTRODUCTION

Recent observations on expansion history of the universe indicate that the universe is currently experiencing a phase of accelerated expansion. This was first observed from high red shift supernova Ia (Reiss et al. [1-2], Perlmutter et al. [3], Astier et al. [4], Spergel et al.[5] etc.) and confirmed later by cross checks from the cosmic microwave background radiation (Bennett et al. [6], Abazajian et al.[7-9], Hawkins et al. [10] etc.). The current accelerating expansion of the universe attributed to the fact that our universe is dominated by an unknown dark energy DE an exotic energy with negative pressure.

The simplest dark energy candidate is the vacuum energy density which is mathematically equivalent to the cosmological constant Λ . As per Copeland et al. [11] "fine tuning" and the cosmic "coincidence" are the two well known difficulties of the cosmological constant problems. There are several alternative theories for the dynamical DE scenario which have been proposed by scientists to interpret the accelerating universe. Wang and Tegmark [12] have shown that the universe is actually undergoing an acceleration with repulsive gravity of some strange energy-form i.e. DE at work. Dark energy is a mysterious substance with negative pressure and accounts for nearly 70% of total matter-energy of universe, but has no clear explanation. Karami et al. [13] introduced a polytropic gas model of DE as an alternative model to explain the accelerated expansion of the universe. Gupta and Pradhan [14] proposed a new candidate known as cosmological nuclear-energy as a possible candidate for the dark energy.

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3. Perlmutter, S., et al.: Astrophys. J. **483**, 565 (1997)

Bianchi types I-IX cosmological models are important in the sense that these are homogeneous and anisotropic, from which the process of isotropization of the universe is studied through the passage of time. Moreover, from the theoretical point of view anisotropic universe have a greater generality than FRW isotropic models. The simplicity of the field equations made Bianchi space-times useful in constructing models of spatially homogeneous and anisotropic cosmologies. Considerable works have been done in obtaining various Bianchi type cosmological models and their inhomogeneous generalization. Bianchi type-VI₀ space-time is of special interest in anisotropic cosmology. Barrow [15] pointed out that Bianchi type-VI₀ models of the universe give a better explanation of some of the cosmological problems like primordial helium abundance and they also isotropize in a special sense. Looking to the importance of Bianchi type-VI₀ universes, many authors [16-20] have studied it in different context. Shri Ram[21, 22] has presented Bianchi type-VI₀ cosmological models filled with dust and perfect fluid in modified Brans-Dicke theory respectively.

Adhav et al. [23] studied Bianchi type-VI₀ cosmological models with anisotropic dark energy. Abdussattar and Prajapati [24] obtained a class of bouncing non-singular FRW models by constraining the deceleration parameter (DP) in the presence of an interacting dark energy represented by a time-varying cosmological constant. They have also discussed the role of deceleration parameter and interacting dark energy in singularity avoidance. Bisabr [25] has shown that an accelerating expansion is possible in a spatially flat universe for large values of the Brans-Dicke parameter consistent with the local gravity experiments. Yadav and Saha [26] studied DE models with variable equation of state (EoS) parameter. Recently, Saha and Yadav [27] presented a general relativistic cosmological model with time-dependent DP in LRS Bianchitype-II space-time which can be described by isotopic and variable EoS parameter. In this paper, We present general relativistic cosmological models with constant and time-dependent DP in Bianchi type-VI₀ space-time which can be described by isotropic constant and variable EoS parameters. This paper is organized as follows: We present the metric and field equations in Sect.2. In Sect.3, we obtain the solutions of the field equations representing Bianchi type-VI₀ cosmological models with perfect fluid by imposing the condition that the shear scalar is proportional to expansion scalar. We also discuss the physical behaviors of the cosmological models with dark energy. Concluding remarks are given in Sect.4.

II. THE METRIC AND FIELD EQUATIONS

We consider the spatially homogeneous and anisotropic Bianchi type-VI₀ space-time in the form

$$ds^2 = -dt^2 + A^2(t)dx^2 + B^2(t)e^{-2mx}dy^2 + C^2(t)e^{2mx}dz^2 \quad (1)$$

where A , B and C are functions of the cosmic time t and m is a constant

The Einstein's field equations, in natural limits ($8\pi G = c = 1$) are

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = -T_{\mu\nu} \quad (2)$$

Ref.

15. Barrow, J.D.: Mon. Not. Astron. Soc. **211**, 221(1984)

where $R_{\mu\nu}$ is the Ricci tensor, R is the Ricci scalar curvature and $T_{\mu\nu}$ is the energy-momentum tensor of matter. For a perfect fluid distribution, the tensor $T_{\mu\nu}$ is given by

$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu + pg_{\mu\nu} \quad (3)$$

where ρ is the energy density of the cosmic matter p is the isotropic pressure and u^μ is the four-velocity vector. In comoving coordinate system $u^\mu = (0, 0, 0, 1)$, the Einstein's field equation (2) together with (3), for the metric (1), yield

$$\frac{\ddot{B}}{B} + \frac{\ddot{C}}{C} + \frac{\dot{B}\dot{C}}{BC} + \frac{m^2}{A^2} = -\omega\rho, \quad (4)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{C}}{C} + \frac{\dot{A}\dot{C}}{AC} - \frac{m^2}{A^2} = -\omega\rho, \quad (5)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} + \frac{\dot{A}\dot{B}}{AB} - \frac{m^2}{A^2} = -\omega\rho, \quad (6)$$

$$\frac{\dot{A}\dot{B}}{AB} + \frac{\dot{A}\dot{C}}{AC} + \frac{\dot{B}\dot{C}}{BC} - \frac{m^2}{A^2} = \rho, \quad (7)$$

$$\left(\frac{\dot{B}}{B} - \frac{\dot{C}}{C}\right) = 0 \quad (8)$$

where ω is the EoS parameter given by

$$p = \omega\rho \quad (9)$$

and a dot denotes ordinary differentiation with respect to t .

The average scalar factor a and volume scalar V are given by

$$a^3 = V = ABC. \quad (10)$$

The generalized mean Hubble parameter H is defined by

$$H = \frac{\dot{a}}{a} = \frac{1}{3}(H_1 + H_2 + H_3) \quad (11)$$

where the directional Hubble parameters H_1 , H_2 and H_3 are given by

$$H_1 = \frac{\dot{A}}{A}, H_2 = \frac{\dot{B}}{B}, H_3 = \frac{\dot{C}}{C}. \quad (12)$$

The expansion scalar θ and shear scalar σ are given by

$$\theta = \frac{\dot{A}}{A} + \frac{\dot{B}}{B} + \frac{\dot{C}}{C}, \quad (13)$$

$$\sigma^2 = \frac{1}{2}\left[\frac{\dot{A}^2}{A^2} + \frac{\dot{B}^2}{B^2} + \frac{\dot{C}^2}{C^2}\right] - \frac{1}{6}\theta^2. \quad (14)$$

The deceleration parameter q is defined by

$$q = -1 + \frac{d}{dt}(H). \quad (15)$$

The sign of q indicates whether the model inflates or not. A positive sign of q corresponds to the standard decelerating model whereas the negative sign of q indicates inflation. The recent observations of SN Ia (Reiss et al.[1], Perlmutter et al.[3]) reveal that the present universe is accelerating and the value of DP lies somewhere in the range $-1 < q < 0$.

III. SOLUTION OF FIELD EQUATIONS

Equation (8), on integration, gives

$$B = C \quad (16)$$

where the constant of integration is absorbed in B or C . Using (16), equations (4) – (7) reduce to

$$\frac{2\ddot{B}}{B} + \frac{\dot{B}^2}{B^2} + \frac{m^2}{A^2} = -\omega\rho, \quad (17)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} + \frac{\dot{A}\dot{B}}{AB} - \frac{m^2}{A^2} = -\omega\rho, \quad (18)$$

$$\frac{2\dot{A}\dot{B}}{AB} + \frac{\dot{B}^2}{B^2} - \frac{m^2}{A^2} = \rho. \quad (19)$$

These are three equations connecting four unknown functions A , B , ρ and ω . In order to solve the above equations we use the physical condition that expansions scalar is proportional to shear scalar, which in our case leads to

$$A = B^n \quad (20)$$

where n is a constant. Roy and Banerjee [28], Bali and Singh [29] have proposed this condition to find exact solutions of cosmological models.

Here we use the procedure of Saha and Yadav [27] to find exact solutions of (17) – (19) combining (10) and (20), we obtain

$$A = V^{\frac{n}{n+1}}, \quad B = V^{\frac{1}{n+1}}. \quad (21)$$

Subtraction of (18) from (17) gives

$$\frac{\ddot{B}}{B} - \frac{\ddot{A}}{A} + \frac{\dot{B}^2}{B^2} - \frac{\dot{A}\dot{B}}{AB} + \frac{2m^2}{A^2} = 0. \quad (22)$$

Substituting (21) into (22), we obtain

$$\ddot{V} = \frac{2m^2(n+2)}{n-1} V^{\frac{2-n}{n+2}}. \quad (23)$$

The first integral of (23) is

$$\int \frac{dV}{V^{\frac{4}{n+2}+C}} = \frac{m(n+2)t}{\sqrt{n-1}} \quad (24)$$

Ref.

3. Perlmutter, S., et al.: *Astrophys. J.* **483**, 565 (1997)

where C is an arbitrary constant. Clearly (24) imposes some restriction on the choice of n namely, $n > 1$. It is not possible to solve equation (24) in general. So, in order to solve the problem completely, we have to choose either C or n in such a way that (24) be integrable. Therefore we consider the following cases.

Case 3.1 When $C=0$

In this case the solution of (24) is

$$V = \left(\frac{mn}{\sqrt{n-1}} \right)^{\frac{n+2}{n}} (t + k_1) \quad (25)$$

where k_1 is an arbitrary constant. From (21) and (25) we obtain the scale factor as

$$A = \frac{mn}{\sqrt{n-1}}(t + k), \quad (26)$$

$$B = \left(\frac{mn}{\sqrt{n-1}} \right)^{\frac{1}{n}} (t + k)^{\frac{1}{n}}. \quad (27)$$

With these scale factors, the metric (1) can be written in form

$$ds^2 = -dT^2 + \left(\frac{mn}{\sqrt{n-1}} \right)^2 dx^2 + \left(\frac{mn}{\sqrt{n-1}} \right)^{\frac{2}{n}} T^{\frac{2}{n}} (e^{-2mx} dy^2 + e^{2mx} dz^2) \quad (28)$$

where $T=t+k$.

The expressions for the energy density ρ and the EOS ω for the model (28) are obtained as

$$\rho = \frac{1+n}{n^2 T^2}, \quad (29)$$

$$\omega = \frac{n-2}{n+1}. \quad (30)$$

The other physical and kinematical parameters are given by

$$nH_1 = H_2 = H_3 = \frac{1}{T}, \quad (31)$$

$$\theta = 3H = \frac{n+2}{nT}, \quad (32)$$

$$\sigma = \frac{1}{\sqrt{3}} \frac{n-1}{nT}, \quad (33)$$

$$q = -\frac{2}{n+2}. \quad (34)$$

The deceleration parameter q is always negative. The EoS parameter is positive when $n > 2$ and is negative if $1 < n < 2$. Thus, the metric (28) represents as ever power-law accelerated expansion universe filled with a perfect fluid. If $1 < n < 2$, $\omega < 0$, we obtain DE cosmological model of Bianchi type-VI₀.

The spatial volume V is zero and all physical parameters diverge at $T = 0$. Therefore, the model has a point-type singularity at $T = 0$. For $0 < T < \infty$, the spatial volume is an increasing function of time. The physical parameters are monotonically decreasing function of time and ultimately tend to zero for large T . The anisotropy in the model is maintained throughout the passage of time. For the physical reality of the model we will have to choose n , greater than 1, in such a way that $\left| \frac{n-2}{n+2} \right| \leq 1$. It deserves mention that we are unable to find n for which $\omega = \pm 1$.

Case 3.2 When $C \neq 0$

When $C \neq 0$ equation (24) is not integrable for general values of n . However, for $n = 2$, it becomes

$$\int \frac{dV}{\sqrt{V+C}} = 4mt \quad (35)$$

which, after integration, yields

$$V = 4m^2t^2 + 2\beta t + \gamma \quad (36)$$

where β and γ are arbitrary constants. The constant C is absorbed in γ . From (21) and (36), we obtain the scale factors as

$$A = (4m^2t^2 + 2\beta t + \gamma)^{\frac{1}{2}}, \quad (37)$$

$$B = (4m^2t^2 + 2\beta t + \gamma)^{\frac{1}{4}}. \quad (38)$$

Therefore, the metric (1) of our solutions can be written in the form

$$ds^2 = -dt^2 + (4m^2t^2 + 2\beta t + \gamma)dx^2 + (4m^2t^2 + 2\beta t + \gamma)^{\frac{1}{2}}(e^{-2mx} dy^2 + e^{2mx} dz^2) \quad (39)$$

The expressions for (H_1, H_2, H_3) , H , ρ , θ and σ are obtained as

$$H_1 = \frac{4m^2t + \beta}{4m^2t^2 + 2\beta t + \gamma}, \quad (40)$$

$$H_2 = H_3 = \frac{1}{2} \left(\frac{4m^2t + \beta}{4m^2t^2 + 2\beta t + \gamma} \right), \quad (41)$$

$$H = \frac{2}{3} \left(\frac{4m^2t + \beta}{4m^2t^2 + 2\beta t + \gamma} \right), \quad (42)$$

$$\theta = 2 \left(\frac{4m^2t + \beta}{4m^2t^2 + 2\beta t + \gamma} \right), \quad (43)$$

$$\sigma = \frac{1}{\sqrt{3}} \left(\frac{4m^2t + \beta}{4m^2t^2 + 2\beta t + \gamma} \right). \quad (44)$$

The energy density, DP and ω are obtained as

$$\rho = \frac{(8m^2t + 2\beta)^2 - (4m^2\gamma - \beta^2)}{4(4m^2t^2 + 2\beta t + \gamma)^2}, \quad (45)$$

$$\omega = -\frac{5(4m^2\gamma - \beta^2)}{(8m^2t + 2\beta)^2 - (4m^2\gamma - \beta^2)}, \quad (46)$$

$$q = -\frac{2m^2(4m^2t^2 + 2\beta t + \gamma)}{(4m^2t + \beta)^2}. \quad (47)$$

The value of DP is always negative since V is never negative. The EoS parameter ω is negative if $\gamma > \frac{\beta^2}{4m^2}$. If this condition holds, the model (39) corresponds to a Bianchi type-VI₀ energy cosmological model with variable q and ω .

If $\gamma > \frac{\beta^2}{4m^2}$, the model (39) has no finite singularity. The physical and kinematical parameters are all decreasing function of time and ultimately tend to zero for large time. The model essentially gives an empty space-time for large time. The anisotropy in the model never dies out.

IV. CONCLUSION

In this paper, we have presented exact solutions of Einstein's field equations for a Bianchi-type VI₀ space-time filled with perfect fluid satisfying the barotropic equation of state under the assumption that the expansion scalar is proportional to shear scalar. Under some specific choice of problem parameters, the present consideration yields singular and non-singular models of the accelerated expansion universe filled with perfect fluid and dark energy. Models with negative EoS parameter ω may be attributed to the current accelerated expansion of universe. The physical and kinematical parameters are all decreasing function of time and ultimately tend to zero for large time. The universe models do not approach to isotropy. The models presented in this paper can be potential tools to describe the present universe as well as the early universe.

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- Consequences, including definite statistics - if the consequences are quantitative in nature, account quantitative data; results of any numerical analysis should be reported
- Significant conclusions or questions that track from the research(es)

Approach:

- Single section, and succinct
- As a outline of job done, it is always written in past tense
- A conceptual should situate on its own, and not submit to any other part of the paper such as a form or table
- Center on shortening results - bound background information to a verdict or two, if completely necessary
- What you account in an conceptual must be regular with what you reported in the manuscript
- Exact spelling, clearness of sentences and phrases, and appropriate reporting of quantities (proper units, important statistics) are just as significant in an abstract as they are anywhere else

Introduction:

The **Introduction** should "introduce" the manuscript. The reviewer should be presented with sufficient background information to be capable to comprehend and calculate the purpose of your study without having to submit to other works. The basis for the study should be offered. Give most important references but shun difficult to make a comprehensive appraisal of the topic. In the introduction, describe the problem visibly. If the problem is not acknowledged in a logical, reasonable way, the reviewer will have no attention in your result. Speak in common terms about techniques used to explain the problem, if needed, but do not present any particulars about the protocols here. Following approach can create a valuable beginning:

- Explain the value (significance) of the study
- Shield the model - why did you employ this particular system or method? What is its compensation? You strength remark on its appropriateness from a abstract point of vision as well as point out sensible reasons for using it.
- Present a justification. Status your particular theory (es) or aim(s), and describe the logic that led you to choose them.
- Very for a short time explain the tentative propose and how it skilled the declared objectives.

Approach:

- Use past tense except for when referring to recognized facts. After all, the manuscript will be submitted after the entire job is done.
- Sort out your thoughts; manufacture one key point with every section. If you make the four points listed above, you will need a least of four paragraphs.
- Present surroundings information only as desirable in order hold up a situation. The reviewer does not desire to read the whole thing you know about a topic.
- Shape the theory/purpose specifically - do not take a broad view.
- As always, give awareness to spelling, simplicity and correctness of sentences and phrases.

Procedures (Methods and Materials):

This part is supposed to be the easiest to carve if you have good skills. A sound written Procedures segment allows a capable scientist to replacement your results. Present precise information about your supplies. The suppliers and clarity of reagents can be helpful bits of information. Present methods in sequential order but linked methodologies can be grouped as a segment. Be concise when relating the protocols. Attempt for the least amount of information that would permit another capable scientist to spare your outcome but be cautious that vital information is integrated. The use of subheadings is suggested and ought to be synchronized with the results section. When a technique is used that has been well described in another object, mention the specific item describing a way but draw the basic



principle while stating the situation. The purpose is to text all particular resources and broad procedures, so that another person may use some or all of the methods in one more study or referee the scientific value of your work. It is not to be a step by step report of the whole thing you did, nor is a methods section a set of orders.

Materials:

- Explain materials individually only if the study is so complex that it saves liberty this way.
- Embrace particular materials, and any tools or provisions that are not frequently found in laboratories.
- Do not take in frequently found.
- If use of a definite type of tools.
- Materials may be reported in a part section or else they may be recognized along with your measures.

Methods:

- Report the method (not particulars of each process that engaged the same methodology)
- Describe the method entirely
- To be succinct, present methods under headings dedicated to specific dealings or groups of measures
- Simplify - details how procedures were completed not how they were exclusively performed on a particular day.
- If well known procedures were used, account the procedure by name, possibly with reference, and that's all.

Approach:

- It is embarrassed or not possible to use vigorous voice when documenting methods with no using first person, which would focus the reviewer's interest on the researcher rather than the job. As a result when script up the methods most authors use third person passive voice.
- Use standard style in this and in every other part of the paper - avoid familiar lists, and use full sentences.

What to keep away from

- Resources and methods are not a set of information.
- Skip all descriptive information and surroundings - save it for the argument.
- Leave out information that is immaterial to a third party.

Results:

The principle of a results segment is to present and demonstrate your conclusion. Create this part a entirely objective details of the outcome, and save all understanding for the discussion.

The page length of this segment is set by the sum and types of data to be reported. Carry on to be to the point, by means of statistics and tables, if suitable, to present consequences most efficiently. You must obviously differentiate material that would usually be incorporated in a study editorial from any unprocessed data or additional appendix matter that would not be available. In fact, such matter should not be submitted at all except requested by the instructor.

Content

- Sum up your conclusion in text and demonstrate them, if suitable, with figures and tables.
- In manuscript, explain each of your consequences, point the reader to remarks that are most appropriate.
- Present a background, such as by describing the question that was addressed by creation an exacting study.
- Explain results of control experiments and comprise remarks that are not accessible in a prescribed figure or table, if appropriate.
- Examine your data, then prepare the analyzed (transformed) data in the form of a figure (graph), table, or in manuscript form.

What to stay away from

- Do not discuss or infer your outcome, report surroundings information, or try to explain anything.
- Not at all, take in raw data or intermediate calculations in a research manuscript.

- Do not present the similar data more than once.
- Manuscript should complement any figures or tables, not duplicate the identical information.
- Never confuse figures with tables - there is a difference.

Approach

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- Put figures and tables, appropriately numbered, in order at the end of the report
- If you desire, you may place your figures and tables properly within the text of your results part.

Figures and tables

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- Make a decision if each premise is supported, discarded, or if you cannot make a conclusion with assurance. Do not just dismiss a study or part of a study as "uncertain."
- Research papers are not acknowledged if the work is imperfect. Draw what conclusions you can based upon the results that you have, and take care of the study as a finished work
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- Make a decision if the tentative design sufficiently addressed the theory, and whether or not it was correctly restricted.
- Try to present substitute explanations if sensible alternatives be present.
- One research will not counter an overall question, so maintain the large picture in mind, where do you go next? The best studies unlock new avenues of study. What questions remain?
- Recommendations for detailed papers will offer supplementary suggestions.

Approach:

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