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Highlights

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# GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH: F MATHEMATICS & DECISION SCIENCES

# GLOBAL JOURNAL OF SCIENCE FRONTIER RESEARCH: F Mathematics & Decision Sciences

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# On Some Properties of the Rising Sun Function

# By Vajha Srinivasa Kumar

Abstract- This paper studies a few interesting properties of the rising sun function of a bounded real function defined on a closed and bounded interval on the real line. An operator on the space of all bounded real functions defined on a closed and bounded interval is introduced and its properties are investigated.

*Keywords:* rising sun function, semi-continuity, darboux continuity, lower (upper) semicontinuity, lower (upper) semi-quasicontinuity, symmetric continuity, cliquishness, quasicontinuity, differentiability.

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Riesz, F. And B.Sz. Nagy, Functional Analysis, Frederick Ungar Publishing Co.,

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# On Some Properties of the Rising Sun Function

Vajha Srinivasa Kumar

Abstract- This paper studies a few interesting properties of the rising sun function of a bounded real function defined on a closed and bounded interval on the real line. An operator on the space of all bounded real functions defined on a closed and bounded interval is introduced and its properties are investigated.

Keywords : rising sun function, semi-continuity, darboux continuity, lower (upper) semicontinuity, lower (upper) semi-quasicontinuity, symmetric continuity, cliquishness, quasicontinuity, differentiability.

#### I. INTRODUCTION

The rising sun function was used as a tool in the proof of the famous Lebesgue's theorem on the differentiability of a real valued monotone function without using the theory of integration [3]. In this paper, some properties of the rising sun function are presented and an operator on the space of all bounded real functions defined on a closed and bounded interval [a,b] is introduced and its properties are investigated.

In what follows X,  $\mathbb{R}$  and  $\mathbb{N}$  stand for a topological space, the real line and the set of all positive integers respectively. Also  $\mathcal{B}$  stands for the Banach space of all bounded real functions defined on a closed and bounded interval [a,b] where  $a,b \in \mathbb{R}$  and a < b under the supremum norm.

#### II. Preliminaries

**1.1 Definition [6]:** The rising sun function of a function  $f \in \mathcal{B}$  is defined by

$$f_{\odot}(x) = \sup \left\{ f(y) \, / \, x \le y \le b \right\}.$$

**1.2 Definition:** For  $f \in \mathcal{B}$  we define the following.

(*i*) 
$$_{\odot} f(x) = \sup \{ f(y) / a \le y \le x \}$$
  
(*ii*)  $f^{\odot}(x) = \inf \{ f(y) / x \le y \le b \}$   
(*iii*)  $^{\odot} f(x) = \inf \{ f(y) / a \le y \le x \}$ 

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- 1.3 Definition [5]: A function f: X → R is said to be semi-continuous at a point p∈ X if for every ε>0 and every neighborhood U of p in X there exists a non-empty open set W⊂U such that |f(x)-f(p)| < ε ∀ x∈ W. We say that a function f is semi-continuous on X if it is semi-continuous at every point of X.</li>
  1.4 Definition [6]: A function f: X → R is said to be lower semicontinuous (lsc) at a
  - point  $x \in X$  if for every  $\varepsilon > 0$  there exists a neighborhood U of x such that

$$f(y) - f(x) > -\mathcal{E} \quad \forall \quad y \in U .$$

We say that a function f is lsc on X if it is lsc at every point of X.

**1.5 Definition [6]:** A function  $f: X \to \mathbb{R}$  is said to be upper semicontinuous (*usc*) at a point  $x \in X$  if for every  $\varepsilon > 0$  there exists a neighborhood G of x such that

 $f(y) - f(x) < \mathcal{E} \quad \forall y \in G.$ 

We say that a function f is usc on X if it is usc at every point of X.

- **1.6 Definition [4]:** A function  $f: X \to \mathbb{R}$  is said to be lower semi-quasicontinuous (lsqc) at a point  $x \in X$  if for every  $\varepsilon > 0$  and every neighborhood U of x there exists a non-empty open set  $W \subset U$  such that  $f(y) f(x) > -\varepsilon \forall y \in W$ . We say that a function f is lsqc on X if it is lsqc at every point of X.
- **1.7 Definition [4]:** A function  $f: X \to \mathbb{R}$  is said to be upper semi-quasicontinuous (usqc) at a point  $x \in X$  if for every  $\varepsilon > 0$  and every neighborhood U of x there exists a non-empty open set  $W \subset U$  such that  $f(y) f(x) < \varepsilon \forall y \in W$ .

We say that a function f is usqc on X if it is usqc at every point of X.

1.8 Definition [5]: A function f: X → R is said to be cliquish at a point x∈ X if for every ε>0 and every neighborhood U of x there exists a non-empty open set W⊂U such that |f(y)-f(z)| < ε ∀ y, z∈ W</li>

We say that a function f is cliquish on X if it is cliquish at every point of X.

**1.9 Definition:** Let  $f:[a,b] \to \mathbb{R}$ . We define f(a-) = f(a) and f(b+) = f(b).

We say that f(p+) exists at  $p \in [a,b)$  and we write f(p+) = L, where

 $L \in \mathbb{R}$  if for every  $\varepsilon > 0$  there exists a  $\delta > 0$  such that

 $|f(x) - L| < \varepsilon \quad \forall x \in (p, p + \delta) \subset [a, b]$ . Similarly for  $p \in (a, b]$  we write

# Notes

 $f(p-) = l \in \mathbb{R}$  if for every  $\varepsilon > 0$  there exists a  $\delta > 0$  such that

 $|f(x)-l| < \varepsilon \quad \forall x \in (p-\delta, p) \subset [a,b].$ 

- **1.10 Definition:** A function  $f : [a,b] \to \mathbb{R}$  is said to be quasicontinuous at a point  $p \in [a,b]$  if f(p+) and f(p-) exist.
- **1.11 Definition:** A function  $f:[a,b] \to \mathbb{R}$  is said to be symmetrically continuous at a point  $x \in [a,b]$  if  $\lim_{h \to 0} [f(x+h) f(x-h)] = 0$ .
- **1.12 Definition:** A function  $f : [a,b] \to \mathbb{R}$  is Darboux continuous if for all  $p,q \in [a,b]$  and for each c between f(p) and f(q) there is an x between p and q such that f(x) = c.
- **1.13 Definition[2]:** An operator *P* on a linear space L is said to be sublinear if (*i*)  $P(x+y) \le P(x) + p(y) \quad \forall x, y \in L \text{ and } (ii) \quad P(\lambda x) = \lambda P(x) \text{ for any}$ positive real number  $\lambda$  and every  $x \in L$ .

III. Relations among  $f_{\odot}, {}_{\odot}f, f^{\circ}$  and  ${}^{\circ}f$ 

In this section the relations between the rising sun function and its analogues that are introduced are presented in the following propositions.

- **2.1 Proposition:** For  $f \in \mathcal{B}$ , (a)  $f^{\circ} = -(-f)_{\circ}$  and (b)  $\circ f = -_{\circ}(-f)$ .
- **2.2 Proposition:** For  $f \in \mathcal{B}$ , (a)  $(f_{\odot})_{\odot} = f_{\odot}$  (b)  $_{\odot}(_{\odot}f) = _{\odot}f$ (c)  $(f^{\circ})^{\circ} = f^{\circ}$  (d)  $^{\circ}(_{\circ}f) = ^{\circ}f$
- **2.3 Proposition:** For  $f \in \mathcal{B}$  and  $x \in [a,b]$ ,
  - $(i) _{\circ} (f_{\circ})(x) = f_{\circ}(a) \qquad (ii) (f_{\circ})^{\circ}(x) = f_{\circ}(b) \qquad (iii) ^{\circ} (f_{\circ})(x) = f_{\circ}(x)$  $(iv) (_{\circ}f)_{\circ}(x) = _{\circ}f(b) (v) (_{\circ}f)^{\circ}(x) = _{\circ}f(x) \qquad (vi) ^{\circ} (_{\circ}f)(x) = _{\circ}f(a)$  $(vii) (f^{\circ})_{\circ}(x) = f^{\circ}(b) \qquad (viii) _{\circ} (f^{\circ})(x) = f^{\circ}(x) \qquad (ix) ^{\circ} (f^{\circ})(x) = f^{\circ}(a)$  $(x) (^{\circ}f)_{\circ}(x) = ^{\circ}f(x) \qquad (xi) _{\circ} (^{\circ}f)(x) = ^{\circ}f(a) \qquad (xii) (^{\circ}f)^{\circ}(x) = ^{\circ}f(b)$
- **2.4 Remark:** In view of the previous propositions it is enough to investigate the properties of the rising sun function and the properties of  $_{\odot}f$ ,  $f^{\odot}$  and  $^{\odot}f$  follow analogously.

# $N_{otes}$

## IV. Characterisations of $f_{\odot}, {}_{\odot}f, f^{\circ}$ and ${}^{\circ}f$

# **3.1 proposition:** For $f \in \mathcal{B}$ , $f_{\odot}$ is the smallest decreasing function dominating f

More precisely,

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- (a)  $f_{\odot}(x) \ge f(x) \quad \forall x \in [a,b]$
- (b)  $f_{\odot}$  is decreasing on [a,b]
- (c) If g satisfies (a) and (b) above, then  $f_{\odot}(x) \le g(x) \quad \forall x \in [a,b]$ .

#### **3.2 Proposition:** For $f \in \mathcal{B}$ , $_{\odot}f$ is the smallest increasing function dominating

f. More precisely, (a)  $_{\odot}f(x) \ge f(x) \forall x \in [a,b]$ 

(b)  $_{\odot}f$  is increasing on [a,b]

(c) If g satisfies (a) and (b) above, then  $_{\odot}f(x) \le g(x) \quad \forall x \in [a,b].$ 

**3.3 Proposition:** For  $f \in \mathcal{B}$ ,

- (a)  $f^{\circ}$  is the largest increasing function such that  $f^{\circ}(x) \le f(x) \quad \forall x \in [a,b]$
- (b)  $^{\circ}f$  is the largest decreasing function such that  $^{\circ}f(x) \le f(x) \quad \forall x \in [a,b]$

V. THE RISING SUN OPERATOR

**4.1 Definition:** Define  $T: \mathcal{B} \to \mathcal{B}$  by  $Tf = f_{\odot}$ . We call this operator T, the rising sun operator on  $\mathcal{B}$ .

- **4.2 Proposition:** The rising sun operator T is sublinear on  $\mathcal{B}$ . More precisely,
  - $(a) T(f+g) \leq Tf + Tg \ \forall \ f,g \in \mathcal{B}$
  - (b)  $T(\lambda f) = \lambda T f$  for every real number  $\lambda > 0$  and every  $f \in \mathcal{B}$ .

**Proof:** Let  $f, g \in \mathcal{B}$  and  $x \in [a, b]$ .

(a) For  $y \in [x,b]$ ,  $(f+g)(y) = f(y) + g(y) \le f_{\odot}(x) + g_{\odot}(x) = (f_{\odot} + g_{\odot})(x)$ 

$$\Rightarrow (f+g)(y) \le (f_{\odot} + g_{\odot})(x) \ \forall \ y \in [x,b]$$
$$\Rightarrow (f+g)_{\odot}(x) \le (f_{\odot} + g_{\odot})(x) \ \forall \ x \in [a,b]$$

Hence  $T(f+g) \leq Tf + Tg \quad \forall f, g \in \mathcal{B}$ .

(b) Suppose that  $\lambda$  is a positive real number and  $f \in \mathcal{B}$ 

Then

$$\begin{aligned} (\lambda f)_{\odot}(x) &= \sup \{ (\lambda f)(y) / x \le y \le b \} \\ &= \lambda \sup \{ f(y) / x \le y \le b \} \\ &= \lambda f_{\odot}(x) \\ T \ \lambda f \ &= \lambda T f. \end{aligned}$$

**4.3 Remark:** From the following example, it is clear that  $T(f+g) \neq Tf + Tg$ .

**4.4 Example:** Define  $f:[0,1] \to \mathbb{R}$  and  $g:[0,1] \to \mathbb{R}$  by f(x) = x and

$$g(x) = 1 - x \quad \forall x \in [0, 1] \quad \Rightarrow (f + g)(x) = 1. \forall x \in [0, 1]$$

Then

$$(f+g)_{\odot}(x) = 1$$
 and  $(f_{\odot} + g_{\odot})(x) = 2 - x \quad \forall x \in [0,1]$ 

 $N_{otes}$ 

Hence  $T(f+g) \neq Tf + Tg$ .

- **4.5 Proposition:** For  $f \in \mathcal{B}$  and  $k \in \mathbb{R}$ , T(f+k) = Tf + k.
- **4.6 Proposition:** For  $f, g \in \mathcal{B}$ , (a)  $f \le g \Rightarrow Tf \le Tg$  (b)  $T(f \lor g) = Tf \lor tg$ where  $(f \lor g)(x) = \max\{f(x), g(x)\}.$
- **4.7 Proposition:** If  $\{f_{\alpha} \mid \alpha \in \Delta\}$  is a collection of functions in  $\mathcal{B}$  and if

$$\bigvee_{\alpha \in \Delta} f_{\alpha} = \sup \{ f_{\alpha} \mid \alpha \in \Delta \} \text{ exists in } \mathcal{B} \text{ then } T\Big(\bigvee_{\alpha \in \Delta} f_{\alpha}\Big) = \bigvee_{\alpha \in \Delta} T(f_{\alpha}).$$

4.8 Remark: From the following example it can be observed that

$$T(f \wedge g) \neq Tf \wedge Tg$$
, where  $(f \wedge g)(x) = \min\{f(x), g(x)\}$ .

**4.9 Example:** Define  $f:[0,1] \to \mathbb{R}$  and  $g:[0,1] \to \mathbb{R}$  by

$$f(x) = \begin{cases} 1 - 2x & \text{if } 0 \le x \le \frac{1}{2} \\ 2x - 1 & \text{if } \frac{1}{2} \le x \le 1 \end{cases} \quad \text{and} \quad g(x) = \begin{cases} 2x & \text{if } 0 \le x \le \frac{1}{2} \\ -2x + 2 & \text{if } \frac{1}{2} \le x \le 1 \end{cases}$$

Then 
$$f_{\odot}(x) = 1 \ \forall \ x \in [0,1]$$
 and  $g_{\odot}(x) = \begin{cases} 1 & \text{if } 0 \le x \le \frac{1}{2} \\ -2x+2 & \text{if } \frac{1}{2} \le x \le 1 \end{cases}$ 

$$\Rightarrow (f_{\odot} \land g_{\odot})(x) = g_{\odot}(x) \ \forall \ x \in [0,1]$$

Also 
$$(f \land g)(x) = \begin{cases} 2x & \text{if } 0 \le x \le \frac{1}{4} \\ -2x+1 & \text{if } \frac{1}{4} \le x \le \frac{1}{4} \\ 2x-1 & \text{if } \frac{1}{2} \le x \le \frac{3}{4} \\ -2x+2 & \text{if } \frac{3}{4} \le x \le 1 \end{cases}$$

$$\Rightarrow (f \land g)_{\odot}(x) = \begin{cases} \frac{1}{2} & \text{if } 0 \le x \le \frac{3}{4} \\ -2x+2 & \text{if } \frac{3}{4} \le x \le 1 \end{cases}$$

Hence  $(f \wedge g)_{\odot}(x) \neq (f_{\odot} \wedge g_{\odot})(x)$ .

**4.10 Proposition:** T is continuous on  $\mathcal{B}$ .

**Proof:** Let  $f_n \in \mathcal{B}$  for  $n \in \mathbb{N}$  and  $f_n \to f$  uniformly on [a,b]. Then  $f \in \mathcal{B}$ .

Let  $\varepsilon > 0$  be given. Then there exists a positive integer N such that

$$\begin{split} n \geq N \Rightarrow \left| f_n(x) - f(x) \right| &< \mathcal{E} \ \forall \ x \in [a, b] \\ \Rightarrow -\mathcal{E} < f_n(x) - f(x) < \mathcal{E} \ \forall \ x \in [a, b] \\ \Rightarrow f(x) - \mathcal{E} < f_n(x) < f(x) + \mathcal{E} \ \forall \ x \in [a, b]. \end{split}$$

Let  $x \in [a,b]$  and choose  $y \in [x,b]$ . Then  $y \in [a,b]$ .

$$\Rightarrow f(y) - \varepsilon < f_n(y) < f(y) + \varepsilon \quad \forall n \ge N$$
$$n \ge N \Rightarrow f_n(y) < f(y) + \varepsilon \le f_{\odot}(x) + \varepsilon$$
$$\Rightarrow f_n(y) < f_{\odot}(x) + \varepsilon \quad \forall y \in [x, b]$$

 $\Rightarrow (f_n)_{\odot}(x) \le f_{\odot}(x) + \mathcal{E}.$ 

Similarly  $f_{\odot}(x) - \mathcal{E} < (f_n)_{\odot}(x)$ .

Hence  $|(f_n)_{\odot}(x) - f_{\odot}(x)| \le \varepsilon \quad \forall n \ge N$  and for every  $x \in [a,b]$ .

Hence  $(f_n)_{\odot} \to f_{\odot}$  uniformly on [a,b].

Hence  $f_n \to f$  in  $\mathscr{B} \Longrightarrow Tf_n \to Tf$  in  $\mathscr{B}$ .

 $\Rightarrow$  *T* is continuous on  $\mathscr{B}$ .

**4.11Proposition:** (a)  $T^n f = f_{\odot} \quad \forall f \in \mathcal{B}$  and for every  $n \in \mathbb{N}$ 

(b) For  $f \in \mathcal{B}$ , the cycle of T is the set  $\{f, f_{\odot}\}$ .

- **4.12 Proposition:** If  $f \in \mathcal{B}$  and f is monotonically decreasing then f is a fixed point of T.
- **4.13 Remark:** The set of all fixed points of T is the set of all monotonically decreasing functions on [a,b]. Let F denote the set of all fixed points of T

Then  $F = \{ f \in \mathcal{B} / Tf = f \} = \{ f \in \mathcal{B} / f \text{ is decreasing} \}.$ 

Notes

- **4.14 Proposition:** The set F of all fixed points of the rising sun operator T is closed in  $\mathcal{B}$ .
- **4.15 Proposition:** Fix  $f \in F$ . Let  $F^* = \{f \in \mathcal{B} / Tf = f\}$ .
  - Then (a)  $F^*$  is closed in  $\mathcal{B}$ .

(b)  $(F^*, \leq)$  is a  $\vee$ -semilattice under the relation  $\leq$  defined on  $F^*$  by

 $f \leq g \iff f(x) \leq g(x) \ \forall x \in [a,b]$ 

**4.16 Proposition:** The operator *T* is bounded. More precisely,

(i)  $||Tf|| \le ||f|| \forall f \in \mathcal{B}$  and (ii) ||T|| = 1

VI. INVARIANT PROPERTIES

**5.1 Proposition:** Let  $f \in \mathcal{B}$ . If f is usc at a point  $x \in [a,b]$  then so

is  $Tf = f_{\odot}$ .

Notes

**Proof:** Let  $\varepsilon > 0$  be given and  $x \in [a,b]$ . Since f is usc at  $x \in [a,b]$ , there exists a  $\delta > 0$  such that  $f(t) - f(x) < \varepsilon \quad \forall t \in (x - \delta, x + \delta) \cap [a,b] = U$ .

$$\Rightarrow f_{\odot}(x) + \mathcal{E} \geq f(x) + \mathcal{E} > f(t) \ \forall t \in U$$

$$\Rightarrow f_{\odot}(x) + \mathcal{E} > f(t) \ \forall t \in U$$

Suppose that x < y.

$$\Rightarrow f_{\odot}(x) + \varepsilon > f_{\odot}(x) \ge f_{\odot}(y) \ge f(y)$$
$$\Rightarrow f_{\odot}(x) + \varepsilon > f(y).$$

If x < t and  $y \in [t, b] \Rightarrow x < y$ 

$$\Rightarrow f_{\odot}(x) + \mathcal{E} > f(y)$$
$$\Rightarrow f_{\odot}(x) + \mathcal{E} \ge f_{\odot}(t)$$

Suppose that  $t \le x$ . Then  $y \in [t,b] \Rightarrow t \le y \le x$  or x < y < b.

If  $t \le y \le x$  then  $y \in U \implies f_{\odot}(x) + \mathcal{E} > f(y) \implies f_{\odot}(x) + \mathcal{E} > f_{\odot}(t)$ .

If  $x < y \le b$  then  $f_{\odot}(x) + \varepsilon > f_{\odot}(t)$ . Hence  $f_{\odot}(x) + \varepsilon > f_{\odot}(t) \quad \forall t \in U$ .

 $\Rightarrow f_{\odot}$  is usc at x.

**5.2 Proposition:** Let  $f \in \mathcal{B}$ . If f is *lsc* at a point  $x \in [a,b]$  then so

is 
$$Tf = f_{\odot}$$
.

**Proof:** Let  $\varepsilon > 0$  be given and  $x \in [a,b]$ . Since f is lsc at  $x \in [a,b]$ , there

exists a  $\delta_1 > 0$  such that  $f(t) - f(x) > \frac{-\varepsilon}{2} \quad \forall t \in (x - \delta_1, x + \delta_1) \cap [a, b] = U$   $\Rightarrow f_0(t) + \frac{\varepsilon}{2} \ge f(t) + \frac{\varepsilon}{2} > f(x) \quad \forall t \in U$  $\Rightarrow f_0(t) + \frac{\varepsilon}{2} > f(x) \quad \forall t \in U$ 

Since  $f_{\odot}(x) - \frac{\varepsilon}{2}$  is not an upper bound of the set  $\{f(y) | x \le y \le b\}$ , there

exists a point  $y \in [x,b]$  such that  $f(y) > f_{\odot}(x) - \frac{\varepsilon}{2}$ .

If 
$$y = x$$
 then  $f_{\odot}(t) + \frac{\varepsilon}{2} > f(x) > f_{\odot}(x) - \frac{\varepsilon}{2} \quad \forall t \in U$   
$$\Rightarrow f_{\odot}(t) + \varepsilon > f_{\odot}(x) \quad \forall t \in U.$$

 $y \neq x$ . If x < z < y then  $f_{\odot}(z) \ge f_{\odot}(y) \ge f(y) > f_{\odot}(x) - \frac{\varepsilon}{2}$ 

$$\Rightarrow f_{\odot}(z) > f_{\odot}(x) - \frac{\varepsilon}{2} > f_{\odot}(x) - \varepsilon \,.$$

If  $z \le x$  then  $f_{\odot}(z) \ge f_{\odot}(x) > f_{\odot}(x) - \mathcal{E}$ .

Choose  $\delta > 0$  such that  $a \le x - \delta < x < x + \delta \le y$ .

Then  $f_{\odot}(z) > f_{\odot}(x) - \mathcal{E} \quad \forall z \in (x - \delta, x + \delta) \cap [a, b]$ . Hence  $f_{\odot}$  is *lsc* at x.

**5.3 Corollary:** Let  $f \in \mathcal{B}$ . If f is continuous at a point  $x \in [a,b]$ , then

 $Tf = f_{\odot}$  is continuous at x.

**5.4 Proposition:** For any  $f \in \mathcal{B}$ ,  $Tf = f_{\odot}$  is lsqc at every  $x \in (a,b]$ .

**Proof:** Let  $f \in \mathcal{B}$  and  $x \in (a,b]$ . Let  $\varepsilon > 0$  be given and let U be a neighborhood of x in [a,b]. Then there exists a  $\delta > 0$  such that  $(x-\delta, x+\delta) \cap [a,b] \subset U$ . Choose  $W = (x-\delta, x) \cap [a,b]$ .

W is a non-empty open subset of U.

$$y \in W \Longrightarrow x - \delta < y < x$$
$$\Rightarrow f_{\odot}(y) \ge f_{\odot}(x)$$
$$\Rightarrow f_{\odot}(y) + \varepsilon > f_{\odot}(y) \ge f_{\odot}(x)$$

Notes

Hence  $f_{\odot}(y) + \mathcal{E} > f_{\odot}(x) \quad \forall y \in W$ .

Thus for every  $\varepsilon > 0$  and every neighborhood U of x there exists a nonempty open set  $W \subset U$  such that  $f_{\odot}(y) + \varepsilon > f_{\odot}(x) \quad \forall y \in W$ .

**5.5 Proposition:** Let  $f \in \mathcal{B}$ . If  $f:[a,b] \to \mathbb{R}$  is *lsqc* at the point *a* then

 $Tf = f_{\odot}$  is *lsqc* at *a*.

**Proof:** Let  $f \in \mathcal{B}$  and  $\varepsilon > 0$  be given. Let U be a neighborhood of a in [a,b]. Since f is lsqc at a, there exists a non-empty open set  $W \subset U$  such that  $f(t) - f(x) > -\varepsilon \quad \forall t \in W$ .

$$\Rightarrow f_{\odot}(t) \ge f(t) > f(a) - \mathcal{E} \ \forall t \in W$$
$$\Rightarrow f_{\odot}(t) > f(a) - \mathcal{E} \ \forall t \in W.$$

Since  $f_{\odot}(a) - \varepsilon$  is not an upper bound of  $\{f(y) | a \le y \le b\}$ , there exists a point  $y \in [a,b]$  such that  $f_{\odot}(a) - \varepsilon < f(y)$ .

If y = a then  $f_{\odot}(t) + \mathcal{E} > f(a) > f_{\odot}(a) - \mathcal{E} \quad \forall t \in W$ .

Suppose that  $y \neq a$ .

Since  $a \in U$  and U is open in [a,b], there exists a  $\delta_1 > 0$  such that

 $[a, a + \delta_1) \subset U$ . Choose  $\delta_2 > 0$  such that  $a < a + \delta_2 < y$ .

Put  $\delta = \min{\{\delta_1, \delta_2\}}$  and  $W_1 = (a, a + \delta)$ . Clearly,  $W_1$  is a non-empty open set such that  $W_1 \subset U$ .

Then  $z \in W_1 \Longrightarrow z < y$ 

$$\Rightarrow f_{\odot}(z) \ge f_{\odot}(y) \ge f(y) > f_{\odot}(a) - \mathcal{E}$$

Hence  $f_{\odot}(z) + \mathcal{E} > f_{\odot}(a) \forall z \in W_1$ 

 $\Rightarrow f_{\odot}$  is *lsqc* at *a*.

**5.6 Proposition:** Let  $f \in \mathcal{B}$ . Then

(a)  $f_{\odot}$  is usqc at every  $x \in [a,b)$ .

(b) If  $f:[a,b] \to \mathbb{R}$  is useque at b then so is  $f_{\odot}$ .

 $N_{otes}$ 

**<sup>5.7</sup> Proposition [5]:** Let  $f:[a,b] \to \mathbb{R}$  and  $p \in [a,b]$ . If f(p+) exists then f is cliquish at p.

- **5.8 Corollary:** For any  $f \in \mathcal{B}$ ,  $Tf = f_{\odot}$  is cliquish on [a,b].
- **5.9 Proposition:** For any  $f \in \mathcal{B}$ ,  $Tf = f_{\odot}$  is quasicontinuous on [a,b].

#### VII. VARIANT PROPERTIES

6.1 Symmetric continuity: It is not necessary that the rising sun function of a symmetrically continuous function is symmetrically continuous. For example, define  $f : [-1,1] \rightarrow \mathbb{R}$  by

$$f(x) = \begin{cases} x^2 & \text{if } x \neq 0\\ 0 & \text{if } x = 0 \end{cases}$$

Then

Clearly f is symmetrically continuous on [-1,1], but  $f_{\odot}$  is not.

 $f_{\odot}(x) = \begin{cases} x^2 & \text{if } -1 \le x \le 0\\ 1 & \text{if } 0 < x \le 1 \end{cases}$ 

**6.2 Semi-continuity:** The semi-continuity of f need not imply the semi-continuity of  $f_{\odot}$  as is evident from the following example.

Define  $f:[-1,1] \to \mathbb{R}$  by

$$f(x) = \begin{cases} 0 & \text{if } x = 0\\ (x+1)2^{\frac{-2}{x}} & \text{if } 0 < x \le 1\\ x+1 & \text{if } -1 \le x < 0 \end{cases}$$

Then

 $f_{\odot}(x) = \begin{cases} 1 & if -1 \le x \le 0\\ \frac{1}{2} & if 0 < x \le 1 \end{cases}$ 

Clearly f is semi-continuous on [-1,1]. But  $f_{\odot}$  is not semi-continuous at x=0.

- 6.3 Darboux continuity: It is not necessary that the rising sun function of a Darboux continuous function is Darboux continuous. The function
  f:[-1,1] → ℝ defined in the above example is Dourbox continuous on[-1,1]
  but its rising sun function is not Dourbox continuous.
- **6.4 Differentiability:** The rising sun function of a differentiable function is not necessarily differentiable as can be observed from the following example.

Notes

Define  $f:[0,1] \to \mathbb{R}$  by

Notes



Clearly f is differentiable at a = 0.3732, but  $f_{\odot}$  is not differentiable at this point.

**6.5 Pointwise Convergence:** If  $\{f_n\}$  converges pointwise to f on [a,b], it is not necessary that  $\{Tf_n\}$  converges to Tf as can be seen from the following example. Define  $f_n:[0,1] \to \mathbb{R}$  by

$$f_n(x) = \begin{cases} 1 & \text{if} \quad x \in \left\{\frac{1}{n} / n \in \mathbb{N}\right\} \\ 0 & \text{if} \quad x \notin \left\{\frac{1}{n} / n \in \mathbb{N}\right\} \end{cases}$$

Then  $\{f_n\}$  converges pointwise to 0.

But  $(f_n)_{\odot}(x) = \begin{cases} 1 & if \quad x = 0 \\ 0 & if \quad 0 < x \le 1 \end{cases}$  does not converge to 0.

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# Some Statistical Properties of Exponentiated Weighted Weibull Distribution

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*Abstract-* This article basically focused on some statistical properties of exponentiated-weighted weibull model which of course numerous authors have written one thing or the other on exponential weibull distribution and not on exponential weighted weibull. This model is established with a view to obtaining a model that is better than both weighted weibull and weibull distribution in terms of the estimate of their characteristics and their parameters using the logit of Beta by Jones (2004). The weighted weibull distribution is proposed by Mahdy (2013) with an additional parameter called "sensitive skewness parameter". Some basic properties of the proposed model including moments and moment generating function (first and second moments about the origin even with standard deviation are derive), survival rate function, hazard rate function, asymptotic behaviours, and the estimation of parameters have been studied. The result from the new model is better representativeness in data and its flexibility and shape.

Keywords: exponentiated-weighted weibull, hazard rate, moments, weighted-weibull, survival rate.

GJSFR-F Classification : MSC 2010: 97K80, 35B40



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Notes

# Some Statistical Properties of Exponentiated Weighted Weibull Distribution

Badmus, N. Idowu<sup> a</sup> & Bamiduro, T. Adebayo<sup>o</sup>

Abstract- This article basically focused on some statistical properties of exponentiated-weighted weibull model which of course numerous authors have written one thing or the other on exponential weibull distribution and not on exponential-weighted weibull. This model is established with a view to obtaining a model that is better than both weighted weibull and weibull distribution in terms of the estimate of their characteristics and their parameters using the logit of Beta by Jones (2004). The weighted weibull distribution is proposed by Mahdy (2013) with an additional parameter called "sensitive skewness parameter". Some basic properties of the proposed model including moments and moment generating function (first and second moments about the origin even with standard deviation are derive), survival rate function, hazard rate function, asymptotic behaviours, and the estimation of parameters have been studied. The result from the new model is better representativeness in data and its flexibility and shape.

Keywords: exponentiated-weighted weibull, hazard rate, moments, weighted-weibull, survival rate.

#### I. INTRODUCTION

In recent time, numerous researchers had used weibull distribution as an alternative to some distribution e.g Gamma and Log-normal distribution in reliability engineering and life testing. The weibull distribution is a well known common distribution and has been a powerful probability distribution in reliability analysis, while weighted distributions are used to adjust the probabilities of the events as observed and recorded. Mahdy applied Azzalini's method to the weibull distribution that produced a new class of weighted weibull distribution as  $WW(\lambda, \beta, \alpha)$  distribution with an additional parameter called "Sensitive Skewness Parameter" and the sensitive skewness parameter governs essentially the shape of the probability density function of the  $WW(\lambda, \beta, \alpha)$  distribution.

The probability density and the cumulative density function (pdf and cdf) of the new class of weighted weibull distribution by Mahdy (2013) is given by

$$f_{y|\{\lambda,\beta,\alpha\}|}(y) = \frac{\lambda\beta(1+\alpha^{\beta})y^{\beta-1}e^{-\lambda y^{\beta}}(1-e^{-\lambda(\alpha y)^{\beta}})}{\alpha^{\beta}}, \text{ for } y > 0$$
(1)

and

$$F_{y|\{\lambda,\beta,\alpha\}|}(y) = \frac{[(1+\alpha^{\beta})(1-e^{-\lambda y^{\beta}})+e^{-\lambda y^{\beta}(1+\alpha^{\beta})}-1]}{\alpha^{\beta}}$$
(2)

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Notes

b=1, 
$$\lambda$$
 =0.5,  $\beta$  = 3.5,  $\alpha$  = 2

Authors on exponenetiated-weighted weibull are very few. The aim of this article is to introduce and investigate this distribution on its statistical properties. The paper is divided as follows: In section 2, we present the proposed distribution exponentiatedweighted weibull distribution. Moments, and moment generating function is studied in section 3, section 4, discussed on the estimation of parameters mathematically and in section 5 we preset the real application to data set and section 6 concluded the research.

#### II. THE PROPOSED EXPONENTIATED-WEIGHTED WEIBULL DISTRIBUTION

Recently, many authors have studied the properties of exponentiated distributions. For instance, Gupta et al (2001) for exponential pareto, Nadarajah and Gupta (2007) for exponential gamma distribution, Mudholkar et al (1995) studied on exponentiated weibull distribution, Salem and Abo-Kasem (2011) based their research on estimation for the parameters of the exponentiated weibull distribution, Gupta and Kundu (2001) they put up a paper on exponentiated exponential etc. Azzalini (1985) first proposed a method of obtaining weighted and the method has been used extensively for several symmetric and non-symmetric distributions. Mahdy (2013) applied the method to study a new class of weighted weibull distribution with an additional parameter called "sensitive skewness parameter". More so, various extensions of weibull and exponential distribution have been proposed in literature. An extension of exponential distribution has been provided by Nadarajah and Kotz (2005) using the logit of Beta distribution and the logit of Beta distribution (the link function of the Beta generalized distribution) is introduced by Jones (2004). Since then extensive work has been done using the logit of beta distribution in literature. For instance, Gupta and Kundu (1999) proposed a generalized exponential distribution which provides an alternative to exponential and weibull distributions. Famoye et al (2005) also introduced the Beta-weibull distribution alongside its major properties and Cordeiro et al (2011) among others.

Now, letting y be a random variable form of the distribution with parameters and defined (1) and (2) using the logit of beta by Jones (2004), we then have

$$f_{EWW}^{(y)} = \frac{1}{B(a,b)} [F(y)]^{a-1} [1 - F(y)]^{b-1} f(y)$$
(3)

by setting b=1, we get

$$f_{EWW}^{(y)} = a[F(y)]^{a-1} f(y)$$
(4)

Putting expressions (1) and (2) in (4) to obtain the probability density function of Exponentiated-weighted weibull distribution



*Figure 2*: The PDF of Exponentiated Weighted Weibull Distribution with values of the parameters ( $a = 1.5, b = 1, c = \lambda = 1.5, d = \alpha = 3.5, e = \beta = 2.$ ) and is rightly skewed

Where  $a > 0, \lambda > 0, \beta > 0, \alpha > 0$  and y > 0 such that  $Y \sim EWW(a, \lambda, \beta, \alpha)$ . Equation (5) is the pdf of Exponentiated-weighted weibull distribution.

$$\operatorname{set} u(x) = \frac{\left[(1+\alpha^{\beta})\left(1-e^{-\lambda y^{\beta}}\right)+e^{-\lambda y^{\beta}}\left(1+\alpha^{\beta}\right)-1\right]}{\alpha^{\beta}}$$
$$\frac{du}{dy} = \frac{\lambda\beta\left(1+\alpha^{\beta}\right)y^{\beta-1}e^{-\lambda y^{\beta}}\left(1-e^{-\lambda(\alpha y)^{\beta}}\right)}{\alpha^{\beta}}\left[\frac{1}{\alpha^{\beta}}\left\{\left(1+\alpha^{\beta}\right)\left(1-e^{-\lambda y^{\beta}}\right)+e^{-\lambda y^{\beta}}\left(1+\alpha^{\beta}\right)-1\right\}\right]$$
(6)

substituting dy into (5), we obtain

$$f_{EWW}^{(y)} = a \left[ \frac{1}{\alpha^{\beta}} \left\{ \left( 1 + \alpha^{\beta} \right) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} \left( 1 + \alpha^{\beta} \right)} - 1 \right]^{a-1} du$$
$$u = \frac{1}{\alpha^{\beta}} \left\{ \left( 1 + \alpha^{\beta} \right) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} \left( 1 + \alpha^{\beta} \right)} - 1 \right\}$$

Now, (5) can becomes

$$f_{EWW}^{(y)} = a[U]^{a-1} \frac{du}{dy}$$
(8)

#### a) Cumulative Density Function (cdf)

The probability density function (pdf) of  $EWW(a, \lambda, \beta, \alpha)$  given in (7), then expression (7) can be written as

$$F_{EWW}^{(y)} = P(Y \le y) = \int_{0}^{y} f(u) du$$
  
=  $\int_{0}^{y} a \left[ \frac{1}{\alpha^{\beta}} \left\{ (1 + \alpha^{\beta}) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} (1 + \alpha^{\beta})} - 1 \right]^{a-1} \frac{\lambda \beta (1 + \alpha^{\beta}) y^{\beta-1} e^{-\lambda y^{\beta} (1 - e^{-\lambda (\alpha y)^{\beta}})}{\alpha^{\beta}} du$  (9)

(7)

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$$F_{EWW}^{(y)} = P(Y \le y) = \int_0^y a[U]^{a-1} du$$

 $=a\int_0^y [U]^{a-1} du$  and the cdf is obtained as

$$F_{EWW}^{(y)} = -(y)^a$$
 (10)

Notes

#### b) The Survival Rate Function

The survival rate function of the Exponentiated-weighted weibull distribution is given by  $S_{EWW}(y) = 1 - F_{EWW}(y) = 1 - \int_0^y f(u) du$ 

$$= 1 - a \int_0^y [U]^{a-1} du = 1 - [-(y)^a]$$
  

$$S_{EWW}(y) = 1 + (y)^a$$
(11)

#### c) The Hazard Rate Function

The hazard rate function of a random variable y with the pdf and cdf is defined by

$$h_{EWW}(y) = \frac{f_{EWW}(y)}{1 - F_{EWW}(y)}$$

Hence, the  $EWW(a, \lambda, \beta, \alpha)$  with  $f_{EWW}(y)$  and  $F_{EWW}(y)$  respectively defined in (4) and (10), the hazard rate function can be expressed as:

$$=\frac{a(u)^{a-1}u'}{S_{EWW}(y)}$$
(12)

where U is expression in (5) To show that  $y \xrightarrow{\lim} \infty h_{EWW}(y) = 0$  and  $y \xrightarrow{\lim} 0 h_{EWW}(y) = 0$ , we have the following

$$y \xrightarrow{lim} \infty h_{EWW}(y) = y \xrightarrow{lim} \infty \frac{a(u)^{a-1}u'}{1-(1+y)^a}$$

where,  $u' = \frac{1}{\alpha^{\beta}} \left\{ \left( 1 + \alpha^{\beta} \right) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} \left( 1 + \alpha^{\beta} \right)} - 1 \right\}$  $= y \xrightarrow{\lim} \infty \frac{a \left[ \frac{1}{\alpha^{\beta}} \left\{ (1 + \alpha^{\beta}) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} \left( 1 + \alpha^{\beta} \right)} - 1 \right]^{a-1} \frac{\lambda^{\beta} \left( 1 + \alpha^{\beta} \right) y^{\beta-1} e^{-\lambda y^{\beta} \left( 1 - e^{-\lambda (\alpha y)^{\beta}} \right)}}{a^{\beta}}}{1 - (1 + x)^{a}}$ 

For simplification on the rigorous mathematics, we take the limit of the following: When  $y \rightarrow \infty = 0$  and  $y \rightarrow 0 = 0$ 

$$= y \xrightarrow{\lim} \infty \frac{\lambda \beta (1 + \alpha^{\beta}) y^{\beta - 1} e^{-\lambda y^{\beta}} (1 - e^{-\lambda (\alpha y)^{\beta}})}{\alpha^{\beta}} = y \xrightarrow{\lim} \infty \frac{\lambda \beta (1 + \alpha^{\beta}) \infty^{\beta - 1} e^{-\lambda \infty^{\beta}} (1 - e^{-\lambda (\alpha \infty)^{\beta}})}{\alpha^{\beta}} = 0$$

$$= y \stackrel{lim}{\longrightarrow} 0 \frac{\lambda \beta (1+\alpha^{\beta}) 0^{\beta-1} e^{-\lambda 0^{\beta}} (1-e^{-\lambda (\alpha 0)^{\beta}})}{\alpha^{\beta}} = 0$$

As  $y \to \infty = 0$  and  $y \to 0 = 0$ , expression (12) above tends to  $\infty$  and 0 and equal to zero.

#### d) Asymptotic Behaviours

Following the steps in hazard function above taken  $y \xrightarrow{\lim} \infty f_{EWW}(y)$  and  $y \xrightarrow{\lim} 0 f_{EWW}(y)$  of the  $EWW(a, \lambda, \beta, \alpha)$  distribution is investigated as follows. Now from expression (5), we have

$$y \xrightarrow{lim} \infty f_{EWW}(y) = \alpha \left[ \frac{\left[ \left( 1 + \alpha^{\beta} \right) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} \left( 1 + \alpha^{\beta} \right)} \right]^{\alpha - 1}}{\alpha^{\beta}} \right]^{\alpha - 1} \frac{\lambda \beta \left( 1 + \alpha^{\beta} \right) y^{\beta - 1} e^{-\lambda y^{\beta}} \left( 1 - e^{-\lambda (\alpha y)^{\beta}} \right)}{\alpha^{\beta}}$$

taking the limit

Notes

$$= y \xrightarrow{\lim} \infty \frac{\lambda \beta (1 + \alpha^{\beta}) y^{\beta - 1} e^{-\lambda y^{\beta}} (1 - e^{-\lambda (\alpha y)^{\beta}})}{\alpha^{\beta}} = y \xrightarrow{\lim} \infty \frac{\lambda \beta (1 + \alpha^{\beta}) \infty^{\beta - 1} e^{-\lambda \infty^{\beta}} (1 - e^{-\lambda (\alpha \infty)^{\beta}})}{\alpha^{\beta}}$$
$$= 0$$
$$= y \xrightarrow{\lim} 0 \frac{\lambda \beta (1 + \alpha^{\beta}) 0^{\beta - 1} e^{-\lambda 0^{\beta}} (1 - e^{-\lambda (\alpha 0)^{\beta}})}{\alpha^{\beta}} = 0$$

From the above results as  $y \to \infty = 0$  and  $y \to 0 = 0$ , this shows that the distribution has at least a mode.

#### III. Moments and Moment Generating Function

Hosking (1990) described in their paper that when a random variable following a generalized beta generated distribution i.e  $y \sim GBG(f, a, c)$  then  $\mu'_r = E[F^{-1}K^{\frac{1}{c}}]^r$  where  $K \sim B(a, 1), c$  is a constant and  $F^{-1}(y)$  is the inverse of CDF of the weighted weibull distribution, since  $EWW(a, \lambda, \beta, \alpha)$  distribution is a special form when a=c=1. We then derive the moment generating function (mgf) of the proposed distribution  $m(t) = E(e^{ty})$  and the general rth moment of a beta generated distribution is defined by

$$\mu'_{r} = \frac{1}{B(a,1)} \int_{0}^{1} [F^{-1}(y)]^{r} [y]^{a-1} du$$
(13)

Also, using the taylor series expansion around the point  $E(y_f) = \mu_f$  to obtain

$$\mu_r' = \sum_{u=0}^r {r \choose k} [F^{-1}(\mu)]^{r-k} [F^{-1(1)}(\mu_f)]^k \sum_{k=0}^n (-1)^i {r \choose i}$$
(14)

Cordeiro et al (2011) gave an alternative series expansion for  $\mu'_r$  in terms of  $r(r, U) = E(U^r F(U)^Y)$  where k follows the parent distribution then for u = 0, 1, ...

$$\mu'_{r} = \frac{1}{B(a,1)} \sum_{i=0}^{\infty} (-1)^{i} {b-1 \choose i} r(r,i-1)$$

They further described another mgf of y for generated beta distribution as

$$M(t) = \frac{1}{B(a,1)} \sum_{i=0}^{\infty} (-1)^{i} {a-1 \choose i} \rho(t, i-1)$$
(15)  
$$\rho(t,r) = \int_{-\infty}^{\infty} e^{ty} [F(y)]^{m} f(y) dy$$

(16)

Notes

Where,

Therefore,

Substituting both probability density and cumulative density function of the weighted weibull distribution into (16), we obtain

 $M_{y}^{(t)} = \frac{1}{B(a,1)} \sum_{i=0}^{\infty} (-1)^{i} {a-1 \choose i} \int_{-\infty}^{\infty} e^{ty} [F(y)]^{(i+1)-1} f(y) \, dy$ 

$$M_{EWW(y)}^{(t)} = a \sum_{i=0}^{\infty} (-1)^{i} {\binom{a-1}{i}} \int^{e^{ty} \left[\frac{1}{\alpha^{\beta}} \left\{ \left(1+\alpha^{\beta}\right) \left(1-e^{-\lambda y^{\beta}}\right) + e^{-\lambda y^{\beta} \left(1+\alpha^{\beta}\right)} - 1\right]^{(i+1)-1} \right]}_{\frac{\lambda\beta \left(1+\alpha^{\beta}\right) y^{\beta-1} e^{-\lambda y^{\beta} \left(1-e^{-\lambda (\alpha y)^{\beta}}\right)}{\alpha^{\beta}}} dy$$

$$(17)$$

Equation (17) becomes the mgf of Exponentiated-weighted weibull distribution. Then, setting a=1 and i = 0, the same expression (17) is reduced to becomes the parent distribution. To obtain the rth moment of  $EWW(a, \lambda, \beta, \alpha)$ , the weighted weibull distribution by Mahdy (2013) and is given by

$$M_{(y)}^{(t)} = \int_{0}^{\infty} e^{ty} \frac{\lambda \beta (1 + \alpha^{\beta}) y^{\beta - 1} e^{-\lambda y^{\beta}} (1 - e^{-\lambda (\alpha y)^{\beta}})}{\alpha^{\beta}} dy$$
$$= \sum_{j=0}^{\infty} \frac{t^{j}}{j! \alpha^{\beta}} \{ \lambda^{-\frac{j}{\beta}} (1 + \alpha^{\beta}) \Gamma(\frac{j + \beta}{\beta}) (1 - (1 + \alpha^{\beta}))^{-\left(\frac{j + \beta}{\beta}\right)} \}$$
(18)

Equation (17) can be re-written as

$$M_{EWW(y)}^{(t)} = a \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} (-1)^{i} {\binom{a-1}{i}}_{j!\alpha^{\beta}} \{\lambda^{-\frac{j}{\beta}} (1+\alpha^{\beta}) \Gamma(\frac{j+\beta}{\beta}) (1-(1+\alpha^{\beta}))^{-(\frac{j+\beta}{\beta})} \}$$
(19)

The rth moment of the  $EWW(a, \lambda, \beta, \alpha)$  distribution can also be written from equation (19) as

$$\mu_r' = E(Y^r) = a \sum_{i=0}^{\infty} (-1)^i {a-1 \choose i} \frac{t^r}{r! \alpha^\beta} \{ \lambda^{-\frac{r}{\beta}} (1+\alpha^\beta) \Gamma(\frac{r+\beta}{\beta}) (1-(1+\alpha^\beta))^{-\left(\frac{r+\beta}{\beta}\right)} \}$$
(20)

Again, putting a = 1 in expression (20) leads to the rth moment of the weighted weibull model by Mahdy (2013) and is given by

$$\mu_r' = E(Y^r) = \frac{1}{\alpha^{\beta}} \{ \lambda^{-\frac{r}{\beta}} (1 + \alpha^{\beta}) \Gamma(\frac{r+\beta}{\beta}) (1 - (1 + \alpha^{\beta}))^{-\left(\frac{r+\beta}{\beta}\right)} \}$$
(21)

$$\mu_{EWW(r)}^{'} = a \sum_{i=0}^{\infty} (-1)^{i} {\binom{a-1}{i}} \frac{r}{\alpha^{\beta}} \{ \lambda^{-\frac{1}{\beta}} (1+\alpha^{\beta}) \Gamma(\frac{r+\beta}{\beta}) (1-(1+\alpha^{\beta}))^{-\left(\frac{r+\beta}{\beta}\right)} \}$$
(22)

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From (22), it is easy to obtain the first and the second mean about the origin e.g when r = 1 and the second moment when r = 2, etc. The first moment of  $EWW(a, \lambda, \beta, \alpha)$  is obtain

$$\mu_{EWW(1)}^{'} = a \sum_{i=0}^{\infty} (-1)^{i} {\binom{a-1}{i}} \frac{1}{\alpha^{\beta}} \{ \lambda^{-\frac{1}{\beta}} (1+\alpha^{\beta}) \Gamma(\frac{1+\beta}{\beta}) (1-(1+\alpha^{\beta}))^{-(\frac{1+\beta}{\beta})} \}$$
(23)

The second moment can also be obtained as follows:

$$\mu_{2}^{'} = V(a, 1, \lambda, \beta, \alpha)^{(y)} = V_{1} - V_{2}$$
(24)

where,

 $N_{otes}$ 

$$V_1 = E(\alpha, \lambda, \beta, \alpha)^{(Y^2)} = \alpha \sum_{i=0}^{\infty} (-1)^i {\binom{\alpha-1}{i}} \frac{1}{\alpha^{\beta}} \{\lambda^{-\frac{2}{\beta}} (1+\alpha^{\beta}) \Gamma(\frac{2+\beta}{\beta}) (1-(1+\alpha^{\beta}))^{-\left(\frac{2+\beta}{\beta}\right)} \}$$

$$V_2 = E(\alpha, \lambda, \beta, \alpha)^{(Y)} = \left(\alpha \sum_{i=0}^{\infty} (-1)^i \binom{\alpha-1}{i} \frac{1}{\alpha^{\beta}} \{\lambda^{-\frac{2}{\beta}} (1+\alpha^{\beta}) \Gamma\left(\frac{1+\beta}{\beta}\right) \left(1 - \left(1+\alpha^{\beta}\right)\right)^{-\left(\frac{1+\beta}{\beta}\right)} \right\} \right)^2$$

Likewise, the standard deviation is given by

$$SD_{EWW}$$
.  $(\alpha, \lambda, \beta, \alpha)^{(Y)} = \sqrt{V_1 - V_2}$ 

#### IV. ESTIMATION OF PARAMETER

We show the maximum likelihood estimate (MLEs) of the parameter of  $EWW(a, \lambda, \beta, \alpha)$  distribution mathematically following Cordeiro et al (2011) and Shittu and Adepoju (2013) studied on the log-likelihood function for  $\omega = (a, c, \varphi)$ , where  $\varphi = (\lambda, \beta, \alpha)$  and setting  $\omega$  to be a vector of parameter and is given by

$$L(\omega) = nlogc - nlog[B(a, 1)] + \sum_{i=1}^{n} \log[f(y; \varphi)] + (a - 1) \sum_{i=1}^{n} \log[F(y; \varphi)]$$
(25)

Note that, 
$$b = c = 1$$
 (24) becomes  $\theta = (a, \varphi)$ 

$$L(\omega) = -nlog(a, 1) + \sum_{i=1}^{n} \log[f(y; \varphi)] + (a - 1) \sum_{i=1}^{n} \log [F(y; \varphi)]$$
(26)

$$f(y;\varphi) = \frac{\lambda\beta(1+\alpha^{\beta})y^{\beta-1}e^{-\lambda y^{\beta}}(1-e^{-\lambda(\alpha y)^{\beta}})}{\alpha^{\beta}} \quad \text{and} \quad$$

$$F(y;\varphi) = \frac{1}{\alpha^{\beta}} \left\{ \left( 1 + \alpha^{\beta} \right) \left( 1 - e^{-\lambda y^{\beta}} \right) + e^{-\lambda y^{\beta} \left( 1 + \alpha^{\beta} \right)} - 1 \right\}$$

$$L_{EWW}^{(\omega)} = -nlog(a,1) + \sum_{i=1}^{n} \log\left[\frac{\lambda\beta(1+\alpha^{\beta})y^{\beta-1}e^{-\lambda y^{\beta}}(1-e^{-\lambda(\alpha y)\beta})}{\alpha^{\beta}}\right] + (a-1)\sum_{i=1}^{n} \log\left[\frac{1}{\alpha^{\beta}}\left\{\left(1+\alpha^{\beta}\right)\left(1-e^{-\lambda y^{\beta}}\right) + e^{-\lambda y^{\beta}(1+\alpha^{\beta})} - 1\right]\right]$$

$$(27)$$

For determining the MLE of  $a, \lambda, \beta, \alpha$ , we take the partial derivative of the (27) with respect to  $(a, \lambda, \beta, \alpha)$  as follows:

$$\frac{L_{EWW}}{\partial a} = -nlog(a,1) + (a-1)\sum_{y=1}^{n} \log\left[\frac{1}{\alpha^{\beta}}\left\{\left(1+\alpha^{\beta}\right)\left(1-e^{-\lambda y^{\beta}}\right) + e^{-\lambda y^{\beta}\left(1+\alpha^{\beta}\right)} - 1\right]\right] (28)$$

$$\frac{L_{EWW}}{\partial \lambda}^{(\omega)} = \sum_{y=1}^{n} log \left[ \frac{\frac{\partial}{\partial \lambda} \frac{\lambda \beta \left(1+\alpha \beta\right) y \beta - 1 e^{-\lambda y \beta} \left(1-e^{-\lambda (\alpha y) \beta}\right)}{\alpha \beta}}{\frac{\lambda \beta \left(1+\alpha \beta\right) y \beta - 1 e^{-\lambda y \beta} \left(1-e^{-\lambda (\alpha y) \beta}\right)}{\alpha \beta}} \right] + (a-1) \sum_{y=1}^{n} log \left[ \frac{\frac{\partial}{\partial \lambda} \left(\frac{1}{\alpha \beta} \left\{ \left(1+\alpha \beta\right) \left(1-e^{-\lambda y \beta}\right) + e^{-\lambda y \beta} \left(1+\alpha \beta\right) - 1\right)}{\frac{1}{\alpha \beta} \left\{ \left(1+\alpha \beta\right) \left(1-e^{-\lambda y \beta}\right) + e^{-\lambda y \beta} \left(1+\alpha \beta\right) - 1} \right]} \right]$$
(29)

Notes

$$\frac{L_{EWW}}{\partial\beta} = \sum_{y=1}^{n} log \left[ \frac{\frac{\partial}{\partial\beta} \lambda\beta \left(1+\alpha\beta\right) ye^{-\lambda y\beta} \left(1-e^{-\lambda(\alpha y)\beta}\right)}{\alpha\beta}}{\lambda\beta \left(1+\alpha\beta\right) y^{\beta-1}e^{-\lambda y\beta} \left(1-e^{-\lambda(\alpha y)\beta}\right)} \right] + a^{2\beta} \left[ \frac{\partial}{\partial\beta} \left(1-\frac{1}{\alpha\beta} \left\{ \left(1+\alpha\beta\right) \left(1-e^{-\lambda y}\right) + e^{-\lambda y\beta} \left(1+\alpha\beta\right) - 1 \right)}{\alpha\beta} \right] \right]$$
(30)

$$(a-1)\sum_{y=1}^{n}\log\left[\frac{\frac{\partial\beta}{\partial\beta}(1-\frac{\alpha}{\alpha\beta}(1+\alpha\beta)(1-e^{-\lambda y\beta})+e^{-\lambda y\beta}(1+\alpha\beta)-1)}{\frac{1}{\alpha\beta}\{(1+\alpha\beta)(1-e^{-\lambda y\beta})+e^{-\lambda y\beta}(1+\alpha\beta)-1\}}\right]$$
(30)

$$\frac{L_{EWW}}{\partial \alpha} = \sum_{y=1}^{n} log[\frac{\frac{\partial \lambda \beta (1+\alpha \beta) y^{\beta-1} e^{-\lambda y^{\beta} (1-e^{-\lambda (\alpha y)^{\beta}})}}{\frac{\lambda \beta (1+\alpha \beta) y^{\beta-1} e^{-\lambda y^{\beta} (1-e^{-\lambda (\alpha y)^{\beta}})}{\alpha \beta}}] +$$

$$(a-1)\sum_{y=1}^{n}\log\left[\frac{\frac{\partial}{\partial\alpha}(1-\frac{1}{\alpha^{\beta}}\{(1+\alpha^{\beta})(1-e^{-\lambda y^{\beta}})+e^{-\lambda y^{\beta}}(1+\alpha^{\beta})-1)}{\frac{1}{\alpha^{\beta}}\{(1+\alpha^{\beta})(1-e^{-\lambda y^{\beta}})+e^{-\lambda y^{\beta}}(1+\alpha^{\beta})-1}\right]$$
(31)

The equations derive above can also be solved using iteration method (Newton Raphson) to obtain the  $\hat{a}, \hat{\lambda}, \hat{\beta}, \hat{\alpha}$  the MLE of  $(a, \lambda, \beta, \alpha)$  respectively.

Taking second derivatives of the said equations 28, 29, 30 and 31 with respect to the parameters above, it is possible to derive the interval estimate and hypothesis tests on the model parameter. This may be shown in further research.

#### V. REAL DATA SET

The data used in this section was studied by Lemonte in his BJPS Accepted Manuscript on the remission times (in months) of a random sample of 128 bladder cancer patients reported in Lee and Wang (2003) to compare and contrast the Exponentiated Weighted Weibull and Weighted Weibull distribution.

R (code) software is used to determine the maximum likelihood estimates and the log-likelihood for the Exponentiated Weighted Weibull distribution are:  $\hat{a} =$ 8.95941,  $\hat{\lambda} = 8.17089$ ,  $\hat{\beta} = 8.36111$ ,  $\hat{\alpha} = 7.91621$  and  $log_{EWW} = 448.5781$  while the maximum likelihood estimates and the log-likelihood for the Weighted Weibull distribution are:  $\hat{\lambda} = 7.90434$ ,  $\hat{\beta} = 8.44282$ ,  $\hat{\alpha} = 8.70033$  and  $log_{WW} = 448.4913$ , where  $lg_{LWW}$  and  $lg_{WW}$  denote log-likelihood of both Exponentiated Weighted Weibull distribution and Weighted Weibull distribution.

	/ 0.802169648	-0.02880174	0.002221849	-0.02674447
l	-0.028801742	0.65433442	-0.025345533	-0.02069611
۱	0.002221849	-0.02534553	0.705642669	-0.02353514
	\−0.026744474	-0.02069611	-0.023535137	0.60907454 /

#### VI. CONCLUSION

We investigated on the statistical properties of the proposed distribution e.g moments, moment generating function, estimation of parameters using R (Code) software for data analysis presented in this article. We also upgraded with an additional parameter to the existing three parameters in the weighted weibull distribution and the results from the estimated parameters show that the Exponentiated Weighted Weibull distribution has a better representation of data than weighted weibull distribution.

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# A Summation Formula Involving Certain Special Functions

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*Abstract-* The main aim of the present paper is to compute a summation formula involving Recurrence relation and Contiguous relation.

Keywords: gauss second summation theorem, recurrence relation, prudnikov.

GJSFR-F Classification : MSC 2010: 33C05 , 33C20 , 33C45 ,33D50 ,33D60

# ASUMMATIONFORMULAINVOLVINGCERTAINSPECIALFUNCTIONS

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Notes

# A Summation Formula Involving Certain Special Functions

Salahuddin  $^{\alpha}$  & Intazar Husain  $^{\sigma}$ 

Abstract- The main aim of the present paper is to compute a summation formula involving Recurrence relation and Contiguous relation.

Keywords : gauss second summation theorem, recurrence relation, prudnikov.

#### I. INTRODUCTION

Generalized Gaussian Hypergeometric function of one variable is defined by

$${}_{A}F_{B}\left[\begin{array}{ccc}a_{1},a_{2},\cdots,a_{A} & ;\\ & & \\ b_{1},b_{2},\cdots,b_{B} & ;\end{array}\right] = \sum_{k=0}^{\infty} \frac{(a_{1})_{k}(a_{2})_{k}\cdots(a_{A})_{k}z^{k}}{(b_{1})_{k}(b_{2})_{k}\cdots(b_{B})_{k}k!}$$
(1)

where the parameters  $b_1, b_2, \dots, b_B$  are neither zero nor negative integers and A, B are non-negative integers and |z| = 1

#### Contiguous Relation is defined by

 $[ \mbox{ Andrews p.363(9.16), E. D. p.51(10)} ]$ 

$$(a-b) {}_{2}F_{1} \begin{bmatrix} a, b; \\ c; \end{bmatrix} = a {}_{2}F_{1} \begin{bmatrix} a+1, b; \\ c; \end{bmatrix} - b {}_{2}F_{1} \begin{bmatrix} a, b+1; \\ c; \end{bmatrix}$$
(2)

Gauss second summation theorem is defined by [Prudnikov., 491(7.3.7.5)]

$${}_{2}F_{1}\left[\begin{array}{cc}a, b \ ; & 1\\\frac{a+b+1}{2} \ ; & 2\end{array}\right] = \frac{\Gamma(\frac{a+b+1}{2}) \ \Gamma(\frac{1}{2})}{\Gamma(\frac{a+1}{2}) \ \Gamma(\frac{b+1}{2})}$$
(3)

$$=\frac{2^{(b-1)} \Gamma(\frac{b}{2}) \Gamma(\frac{a+b+1}{2})}{\Gamma(b) \Gamma(\frac{a+1}{2})}$$
(4)

In a monograph of Prudnikov et al., a summation theorem is given in the form [Prudnikov.,  $\mathrm{p.491}(7.3.7.8)]$ 

$${}_{2}F_{1}\left[\begin{array}{cc}a, \ b \\ \frac{a+b-1}{2} \ ; & \frac{1}{2}\right] = \sqrt{\pi} \left[\frac{\Gamma(\frac{a+b+1}{2})}{\Gamma(\frac{a+1}{2}) \ \Gamma(\frac{b+1}{2})} + \frac{2 \ \Gamma(\frac{a+b-1}{2})}{\Gamma(a) \ \Gamma(b)}\right]$$
(5)

Author α: P. D. M College of Engineering, Bahadurgarh, Haryana, India. e-mail: vsludn@gmail.com Author α: Department of Applied Sciences & Humanities, Jamia Millia Islamia, New Delhi, India. Now using Legendre's duplication formula and Recurrence relation for Gamma function, the above theorem can be written in the form

$${}_{2}F_{1}\left[\begin{array}{c}a,\ b\\\frac{a+b-1}{2}\ ; \\ \end{array}\right] = \frac{2^{(b-1)}\ \Gamma(\frac{a+b-1}{2})}{\Gamma(b)}\left[\frac{\Gamma(\frac{b}{2})}{\Gamma(\frac{a-1}{2})} + \frac{2^{(a-b+1)}\ \Gamma(\frac{a}{2})\ \Gamma(\frac{a+1}{2})}{\{\Gamma(a)\}^{2}} + \frac{\Gamma(\frac{b+2}{2})}{\Gamma(\frac{a+1}{2})}\right]$$
(6)

Recurrence relation is defined by

$$\Gamma(\zeta + 1) = \zeta \ \Gamma(\zeta) \tag{7}$$

Notes

II. MAIN SUMMATION FORMULA

$${}_{2}F_{1}\left[\begin{array}{cc}a, \ b \ ; & \frac{1}{2}\\\frac{a+b+45}{2} \ ; & \frac{1}{2}\right] = \\ = \frac{2^{b} \ \Gamma(\frac{a+b+45}{2})}{(a-b) \ \Gamma(b)} \left[\frac{\Gamma(\frac{b}{2})}{\Gamma(\frac{a+1}{2})} \left\{\frac{1}{\left[\prod\limits_{\zeta=1}^{21} \left\{a-b-(2\zeta-1)\right\}\right] \left[\prod\limits_{\eta=1}^{22} \left\{a-b+(2\eta-1)\right\}\right]} \times \right.$$

 $-33100902479081461753011150a^2 + 18852771841456100048901630a^3 \\$ 

 $-6998152417792503243516261a^4 + 1830981648088050432124941a^5 \\$ 

 $-117140064945846306a^{12} + 4552422878866226a^{13} - 142380641710440a^{14} + 3551773672360a^{15} + 142380641710440a^{15} + 142380641710440a^{14} + 142380641710440a^{15} + 14238064171040a^{15} + 14238064171040a^{15} + 14238064171040a^{15} + 14238064171040a^{15} + 14238064171040a^{15} + 14238064171040a^{15} + 142380641710440a^{15} + 14238064171040a^{15} + 142380641700a^{15} + 142380641700a^{15} + 142380641700a^{15} + 142380641700a^{15} + 14238064180a^{15} + 142380641700a^{15}$ 

 $-69638530941a^{16} + 1048818981a^{17} - 11702670a^{18} + 91070a^{19} - 441a^{20} + a^{21}$ 

$$+53745962882620608893167875b-45053743668222224004344100ab$$

 $+ 22213537883171068151439531a^4b - 3437881821120782582829456a^5b \\$ 

 $+ 1069513403791665845357128a^{6}b - 102017962030642073210480a^{7}b$ 

 $-4742709531308534760a^{11}b + 317038686494337262a^{12}b - 8612153544899408a^{13}b - 861215354489408a^{13}b - 861215354489408a^{13}b - 861215354486b - 861215354486b - 8612153546b - 861215566b - 861215566b - 861215566b - 86121556b - 86121556b - 86121556b - 86121556b - 86121556b - 86121556b - 861256b - 861$ 

 $+ 370669544999304a^{14}b - 6357587526960a^{15}b + 172733256855a^{16}b - 1689533124a^{17}b + 27096622a^{18}b + 1866b + 186b + 1866b + 1866b + 186b +$ 

 $-29662057538012999895338574a^2b^2+73924679489058918046015938a^3b^2$ 

 $-9095261732430392323657176a^4b^2+6475045382105375335889480a^5b^2$ 

 $-523071699408050859929016a^{6}b^{2}+171242942351303632910536a^{7}b^{2}$
$-9336842670271654008484a^8b^2 + 1689474770765990596076a^9b^2 - 62588562419040481252a^{10}b^2 - 625885624190b^2 - 625885624190b^2 - 625885624190b^2 - 625885624190b^2 - 625885624190b^2 - 625885624190b^2 - 62588562419b^2 - 625885624190b^2 - 62588562419b^2 - 625885662419b^2 - 62588566240 - 6258856624 - 6258856626 - 6258856626 - 6258856626 - 625885662 - 625885662 - 625885662 - 625885662 - 625885662 - 625885662 - 625885662 - 62588566 - 625866 - 62588566 - 62588566 - 625866 - 62588566 - 6258666 - 625866 - 62588566 - 62586$  $+ 6843044908448608444a^{11}b^2 - 169609751776172344a^{12}b^2 + 11684556963374568a^{13}b^2$  $-185095028637272a^{14}b^2 + 8105951624872a^{15}b^2 - 73945230446a^{16}b^2 + 2006403306a^{17}b^2 - 66666a^{16}b^2 + 1006666a^{16}b^2 + 1006666a^{16}b^2 + 100666a^{16}b^2 + 100666b^2 + 100666a^{16}b^2 + 10066a^{16}b^2 + 10066a^{16}b^2 + 10066a^{16}b^2 + 10066a^{16}b^2 + 100666a^{16}b^2 + 10066a^{16}b^2 + 100666a^{16}b^2 + 10066a^{16}b^2 + 10066a$  $+91735804724807821121609526a^{2}b^{3} - 6747725269089701729734368a^{3}b^{3}$  $+ 16438427349229328982545896a^4b^3 - 1038876114726270005445776a^5b^3$  $+735240415476876038312632a^6b^3-34241820697597088722848a^7b^3$  $+ 11226186945118492939852a^8b^3 - 371078402291905961224a^9b^3 + 67240314521518855860a^{10}b^3 + 6724056b^3 + 672405b^3 + 67240b^3 + 6724b^3 + 674b^3 +$  $+ 167607787407128a^{14}b^3 - 1445744454752a^{15}b^3 + 61843300298a^{16}b^3 - 235148940a^{17}b^3 - 2351480a^{17}b^3 - 235148b^{17}b^3 - 23514b^{17}b^3 - 23514b$  $+ 1858191737565276546672360a^2b^4 + 19566549200385431103972216a^3b^4 \\$  $-673581690360746582755756a^4b^4 + 1596002262124887064872892a^5b^4 \\$  $+ 352022359346404662630a^9b^4 - 7157179769538448936a^{10}b^4 + 1266877992724507592a^{11}b^4 + 1266877992a^{11}b^4 + 1266877992a^{11}b^4 + 1266876b^{11}b^4 + 1266876b^{11}b^{$  $-3560764597a^{16}b^4 + 145008513a^{17}b^4 + 6412982275781442369534447b^5 \\$  $+ 3020391903873160874624880ab^5 + 11072544506717425040943320a^2b^5$  $+251409502832929940580848a^{3}b^{5}+1850900598832213039588308a^{4}b^{5}$  $-33362183528683696143184a^{5}b^{5} + 76635144078132798484872a^{6}b^{5} - 1614107103242780383952a^{7}b^{5} + 7663514a^{5}b^{5} + 7663514b^{5} + 7663514a^{5}b^{5} + 7663514b^{5} + 7663514a^{5}b^{5} + 7663514b^{5} + 766354b^{5} + 7663514b^{5} + 766354b^{5} + 7664554b^{5} + 7664554b^{5} + 7664554b^{5} + 766454b^{5} + 7664554b^{5} +$  $-68669517587706672a^{11}b^5 + 11692563911936180a^{12}b^5 - 87796234586480a^{13}b^5 + 8925930851320a^{14}b^5 + 89259560a^{14}b^5 + 89259560a^{14}b^5 + 892560a^{14}b^5 + 892560a^{14}b^{14}b^{14}b^5 + 892560a^{14}b^{14}$  $-30677356528a^{15}b^5 + 1917334783a^{16}b^5 + 604573745934401191005912b^6$  $+2761535361221837965136664ab^{6}+446409836411046907751368a^{2}b^{6}$  $+1155941471891438860601896a^{3}b^{6} + 15624214876016679739352a^{4}b^{6}$  $-24520882366446890424a^{8}b^{6}+15396021766853800072a^{9}b^{6}-157355447527654376a^{10}b^{6}$  $+15338678264a^{15}b^6 + 231724339078175706150136b^7 + 153811376527605888117264ab^7$  $+374921775760059309890200a^{2}b^{7}+28831347254232216398432a^{3}b^{7}$  $+58118545317700335413400a^{4}b^{7}+490468210861184892336a^{5}b^{7}+2185851421255096595000a^{6}b^{7}+2185656b^{7}+218566b^{7}+21856b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2186b^{7}+2$ 

2014

Year

 $+ 13039358045721374943420a^2b^8 + 22139104917821671572676a^3b^8 + 921471943096864139266a^4b^8 + 9214719430966a^4b^8 + 9214719406a^4b^8 + 92146a^4b^8 + 9214b^8 + 920b^8 + 920b^$  $+434411052282116a^{11}b^8 - 1013931747010a^{12}b^8 + 265182149218a^{13}b^8 + 3024008960876320420246b^9 + 265182149218a^{13}b^8 + 26518214b^8 + 26518214b^8 + 26518214b^8 + 26518214b^8 + 26518214b^8 + 26518214b^8 + 2651824b^8 + 265182b^8 + 265182b^8 + 26518b^8 + 265$  $+ 69758491996612848a^7b^9 + 215842089833789562a^8b^9 - 353214178214040a^9b^9 + 21584208983789562a^8b^9 - 353214178214040a^9b^9 + 21584208983789562a^8b^9 - 353214178214040a^9b^9 + 21584208983789562a^8b^9 + 21584208983789562a^8b^9 - 353214178214040a^9b^9 + 21584208983789562a^8b^9 - 353214178214040a^9b^9 + 2158420898562a^8b^9 - 353214178214040a^9b^9 + 215842089866a^8 + 21584208966a^8 + 21584208966a^8 + 2158666a^8 + 2158666a^8 + 2158666a^8 + 215866a^8 + 2158666a^8 + 2158666a^8 + 215866a^8 + 2158666a^8 + 215866a^8 + 215$  $+702101231167908a^{10}b^9 - 1152680301864a^{11}b^9 + 608359048206a^{12}b^9 + 137104871154095942604b^{10} + 108466a^{11}b^9 + 108466a^{11$  $+ 9180121926490903960a^{4}b^{10} + 9807216639522886584a^{5}b^{10} + 132405723186144408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 132405723186144408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 132405723186144408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 132405728a^{6}b^{10} + 13240572318614408a^{6}b^{10} + 1324057231864b^{10} + 1324057231864b^{10} + 1324057231864b^{10} + 1324057231864b^{10} + 132405724b^{10} + 1324056b^{10} + 1324056b^{10} + 1324056b^{10} + 1324056b^{10} + 1324056b^{10} + 1324056b^{10} + 132405b^{10} + 1324056b^{10} + 13240b^{10} + 13240b^{10} + 13240b^{10} + 13240b^{10} + 13240b^{10} + 13240b^{10} + 13240$  $-503154100020a^{10}b^{10}+960566918220a^{11}b^{10}+16116280063842269532b^{11}$  $+ 13266524009934170456ab^{11} + 22468630494292536788a^2b^{11} + 2744654110430536928a^3b^{11} + 274465411b^{11} + 27446541b^{11} + 27446541b^{11} + 27446541b^{11} + 2744b^{11} + 274b^{11} + 27$  $+ 552383916933472a^7b^{11} + 579276741224844a^8b^{11} + 457412818200a^9b^{11} + 1052049481860a^{10}b^{11} + 1052049484a^{10}b^{11} + 1052049484a^{10}b^{11} + 1052049484a^{10}b^{11} + 1052049484a^{10}b^{11} + 1052049484a^{11} + 10520494b^{11} + 10520494b^{11} + 105204b^{11} + 10$  $+483431815242444952a^{3}b^{12}+29546096072985492a^{4}b^{12}+24458385708121820a^{5}b^{12}$  $+ 381004267695688a^{6}b^{12} + 294616546783832a^{7}b^{12} + 864510226398a^{8}b^{12} + 800472431850a^{9}b^{12} + 8004766a^{12} + 8004766a^{12} + 800476a^{12} + 800476a^{12} + 800476a^{12} + 800476a^{$  $+ 36842919196739110b^{13} + 30556991340393392ab^{13} + 46327319945209016a^2b^{13} + 66327319945209016a^2b^{13} + 66327319b^{13} + 6632756b^{13} + 66$  $+5648934122340208a^{3}b^{13}+4881179381635732a^{4}b^{13}+144878605202576a^{5}b^{13}$  $+ 22313057280808a^{5}b^{14} + 256037937176a^{6}b^{14} + 151532656696a^{7}b^{14} + 35857948972408b^{15} + 3585794886b^{15} + 358579486b^{15} + 358579486b^{15} + 358579486b^{15} + 358579486b^{15} + 358579486b^{15} + 358579486b^{15} + 35857948b^{15} + 35857948b^{15} + 35857948b^{15} + 35857948b^{15} + 35857948b^{15} + 35857948b^{15} + 3585786b^{15} + 3585786b^{15} + 3585786b^{15} + 3585786b^{15} + 3585786b^{15} + 358578b^{15} + 35858b^{15} + 3$  $+28280942023504ab^{15}+39037774410808a^{2}b^{15}+4143128700320a^{3}b^{15}+3132630060840a^{4}+3132630060840a^{4}+316b^{15}+316b^{15}+316b^{15}+315b^{15}+316b^{15}+315b^{1$  $+ 63714509712a^{5}b^{15} + 36576848168a^{6}b^{15} + 484185248451b^{16} + 1250881854197ab^{16} + 12508818548485484866b^{16} + 12508881866b^{16} + 12508881866b^{16} + 1250888185486b^{16} + 12508881856b^{16} + 12508881856b^{16} + 12508881856b^{16} + 12508886b^{16} + 12508886b^{16} + 1250886b^{16} + 1250886b^{16} + 1250886b^{16} + 1250886b^{16} + 1250886b^{16} + 125086b^{16} + 125086b^{16} + 125086b^{16} + 125086b^{16} + 125086b^{16} + 12508$  $+9489432124ab^{17}+11986499486a^{2}b^{17}+837826964a^{3}b^{17}+563921995a^{4}b^{17}+106564578b^{18}$  $+254376218ab^{18}+39191490a^2b^{18}+32224114a^3b^{18}+1743994b^{19}+839188ab^{19}+962598a^2b^{19}+1743994b^{19}+839188ab^{19}+962598a^2b^{19}+1743994b^{19}+839188ab^{19}+962598a^2b^{19}+1743994b^{19}+19444b^{19}+1944b^{19}+1944b^{19}+1944b^{19}+1744b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19}+194b^{19$ 

$$+5719b^{20} + 12341ab^{20} + 43b^{21}) + \frac{1}{\left[\prod_{\mu=1}^{22} \left\{a - b - (2\mu - 1)\right\}\right] \left[\prod_{\xi=1}^{21} \left\{a - b + (2\xi - 1)\right\}\right]} \times$$

### $\mathbf{N}_{\mathrm{otes}}$

$\Big(2097152b(-13113070457687988603440625+53745962882620608893167875a)\Big)$	
$+3418740929846540363004450a^2 + 49302844553810665888575930a^3$	
$+ 6523406582853698567485227a^4 + 6412982275781442369534447a^5 + 604573745934401191005912a^3 + 604573745984605944059546056666666666666666666666666$	$a^6$
$+ 231724339078175706150136a^7 + 15128932813380875312110a^8 + 3024008960876320420246a^9 + 15128932813380875312110a^8 + 15128932813380875312110a^8 + 15128932813380875312110a^8 + 15128932813380875312110a^8 + 15128932813380875312110a^8 + 15128932813380875312110a^8 + 1512893281338087531210a^8 + 15128932813380875312804008960876320420246a^9 + 1512893281338087531280408960876320420246a^9 + 1512893281338087531280408960876320420246a^9 + 1512893281338087531280408960876320420246a^9 + 15128932881380875312804089608763204200896087632042008960876320400896087632040089608763204008960876320400896086666666666666666666666666666666$	
$+ 137104871154095942604a^{10} + 16116280063842269532a^{11} + 504464808447965118a^{12} + 5046866666666666666666666666666666666666$	
$+ 36842919196739110a^{13} + 777692306266456a^{14} + 35857948972408a^{15} + 484185248451a^{16} + 100000000000000000000000000000000000$	
$+ 13811946279a^{17} + 106564578a^{18} + 1743994a^{19} + 5719a^{20} + 43a^{21} + 32835960506324395192397625b^{10} + 5719a^{10} + 5719a$	5
$-45053743668222224004344100ab + 152628207170808204647275770a^2b$	
$+5196235439832751151572788a^{3}b + 43689757687398414790547553a^{4}b$	
$+ 3020391903873160874624880a^5b + 2761535361221837965136664a^6b$	
$+ 153811376527605888117264a^7b + 56881481026993911310578a^8b + 2332408473401733698504a^9b + 233240847340047340b + 23324086b + 23324086b + 23324086b + 23324086b + 23324086b + 233266b + 2386b + 23326b + 23326b + 23326b + 23326b + 2386b +$	)
$+454222013907202843900a^{10}b+13266524009934170456a^{11}b+1520074513892383306a^{12}b$	
$+ 30556991340393392a^{13}b + 2164283489573848a^{14}b + 28280942023504a^{15}b + 1250881854197a^{16}b^{16}b + 1250881854197a^{16}b^{$	)
$+9489432124a^{17}b + 254376218a^{18}b + 839188a^{19}b + 12341a^{20}b - 33100902479081461753011150b^{20}b + 12341a^{20}b + 12341a^{$	2
$+ 113826458136107026075141710ab^2 - 29662057538012999895338574a^2b^2 - 29662057538012995600000000000000000000000000000000000$	
$+91735804724807821121609526a^{3}b^{2}+1858191737565276546672360a^{4}b^{2}$	
$+ 11072544506717425040943320a^{5}b^{2} + 446409836411046907751368a^{6}b^{2}$	
$+ 374921775760059309890200a^7b^2 + 13039358045721374943420a^8b^2 + 4569477008635213654148a^9b^2 + 13039358045721374943420a^8b^2 + 1303935804572137494440a^8b^2 + 130393580a^8 + 1303956a^8 + 1303956a^8 + 1303956a^8 + 130396a^8 + 13036a^8 + 13036a^8 + 13036a^8 + 1306a^8 + $	$b^2$
$+ 120615054891875887772a^{10}b^2 + 22468630494292536788a^{11}b^2 + 422390481555014024a^{12}b^2 + 22468630494292536788a^{11}b^2 + 22468630494292536788a^{11}b^2 + 22468630494292536788a^{11}b^2 + 22468630494292536788a^{11}b^2 + 22468630494292536788a^{11}b^2 + 22468630494292536788a^{11}b^2 + 224686304942925566786a^{11}b^2 + 224686304942925566666666666666666666666666666666$	
$+46327319945209016a^{13}b^2+578915870956584a^{14}b^2+39037774410808a^{15}b^2+288251667106a^{16}b^2+28825665666666666666666666666666666666$	2
$+ 11986499486a^{17}b^2 + 39191490a^{18}b^2 + 962598a^{19}b^2 + 18852771841456100048901630b^3 + 188527718414561000489006565656565656565656565656565656565656$	
$-32581896347358559999317420ab^3+73924679489058918046015938a^2b^3$	
$-6747725269089701729734368a^3b^3 + 19566549200385431103972216a^4b^3 \\$	
$+ 251409502832929940580848a^5b^3 + 1155941471891438860601896a^6b^3 \\$	
$+ 28831347254232216398432a^7b^3 + 22139104917821671572676a^8b^3 + 492678254183696354488a^9b^3 + 492678254183696354488a^9b^3 + 22139104917821671572676a^8b^3 + 2213910491782167566666666666666666666666666666666666$	3
$+ 161949836634482706524a^{10}b^3 + 2744654110430536928a^{11}b^3 + 483431815242444952a^{12}b^3 + 2744654110430536928a^{11}b^3 + 483431815242444952a^{12}b^3 + 283431815242444952a^{12}b^3 + 283431845444952a^{12}b^3 + 2834318454444952a^{12}b^3 + 28343184b^3 + 28343184b^3 + 28343184b^3 + 28344444b^3 + 283444b^3 + 28344b^3 + 2834b^3 + 2834b^3 + 2834b^3 + 2834b^3 + 284b^3 + 284$	
$+5648934122340208a^{13}b^3 + 585771144020552a^{14}b^3 + 4143128700320a^{15}b^3 + 262400960302a^{16}b^3 + 262400000000000000000000000000000000000$	
$+837826964a^{17}b^3+32224114a^{18}b^3-6998152417792503243516261b^4$	
$+ 22213537883171068151439531ab^4 - 9095261732430392323657176a^2b^4 \\$	
$+ 16438427349229328982545896a^{3}b^{4} - 673581690360746582755756a^{4}b^{4} - 673586a^{4}b^{4} - 673586a^{4} - 673586a^{4} - 673586a^{4}b^{4} - 673586a^{4} - 67358$	
$+ 1850900598832213039588308a^5b^4 + 15624214876016679739352a^6b^4 \\$	
$+58118545317700335413400a^7b^4 + 921471943096864139266a^8b^4 + 645945537218444676642a^9b^4 + 64594553721844676642a^9b^4 + 64594553721864553721844676642a^9b^4 + 64594553721864565666666666666666666666666666666666$	

### $N_{\rm otes}$

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 $+9180121926490903960a^{10}b^4 + 2812829209563463896a^{11}b^4 + 29546096072985492a^{12}b^4 \\ +4881179381635732a^{13}b^4 + 32252520795880a^{14}b^4 + 3132630060840a^{15}b^4 + 9586673915a^{16}b^4 \\ +563921995a^{17}b^4 + 1830981648088050432124941b^5 - 3437881821120782582829456ab^5 \\ +6475045382105375335889480a^2b^5 - 1038876114726270005445776a^3b^5$ 

 $+ 1596002262124887064872892a^4b^5 - 33362183528683696143184a^5b^5 \\$ 

### Notes

 $+87313680148565150111832a^{6}b^{5} + 490468210861184892336a^{7}b^{5} + 1520878767836758616286a^{8}b^{5} \\ +15329149883995675728a^{9}b^{5} + 9807216639522886584a^{10}b^{5} + 85999358210064208a^{11}b^{5} \\ +24458385708121820a^{12}b^{5} + 144878605202576a^{13}b^{5} + 22313057280808a^{14}b^{5} + 63714509712a^{15}b^{5} \\ +5752004349a^{16}b^{5} - 354501246907809948405480b^{6} + 1069513403791665845357128ab^{6} \\ -523071699408050859929016a^{2}b^{6} + 735240415476876038312632a^{3}b^{6} \\ +266864612882774626831720a^{4}b^{6} + 76635144078132798484872a^{5}b^{6} - 874649607643712372472a^{6}b^{6} \\ +2185851421255096595000a^{7}b^{6} + 8095043487442343432a^{8}b^{6} + 21357540300644860312a^{9}b^{6} \\ +132405723186144408a^{10}b^{6} + 77331167721172456a^{11}b^{6} + 381004267695688a^{12}b^{6} \\ +100261443987096a^{13}b^{6} + 256037937176a^{14}b^{6} + 36576848168a^{15}b^{6} + 52462262753503539155240b^{7} \\ -102017962030642073210480ab^{7} + 171242942351303632910536a^{2}b^{7} \\ -34241820697597088722848a^{3}b^{7} + 39096473046294723020936a^{4}b^{7} - 1614107103242780383952a^{5}b^{7} \\ +1943115716512939776808a^{6}b^{7} - 12438236313336433600a^{7}b^{7} + 29720901376476436984a^{8}b^{7} \\ +69758491996612848a^{9}b^{7} + 158910354362756568a^{10}b^{7} + 552383916933472a^{11}b^{7} \\ +294616546783832a^{12}b^{7} + 623957998160a^{13}b^{7} + 151532656696a^{14}b^{7} - 6068045552614043865426b^{8} \\$ 

 $+17605238308074264053574ab^{8} - 9336842670271654008484a^{2}b^{8} + 11226186945118492939852a^{3}b^{8} \\ -1099649665242617000798a^{4}b^{8} + 1064600574662833613834a^{5}b^{8} - 24520882366446890424a^{6}b^{8} \\ +26672527585599011688a^{7}b^{8} - 94427080896983406a^{8}b^{8} + 215842089833789562a^{9}b^{8}$ 

 $+291326223911580a^{10}b^{8} + 579276741224844a^{11}b^{8} + 864510226398a^{12}b^{8} + 421171648758a^{13}b^{8} + 557078686778195952226b^{9} - 1094032303287108057912ab^{9} + 1689474770765990596076a^{2}b^{9} - 371078402291905961224a^{3}b^{9} + 352022359346404662630a^{4}b^{9} - 18285405023328014000a^{5}b^{9}$ 

 $+15396021766853800072a^{6}b^{9} - 196990298922392976a^{7}b^{9} + 195113449113625614a^{8}b^{9} \\ -353214178214040a^{9}b^{9} + 772341543530700a^{10}b^{9} + 457412818200a^{11}b^{9} + 800472431850a^{12}b^{9} \\ -41014337504927410260b^{10} + 114934366191849037236ab^{10} - 62588562419040481252a^{2}b^{10} \\ +67240314521518855860a^{3}b^{10} - 7157179769538448936a^{4}b^{10} + 5631584078513951272a^{5}b^{10} \\ -157355447527654376a^{6}b^{10} + 117137843987075912a^{7}b^{10} - 774287072717444a^{8}b^{10} \\ +702101231167908a^{9}b^{10} - 503154100020a^{1}0b^{10} + 1052049481860a^{11}b^{10} + 2436777594647459700b^{11} \\ \end{array}$ 

 $-4742709531308534760ab^{11}+6843044908448608444a^2b^{11}-1526743180734652192a^3b^{11}-152674318b^{11}-152674318b^{11}-156674318b^{11}-15667431b^{11}-156674318b^{11}-15667431b^{11}-15667431b^{11}-15667431b^{11}-15667431b^{11}-15667431b^{11}-15667431b^{11}-15667431b^{11}-15667431b^{11}-15667445b^{11}-15667445b^{11}-15667445b^{11}-15667445b^{11}-15667445b^{11}-15667445b^{11}-15667445b^{11}-15667445b^{11}-1566745b^{11}-15665b^{11}-1566b^{11}-1566b^{11}-1566b^{11}-1566b^{11}-1566b^{11}-156b^{11}$ 

 
$$\begin{split} &+1266877992724507592a^4b^{11}-68669517587706672a^5b^{11}+46179770718580408a^6b^{11}\\ &-652631982303392a^7b^{11}+434411052282116a^8b^{11}-1152680301864a^9b^{11}+960566918220a^{10}b^{11}\\ &-117140064945846306b^{12}+317038686494337262ab^{12}-169609751776172344a^2b^{12}\\ &+166445736776215624a^3b^{12}-17271058034689004a^4b^{12}+11692563911936180a^5b^{12}\\ &-309271965286648a^6b^{12}+181277011099912a^7b^{12}-1013931747010a^8b^{12}+608359048206a^9b^{12}\\ &+4552422878866226b^{13}-8612153544899408ab^{13}+11684556963374568a^2b^{13}\\ &-2459287660143888a^3b^{13}+1832663509135228a^4b^{13}-87796234586480a^5b^{13}+50125620668424a^6b^{13}\\ &-512075874352a^7b^{13}+265182149218a^8b^{13}-142380641710440b^{14}+370669544999304ab^{14}\\ &-185095028637272a^2b^{14}+167607787407128a^3b^{14}-14882731838680a^4b^{14}+8925930851320a^5b^{14}\\ &-159286274280a^6b^{14}+78378960360a^7b^{14}+3551773672360b^{15}-6357587526960ab^{15}\\ &+8105951624872a^2b^{15}-1445744454752a^3b^{15}+979484169144a^4b^{15}-30677356528a^5b^{15}\\ &+15338678264a^6b^{15}-69638530941b^{16}+172733256855ab^{16}-73945230446a^2b^{16}+61843300298a^3b^{16}\\ &-3560764597a^4b^{16}+1917334783a^5b^{16}+1048818981b^{17}-1689533124ab^{17}+2006403306a^2b^{17}\\ &-235148940a^3b^{17}+145008513a^4b^{17}-11702670b^{18}+27096622ab^{18}-7972286a^2b^{18}+6096454a^3b^{18}\\ \end{split}$$

 $\left. +91070b^{19} - 114380ab^{19} + 123410a^2b^{19} - 441b^{20} + 903ab^{20} + b^{21}) \right) \right\} -$ 

$$-\frac{\Gamma(\frac{b+1}{2})}{\Gamma(\frac{a}{2})} \Biggl\{ \frac{4194304}{\left[\prod\limits_{\zeta=1}^{21} \left\{a-b-(2\zeta-1)\right\}\right] \left[\prod\limits_{\eta=1}^{22} \left\{a-b+(2\eta-1)\right\}\right]} (13113070457687988603440625)$$

 $+53745962882620608893167875a-3418740929846540363004450a^2$ 

 $+49302844553810665888575930a^3-6523406582853698567485227a^4$ 

 $+ 6412982275781442369534447a^5 - 604573745934401191005912a^6 + 231724339078175706150136a^7 - 15128932813380875312110a^8 + 3024008960876320420246a^9 - 137104871154095942604a^{10}$ 

 $+16116280063842269532a^{11} - 504464808447965118a^{12} + 36842919196739110a^{13} \\ -777692306266456a^{14} + 35857948972408a^{15} - 484185248451a^{16} + 13811946279a^{17} - 106564578a^{18} \\ +1743994a^{19} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{19} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{19} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{10} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{10} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{10} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{10} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{10} - 5719a^{20} + 43a^{21} + 32835960506324395192397625b + 45053743668222224004344100ab^{10} \\ +1743994a^{10} - 5719a^{10} + 5719a^$ 

 $+2761535361221837965136664a^{6}b - 153811376527605888117264a^{7}b$ 

 $+56881481026993911310578a^{8}b - 2332408473401733698504a^{9}b + 454222013907202843900a^{10}b \\ +2164283489573848a^{14}b - 28280942023504a^{15}b + 1250881854197a^{16}b - 9489432124a^{17}b \\ +254376218a^{18}b - 839188a^{19}b + 12341a^{20}b + 33100902479081461753011150b^{2}$ 

 $+ 113826458136107026075141710ab^2 + 29662057538012999895338574a^2b^2$  $+91735804724807821121609526a^{3}b^{2} - 1858191737565276546672360a^{4}b^{2}$  $+ 11072544506717425040943320a^{5}b^{2} - 446409836411046907751368a^{6}b^{2} - 44640866b^{2} - 4464008b^{2} - 44$  $+ 374921775760059309890200a^7b^2 - 13039358045721374943420a^8b^2 \\$ Notes  $+4569477008635213654148a^9b^2-120615054891875887772a^{10}b^2+22468630494292536788a^{11}b^2$  $-422390481555014024a^{12}b^2+46327319945209016a^{13}b^2-578915870956584a^{14}b^2$  $+ 39037774410808a^{15}b^2 - 288251667106a^{16}b^2 + 11986499486a^{17}b^2 - 39191490a^{18}b^2 + 962598a^{19}b^2 - 966566a^{19}b^2 - 96656b^{19}b^2 - 9665b^{19}b^2 - 965b^{19}b^2 - 96b^{19}b^2 - 96b^{19}b^2 - 96b^{19}b^{19}b^2 - 96b^{$  $+ 18852771841456100048901630b^3 + 32581896347358559999317420ab^3 \\$  $+73924679489058918046015938a^{2}b^{3}+6747725269089701729734368a^{3}b^{3}$  $+ 19566549200385431103972216a^4b^3 - 251409502832929940580848a^5b^3 \\$  $+ 1155941471891438860601896a^6b^3 - 28831347254232216398432a^7b^3 \\$  $+ 22139104917821671572676a^8b^3 - 492678254183696354488a^9b^3 + 161949836634482706524a^{10}b^3 + 1619498b^3 + 161949856634482706524a^{10}b^3 + 16194b^3 + 1614b^3 + 1614b^3 + 1614b^3 + 161b^3 + 160b^3 + 160b^3 + 16b^3 +$  $-2744654110430536928a^{11}b^3 + 483431815242444952a^{12}b^3 - 5648934122340208a^{13}b^3 - 564894b^3 - 564894b^3 - 56489b^3 - 56489b^3 - 56489b^3 - 5648b^3 - 564b^3 - 5648b^3 - 5648b^3 - 564b^3 - 564b^3$  $+585771144020552a^{14}b^3-4143128700320a^{15}b^3+262400960302a^{16}b^3-837826964a^{17}b^3$  $+ 32224114a^{18}b^3 + 6998152417792503243516261b^4 + 22213537883171068151439531ab^4 + 22213537883164 + 222135378831864 + 2221353788831710681514844 + 22213537888318844 + 2221353788844 + 2221353788844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 222135378844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213578844 + 22213844 + 22213844 + 22213844 + 22213844 + 22213844 + 22213844 + 22213844 + 22213844 + 22213844 + 22213844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 222135844 + 2221844 + 222135844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 2221844 + 22218444 + 2284844 + 22218444 + 2221844 + 2221844 + 22218444 + 22218$  $+9095261732430392323657176a^2b^4+16438427349229328982545896a^3b^4$  $+673581690360746582755756a^4b^4 + 1850900598832213039588308a^5b^4$  $-15624214876016679739352a^{6}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}-921471943096864139266a^{8}b^{4}+58118545317700335413400a^{7}b^{4}+58118545317700335413400a^{7}b^{4}+5811854531700335413400a^{7}b^{4}+581185453170035414b^{4}+5811854553170035414b^{4}+5811854553170035414b^{4}+5811854553170035413400a^{7}b^{4}+58118545531700355415b^{4}+581185455556a^{6}+58066a^{$  $+ 645945537218444676642a^9b^4 - 9180121926490903960a^{10}b^4 + 2812829209563463896a^{11}b^4 + 2812866a^{11}b^4 + 281286a^{11}b^4 + 281286a^{11}b^{11}b^{11}b^{11}b^{11}b^{11}b^{11}b^{11}b^{11}b^{11$  $-29546096072985492a^{12}b^4+4881179381635732a^{13}b^4-32252520795880a^{14}b^4$  $+ 3132630060840a^{15}b^4 - 9586673915a^{16}b^4 + 563921995a^{17}b^4 + 1830981648088050432124941b^5 + 563921995a^{17}b^4 + 56392196a^{17}b^4 + 5639216a^{17}b^4 + 563921b^{17}b^4 + 563921b^{17}b$  $+3437881821120782582829456ab^5 + 6475045382105375335889480a^2b^5$  $+ 1038876114726270005445776a^{3}b^{5} + 1596002262124887064872892a^{4}b^{5} \\$  $-85999358210064208a^{11}b^5 + 24458385708121820a^{12}b^5 - 144878605202576a^{13}b^5 - 1448786b^5 - 1448786b^5 - 14686b^5 - 14686b^5 - 14686b^5 - 14686b^5 - 14686b^5 - 1468b^5 - 146b^5 - 146b^5$  $+ 22313057280808a^{14}b^5 - 63714509712a^{15}b^5 + 5752004349a^{16}b^5 + 354501246907809948405480b^6 + 5752004349a^{16}b^5 + 354501246907809948405480b^6 + 5752004349a^{16}b^5 + 575200444b^5 + 57520044b^5 + 57520044b^5 + 57520044b^5 + 57520b^5 + 57520b$  $+ 1069513403791665845357128ab^6 + 523071699408050859929016a^2b^6 \\$  $+735240415476876038312632a^{3}b^{6}+56864612882774626831720a^{4}b^{6}$  $+76635144078132798484872a^{5}b^{6}+874649607643712372472a^{6}b^{6}+2185851421255096595000a^{7}b^{6}$  $-8095043487442343432a^{8}b^{6}+21357540300644860312a^{9}b^{6}-132405723186144408a^{10}b^{6}$  $+77331167721172456a^{11}b^6 - 381004267695688a^{12}b^6 + 100261443987096a^{13}b^6 - 256037937176a^{14}b^6 - 2560379576a^{14}b^6 - 2560379576a^{14}b^6 - 256037656a^{14}b^6 - 256037656a^{14}b^6 - 256037656a^{14}b^{14}b^{14$ 

 $+171242942351303632910536a^{2}b^{7} + 34241820697597088722848a^{3}b^{7}$  $+158910354362756568a^{10}b^7 - 552383916933472a^{11}b^7 + 294616546783832a^{12}b^7$  $-623957998160a^{13}b^7 + 151532656696a^{14}b^7 + 6068045552614043865426b^8$  $+ 1099649665242617000798a^{4}b^{8} + 1064600574662833613834a^{5}b^{8} + 24520882366446890424a^{6}b^{8} + 24520882666466b^{8} + 2452088266666b^{8} + 2452086666b^{8} + 245208666b^{8} + 24520866b^{8} + 2452086b^{8} + 245208b^{8} + 2452086b^{8} + 245208b^{8} + 24520b^{8} + 24520b^{8} + 24520b^{8} + 24520b^{8} + 2450b^{8} + 2450b^{8} + 2450b^{8} +$  $+26672527585599011688a^7b^8+94427080896983406a^8b^8+215842089833789562a^9b^8$  $+371078402291905961224a^{3}b^{9}+352022359346404662630a^{4}b^{9}+18285405023328014000a^{5}b^{9}$  $+ 15396021766853800072a^{6}b^{9} + 196990298922392976a^{7}b^{9} + 195113449113625614a^{8}b^{9} + 1951134b^{9} + 195114b^{9} + 19514b^{9} + 1954b^{9} + 1954$  $+ 353214178214040a^9b^9 + 772341543530700a^{10}b^9 - 457412818200a^{11}b^9 + 800472431850a^{12}b^9 + 800472450a^{12}b^9 + 80047240a^{12}b^{12}b^{12}b^{12}b^{12}b^{12}b^{12}b^{12}b^{$  $+41014337504927410260b^{10}+114934366191849037236ab^{10}+62588562419040481252a^2b^{10}$  $+ 67240314521518855860a^{3}b^{10} + 7157179769538448936a^{4}b^{10} + 5631584078513951272a^{5}b^{10} + 56315860a^{5}b^{10} + 5631560a^{5}b^{10} + 5631560a^{5}b^{10}$  $+157355447527654376a^{6}b^{10} + 117137843987075912a^{7}b^{10} + 774287072717444a^{8}b^{10}$  $+4742709531308534760ab^{11}+6843044908448608444a^2b^{11}+1526743180734652192a^3b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+152674318b^{11}+15267448b^{11}+15267448b^{11}+15267448b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1526748b^{11}+1566b^{11}+156b^{11}$  $+ 652631982303392a^7b^{11} + 434411052282116a^8b^{11} + 1152680301864a^9b^{11} + 960566918220a^{10}b^{11} + 96056691820a^{10}b^{11} + 96056691820a^{11} + 96056691820a^{10}b^{11} + 96056691820a^{10}b^{11} + 96056691820a^{10}b^{11} + 96056691820a^{11} + 96056669180a^{11} + 9605666914a^{11} + 96056669$  $+ 117140064945846306b^{12} + 317038686494337262ab^{12} + 169609751776172344a^2b^{12} + 169609751776172344a^{12} + 169609751761760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609751760760 + 169609750760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 169609760 + 1696097600 + 1696097600 + 1696000 + 16960000 + 1696000 + 1696000 + 169600 + 1696000 +$  $+ 166445736776215624a^{3}b^{12} + 17271058034689004a^{4}b^{12} + 11692563911936180a^{5}b^{12} + 11692563911960a^{5}b^{12} + 11692563910a^{5}b^{12} + 11692563911960a^{5}b^{12} + 116925639160a^{5}b^{12} + 11692563911960a^{5}b^{12} + 11692563910a^{5}b^{12} + 116925663910a^{5}b^{12} + 11692566390a^{5}b^{12} + 11692566390$  $+ 309271965286648a^{6}b^{12} + 181277011099912a^{7}b^{12} + 1013931747010a^{8}b^{12} + 608359048206a^{9}b^{12} + 608359048006a^{9}b^{12} + 608359048006a^{9}b^{12} + 608359048006a^{9}b^{12} + 60835904a^{9}b^{12} + 6083590b^{12} + 6083590b^{12} + 6083590b^{12} + 6083590$  $+4552422878866226b^{13}+8612153544899408ab^{13}+11684556963374568a^{2}b^{13}$  $+2459287660143888a^{3}b^{13}+1832663509135228a^{4}b^{13}+87796234586480a^{5}b^{13}$  $+ 50125620668424a^{6}b^{13} + 512075874352a^{7}b^{13} + 265182149218a^{8}b^{13} + 142380641710440b^{14} + 1423806417104b^{14} + 1423806417104b^{14} + 1423806417104b^{14} + 1423806417104b^{14} + 1423806417104b^{14} + 1423806417104b^{14} + 142380641710b^{14} + 142380641710b^{14} + 1423806b^{14} + 142380b^{14} + 1428b^{14} + 1428b^{14$  $+ 370669544999304ab^{14} + 185095028637272a^2b^{14} + 167607787407128a^3b^{14} + 14882731838680a^4b^{14} + 167607787407128a^3b^{14} + 14882731838680a^4b^{14} + 167607787407128a^3b^{14} + 167607787407128a^{14} + 16760778740718840 + 1676077874071840 + 16760778740718840 + 1676077874071840 + 1676077874071880740 + 1676077874071880740 + 1676077874071880740 + 16760778740718800 + 16760778740718800 + 16760778740718800 + 16760778740718800 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 16760778740 + 167607787760 + 16760778740 + 16760778740 + 16760778740 + 16760778778740 +$  $+8925930851320a^{5}b^{14}+159286274280a^{6}b^{14}+78378960360a^{7}b^{14}+3551773672360b^{15}$  $+6357587526960ab^{15} + 8105951624872a^{2}b^{15} + 1445744454752a^{3}b^{15} + 979484169144a^{4}b^{15}$  $+30677356528a^{5}b^{15} + 15338678264a^{6}b^{15} + 69638530941b^{16} + 172733256855ab^{16}$  $+73945230446a^{2}b^{16} + 61843300298a^{3}b^{16} + 3560764597a^{4}b^{16} + 1917334783a^{5}b^{16} + 1048818981b^{17}$ 

### Notes

2014

Year

 $+ 1689533124ab^{17} + 2006403306a^2b^{17} + 235148940a^3b^{17} + 145008513a^4b^{17} + 11702670b^{18}$  $+ 27096622ab^{18} + 7972286a^2b^{18} + 6096454a^3b^{18} + 91070b^{19} + 114380ab^{19} + 123410a^2b^{19} + 441b^{20}b^{19} + 441b^{20}b^{$  $+903ab^{20} + b^{21}\Big) + \frac{4194304}{\Big[\prod_{\mu=1}^{22} \{a - b - (2\mu - 1)\}\Big] \Big[\prod_{\ell=1}^{21} \{a - b + (2\xi - 1)\}\Big]} \times$ Notes  $\times (13113070457687988603440625 + 32835960506324395192397625a$  $+ 33100902479081461753011150a^2 + 18852771841456100048901630a^3 \\$  $+6998152417792503243516261a^{4}+1830981648088050432124941a^{5}$  $+ 69638530941a^{16} + 1048818981a^{17} + 11702670a^{18} + 91070a^{19} + 441a^{20} + a^{21}a^{10} + 441a^{10} + a^{10}a^{10} + a^{10}a^{10}$ +53745962882620608893167875b + 45053743668222224004344100ab $+113826458136107026075141710a^{2}b + 32581896347358559999317420a^{3}b$  $+22213537883171068151439531a^{4}b+3437881821120782582829456a^{5}b$  $+1069513403791665845357128a^{6}b + 102017962030642073210480a^{7}b$  $+ 17605238308074264053574a^{8}b + 1094032303287108057912a^{9}b + 114934366191849037236a^{10}b + 1094032303287108057912a^{9}b + 1094032303287108b + 10940328057912a^{9}b + 1094032303287108b + 1094032806b + 1094032806b + 1094086b + 1094086b + 1094086b + 1094086b + 1094086b + 1004086b + 1004086b + 1004086b + 100408b + 1004086b + 100408$  $+4742709531308534760a^{11}b + 317038686494337262a^{12}b + 8612153544899408a^{13}b + 317038686494337262a^{12}b + 31703866864943372662a^{12}b + 3170386686494337262a^{12}b + 3170386686494337262a^{12}b + 3170386686494337262a^{12}b + 3170386686494337262a^{12}b + 3170386686494337262a^{12}b + 31703866864943372662a^{12}b + 31703866864943372662a^{12}b + 31703866864943372662a^{12}b + 3170386686494a^{12}b + 317066a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 3170386686464a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 3170386686494a^{12}b + 31703866864666a^{12}b + 3170386686666a^{12}b + 3170666a^{12}b + 317066a^{12}b + 317066a^{12}b$  $+370669544999304a^{14}b + 6357587526960a^{15}b + 172733256855a^{16}b + 1689533124a^{17}b$  $+27096622a^{18}b + 114380a^{19}b + 903a^{20}b - 3418740929846540363004450b^2$  $+152628207170808204647275770ab^{2} + 29662057538012999895338574a^{2}b^{2}$  $+73924679489058918046015938a^{3}b^{2}+9095261732430392323657176a^{4}b^{2}$  $+ 6475045382105375335889480a^5b^2 + 523071699408050859929016a^6b^2 \\$  $+171242942351303632910536a^7b^2 + 9336842670271654008484a^8b^2$  $+ 1689474770765990596076a^9b^2 + 62588562419040481252a^{10}b^2 + 6843044908448608444a^{11}b^2 + 684304490848608444a^{11}b^2 + 684304490866b^2 + 68430449086b^2 + 6843046b^2 + 6843046b^2 + 6843046b^2 + 6843046b^2 + 684304b^2 + 684304b^2 + 684304b^2 + 68430b^2 + 6840b^2 + 6840b^2$  $+8105951624872a^{15}b^2 + 73945230446a^{16}b^2 + 2006403306a^{17}b^2 + 7972286a^{18}b^2 + 123410a^{19}b^2$  $+49302844553810665888575930b^3-5196235439832751151572788ab^3$  $+91735804724807821121609526a^{2}b^{3}+6747725269089701729734368a^{3}b^{3}$  $+ 16438427349229328982545896a^4b^3 + 1038876114726270005445776a^5b^3$  $+735240415476876038312632a^{6}b^{3}+34241820697597088722848a^{7}b^{3}$  $+ 11226186945118492939852a^{8}b^{3} + 371078402291905961224a^{9}b^{3} + 67240314521518855860a^{10}b^{3} + 6724056b^{3} + 672405b^{3} + 67240b^{3} + 6$ 

 $+1526743180734652192a^{11}b^3 + 166445736776215624a^{12}b^3 + 2459287660143888a^{13}b^3 + 167607787407128a^{14}b^3 + 1445744454752a^{15}b^3 + 61843300298a^{16}b^3 + 235148940a^{17}b^3 + 6096454a^{18}b^3 - 6523406582853698567485227b^4 + 43689757687398414790547553ab^4$ 

 $-1858191737565276546672360a^2b^4+19566549200385431103972216a^3b^4$ 

Notes

 $+ 673581690360746582755756a^4b^4 + 1596002262124887064872892a^5b^4 \\$ 

 $+ 56864612882774626831720a^{6}b^{4} + 39096473046294723020936a^{7}b^{4} \\$ 

 $+ 1266877992724507592a^{11}b^4 + 17271058034689004a^{12}b^4 + 1832663509135228a^{13}b^4$ 

 $+ 14882731838680a^{14}b^4 + 979484169144a^{15}b^4 + 3560764597a^{16}b^4 + 145008513a^{17}b^4 + 145008515a^{17}b^4 + 145008515a^{17}b^4 + 145008515a^{17}b^4 + 145008515b^4 + 145008505b^4 + 145008505b^4 + 1450085b^4 + 1450085$ 

 $+ 6412982275781442369534447b^5 - 3020391903873160874624880ab^5 \\$ 

 $+ 11072544506717425040943320a^2b^5 - 251409502832929940580848a^3b^5 \\$ 

 $+ 1850900598832213039588308a^4b^5 + 33362183528683696143184a^5b^5$ 

 $+76635144078132798484872a^{6}b^{5} + 1614107103242780383952a^{7}b^{5} + 1064600574662833613834a^{8}b^{5} \\ + 18285405023328014000a^{9}b^{5} + 5631584078513951272a^{10}b^{5} + 68669517587706672a^{11}b^{5} \\ + 11692563911936180a^{12}b^{5} + 87796234586480a^{13}b^{5} + 8925930851320a^{14}b^{5} + 30677356528a^{15}b^{5} \\ + 1917334783a^{16}b^{5} - 604573745934401191005912b^{6} + 2761535361221837965136664ab^{6} \\ \end{array}$ 

 $-446409836411046907751368a^{2}b^{6} + 1155941471891438860601896a^{3}b^{6} \\ -15624214876016679739352a^{4}b^{6} + 87313680148565150111832a^{5}b^{6} + 874649607643712372472a^{6}b^{6} \\ +1943115716512939776808a^{7}b^{6} + 24520882366446890424a^{8}b^{6} + 15396021766853800072a^{9}b^{6} \\ +157355447527654376a^{10}b^{6} + 46179770718580408a^{11}b^{6} + 309271965286648a^{12}b^{6} \\ \end{array}$ 

 $+50125620668424a^{13}b^6 + 159286274280a^{14}b^6 + 15338678264a^{15}b^6 + 231724339078175706150136b^7 \\ -153811376527605888117264ab^7 + 374921775760059309890200a^2b^7$ 

 $-28831347254232216398432a^{3}b^{7} + 58118545317700335413400a^{4}b^{7} - 490468210861184892336a^{5}b^{7} + 2185851421255096595000a^{6}b^{7} + 12438236313336433600a^{7}b^{7} + 26672527585599011688a^{8}b^{7}$ 

 $+196990298922392976a^{9}b^{7} + 117137843987075912a^{10}b^{7} + 652631982303392a^{11}b^{7} \\ +181277011099912a^{12}b^{7} + 512075874352a^{13}b^{7} + 78378960360a^{14}b^{7} - 15128932813380875312110b^{8} \\ +56881481026993911310578ab^{8} - 13039358045721374943420a^{2}b^{8} + 22139104917821671572676a^{3}b^{8} \\ -921471943096864139266a^{4}b^{8} + 1520878767836758616286a^{5}b^{8} - 8095043487442343432a^{6}b^{8} \\ \end{array}$ 

 $+29720901376476436984a^{7}b^{8}+94427080896983406a^{8}b^{8}+195113449113625614a^{9}b^{8}\\+774287072717444a^{10}b^{8}+434411052282116a^{11}b^{8}+1013931747010a^{12}b^{8}+265182149218a^{13}b^{8}\\+3024008960876320420246b^{9}-2332408473401733698504ab^{9}+4569477008635213654148a^{2}b^{9}\\-492678254183696354488a^{3}b^{9}+645945537218444676642a^{4}b^{9}-15329149883995675728a^{5}b^{9}$ 

 $+ 21357540300644860312a^{6}b^{9} - 69758491996612848a^{7}b^{9} + 215842089833789562a^{8}b^{9} + 21584208983789562a^{8}b^{9} + 215842089833789562a^{8}b^{9} + 21584208983789562a^{8}b^{9} + 2158420898562a^{8}b^{9} + 2158466a^{8}b^{9} + 215846a^{8}b^{9} + 21584b^{8}b^{9} + 215846a^{8}b^{9} + 21584b^{8}b^{9} + 21584b^{8}b^{8} + 21584b^{8}b^{8} + 21584b^{8}b^{8} + 21584b^{8}b^{8} + 21584b^{8} +$  $+353214178214040a^{9}b^{9}+702101231167908a^{10}b^{9}+1152680301864a^{11}b^{9}+608359048206a^{12}b^{9}$  $- 137104871154095942604b^{10} + 454222013907202843900ab^{10} - 120615054891875887772a^2b^{10} + 454222013907202843900ab^{10} + 4542220013907202843900ab^{10} + 4542220013907202843900ab^{10} + 4542220013907202843900ab^{10} + 4542220013907202843900ab^{10} + 4542220013907202843900ab^{10} + 45422000ab^{10} + 4542200ab^{10} + 4542200ab^{10} + 4542200ab^{10} + 4542200ab^{10} + 454220ab^{10} + 454220$  $-132405723186144408a^{6}b^{10}+158910354362756568a^{7}b^{10}-291326223911580a^{8}b^{10}$  $- 13266524009934170456ab^{11} + 22468630494292536788a^2b^{11} - 2744654110430536928a^3b^{11} - 2744654110430536928a^{11} - 274465411 - 2744654110430536928a^{11} - 27446541 - 27446541 - 27446541 - 27446541 - 27446541 - 27446541 - 27446541 - 27446541 - 27446544 - 27446544 - 27446544 - 2744654 - 2744654 - 27466 - 27$  $+ 2812829209563463896a^4b^{11} - 85999358210064208a^5b^{11} + 77331167721172456a^6b^{11} + 77331167721172456a^{11} + 773311677214746a^{11} + 7733116772146a^{11} + 773311677246a^{11} + 77331167726a^{11} + 773316a^{11} + 773316a^{11} + 773316a^{11} + 7733$  $-552383916933472a^7b^{11} + 579276741224844a^8b^{11} - 457412818200a^9b^{11} + 1052049481860a^10b^{11} + 1052049481860a^{10} + 1052049481860a^{11} + 1052049481860a^{11} + 1052049481860a^{11} + 1052049481860a^{11} + 1052049481860a^{11} + 105204948860a^{11} + 10520498860a^{11} + 10520498860a^{11} + 1052049866a^{11} + 10520666a^{11} + 1052066a^{11} + 1052066a^$  $-504464808447965118b^{12} + 1520074513892383306ab^{12} - 422390481555014024a^2b^{12} + 1520074513892383306ab^{12} - 422390481555014024a^2b^{12} + 1520074513892383306ab^{12} - 422390481555014024a^2b^{12} + 1520074513892383306ab^{12} + 152007451389238306ab^{12} + 15200745186ab^{12} + 15200745186ab^{12} + 15200745186ab^{12} + 15200745186ab^{12} + 15200745186ab^{12} + 15200745186ab^{12} + 152007456ab^{12} +$  $+483431815242444952a^{3}b^{12}-29546096072985492a^{4}b^{12}+24458385708121820a^{5}b^{12}$  $+ 36842919196739110b^{13} - 30556991340393392ab^{13} + 46327319945209016a^2b^{13} + 46327319b^{13} + 4632739b^{13} + 4632739b^{13} + 46327319b^{13} + 4632739b^{13} + 4632739b^{13} + 4632739b^{13} + 4632739b^{13} + 4632739b^{13} + 463275b^{13} + 4655b^{13} + 4655b^{13} + 4655b^{13} + 4655b^{13$  $-5648934122340208a^{3}b^{13}+4881179381635732a^{4}b^{13}-144878605202576a^{5}b^{13}$  $+ 100261443987096a^{6}b^{13} - 623957998160a^{7}b^{13} + 421171648758a^{8}b^{13} - 777692306266456b^{14} + 421176b^{13} + 421176b^{14} + 42116b^{14} + 421176b^{14} + 421176b^{14} + 421$  $-32252520795880a^4b^{14}+22313057280808a^5b^{14}-256037937176a^6b^{14}+151532656696a^7b^{14}+150566a^7b^{14}+15056a^7b^{14}+150$  $+ 3132630060840a^4b^{15} - 63714509712a^5b^{15} + 36576848168a^6b^{15} - 484185248451b^{16} + 63714509712a^5b^{16} + 63714509712a^5b^{15} + 36576848168a^6b^{15} - 484185248451b^{16} + 63714509712a^5b^{15} + 36576848168a^6b^{15} - 484185248451b^{16} + 63714509712a^5b^{15} + 36576848168a^6b^{15} - 484185248451b^{16} + 63714509712a^5b^{16} + 6371660760715$  $+ 5752004349a^{5}b^{16} + 13811946279b^{17} - 9489432124ab^{17} + 11986499486a^{2}b^{17} - 837826964a^{3}b^{17} - 94894a^{2}b^{17} - 9484a^{2}b^{17} + 563921995a^4b^{17} - {106564578}b^{18} + {254376218}ab^{18} - {39191490}a^2b^{18} + {32224114}a^3b^{18} - {39191490}a^2b^{18} - {39191490}a^2b^{18} + {39191490}a^2b^{18} - {3919140}a^2b^{18} - {3919140}$  $+1743994b^{19} - 839188ab^{19} + 962598a^{2}b^{19} - 5719b^{20} + 12341ab^{20} + 43b^{21} \Big) \Big\} \Big|$ (8)III. Derivation of the Summation Formula

Substituting  $c = \frac{a+b+45}{2}$  and  $z = \frac{1}{2}$  in equation (2), we get

$$(a-b) {}_{2}F_{1}\left[\begin{array}{cc}a, b \\ \frac{a+b+45}{2} \\ \vdots \end{array}; \quad \frac{1}{2}\right] = a {}_{2}F_{1}\left[\begin{array}{cc}a+1, b \\ \frac{a+b+45}{2} \\ \vdots \end{array}; \quad \frac{1}{2}\right] - b {}_{2}F_{1}\left[\begin{array}{cc}a, b+1 \\ \frac{a+b+45}{2} \\ \vdots \end{array}; \quad \frac{1}{2}\right]$$

Now involving the formula [Salahuddin et. al. p.12-41(8)], the summation formula is obtained.

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# Global Existence and Uniqueness of the Weak Solution in Keller Segel Model

By C. Messikh, A. Guesmia & S. Saadi

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Abstract- This paper deals with the global existence, uniqueness and boundedness of the weak solution for the chemotaxis system (P) defined as

$$\begin{cases} u_t - \Delta u + \operatorname{div} (u \nabla c) = 0 & (t, x) \in \mathbb{R}^+ \times \Omega \\ -\Delta c + \tau c = 0 & x \in \Omega \\ u = 0, \ c = g & \Gamma \\ u (0, x) = u_0 & x \in \Omega \end{cases}$$
(P)

The system (P) is under homogeneous Dirichlet boundary conditions in a convex smooth bounded domain  $\Omega \in \mathbb{R}^n$  with smooth boundary  $\Gamma$  ( $\in H^{\frac{3}{2}}(\Gamma)$  and  $u_0 \in H^{\frac{1}{2}}(\Omega)$  Based on Galerkin's method, Lax -Milgran's and maximum principle, a prove of the existence and uniqueness of a global solution for the system (P) is determined. Moreover we show that the unique solution is positive.

Keywords and Phrases: chemotaxis, global existence, boundedness, positive solution.

GJSFR-F Classification : MSC 2010: 35K58, 45G05, 65C35, 82C22, 82C31,82C80, 92C17







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## Global Existence and Uniqueness of the Weak Solution in Keller Segel Model

C. Messikh °, A. Guesmia ° & S. Saadi  $^{\rho}$ 

Abstract- This paper deals with the global existence, uniqueness and boundedness of the weak solution for the chemotaxis system (P) defined as

 $\begin{cases} u_t - \Delta u + \operatorname{div} (u \nabla c) = 0 & (t, x) \in \mathbb{R}^+ \times \Omega \\ -\Delta c + \tau c = 0 & x \in \Omega \\ u = 0, \ c = g & \Gamma \\ u (0, x) = u_0 & x \in \Omega \end{cases}$ (P)

The system (P) is under homogeneous Dirichlet boundary conditions in a convex smooth bounded domain  $\Omega \in \mathbb{R}^n$  with smooth boundary  $\Gamma$  ( $\in H^{\frac{3}{2}}(\Gamma)$  and  $u_0 \in H^{\frac{1}{2}}(\Omega)$  Based on Galerkin's method, Lax-Milgran's Theorem and maximum principle, a prove of the existence and uniqueness of a global solution for the system (P) is determined. Moreover we show that the unique solution is positive.

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#### I. INTRODUCTION

Chemotaxis is an important means for cellular communication. It is the influence of chemical substances in the environment on the movement of mobile species. This can lead to strictly oriented movement or to partially oriented and partially tumbling movement. The movement towards a higher concentration of the chemical substance is called positive chemotaxis whereas the movement towards regions of lower chemical concentration is called negative chemotactical movement.

The classical chemotaxis model — the so-called Keller–Segel model — system defined in (0.1) was first introduced by Paltak [11] (1953), E. Keller and L. Segel [9] (1970)

$$u_t - \nabla (a\nabla u) + \nabla (\chi u\nabla c) = 0 \qquad (t,x) \in \mathbb{R}^+ \times \mathbb{R}^d$$
  

$$\alpha c_t - \Delta c + \tau c + \beta u = 0 \qquad x \in \mathbb{R}^d \qquad (0.1)$$

where u(t, x) denotes the density of bacteria in the position  $x \in \mathbb{R}^d$  and at time t, c the concentration of chemical signal substance,  $\alpha \geq 0$  the relaxation time, the parameter  $\chi$  the sensitivity of cells to the chemoattractant and  $a, \tau, \beta$  are given smooth functions. As it can be seen, when  $\alpha \neq 0$  the model is Global Journal of Science Frontier Research (F) Volume XIV Issue II Version I

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called Parabolic-Parabolic while it is an Elliptic-Parabolic model when  $\alpha = 0$ . This modelling is very simple, it exhibits a profound mathematical structure and mostly only dimension 2 is understood, especially chemotactic collapse. The proposed model has been extensively studied in the last few years (see ([7]-[8],[12],[13]) for a recent survey articles).

The Parabolic-Parabolic model has been investigated by many authers (see for examples Refs [8] and [13]), I. Fatkullin [10] had developed numerical method ( a composite particle-grid ) with adaptive time stepping which allows us to resolve and propagates singular solutions when with Neumann boundary condition.

The Elliptic-Parabolic model has been investigated by many authers (see for examples Refs [3] and [4]). This model have been carried out where the main concern is whether the solution of model is bounded or blow-up. It has been proved that the solution strongly depends on the spatial dimension. It does not occur in one-dimensional problems, and it occurs conditionally in higher dimensional situations. More precisely, see [2] in case higher dimensions  $(n \ge 3)$ , if the norm of initial condition  $u_0$  is small in space  $L^{\frac{n}{2}}(\mathbb{R}^n)$ , then there are global weak solutions and if  $(\int x^2 u_0)^{d-2} \le C \|u_0\|_{L^1(\mathbb{R}^n)}^n$  with C is small ,then there is blow up in a finite time  $T^*$ . But in two dimension, (see [5]), if  $\|u_0\|_{L^1(\mathbb{R}^2)} < \frac{8\pi}{\chi}$ , there are smooth solutions, and if  $\|u_0\|_{L^1(\mathbb{R}^2)} > \frac{8\pi}{\chi}$ , there is creation of a singular measure (blow-up) in finite time.

In this paper we demonstrate the global existence and uniqueness of weak positive solution for the elliptic-parabolic model's problem defined as

$$(P) \begin{cases} (P_1) \begin{cases} u_t - \Delta u + \operatorname{div} (u\nabla c) = 0 & (t, x) \in \mathbb{R}^+ \times \Omega \\ u = 0 & \Gamma \\ u (0, x) = u_0 & x \in \Omega \\ (P_2) \begin{cases} -\Delta c + \tau c = 0 & x \in \Omega \\ c = g & \Gamma \end{cases} \end{cases}$$

Where u(t, x) is a function denotes the density of bacteria in the position  $x \in \Omega \subset \mathbb{R}^2$  or  $\mathbb{R}^3$ ,  $\Omega$  is a bounded convex domain with smooth boundary  $\Gamma$ , c denotes the concentration of chemical signal that stimulates the bacteria. The parameter  $\tau$  is a time constant and it is expressed on the one hand the movement of bacteria (representing a random distribution side and a deterministic drift in the direction of high concentrations) and secondly the diffusion degradation of c.

To simplify the solution of the system (P), a decomposition of (P) into two subsystems (P<sub>1</sub>) and (P<sub>2</sub>) are adopted. Lax-Milgram's Theorem is very important theorem which we help us to demonstrate the existence and uniqueness of a weak solution for the system (P<sub>2</sub>). However this theorem can not be applied directly because it is a nonhomogenous system. For this raison an adoptation of Trace Theorem is used to simplify the system (P<sub>2</sub>), and together with Galerking method we can demonstrate the existence and uniqueess of a weak solution for the system (P<sub>1</sub>). Therefore we have the existence and uniqueess for the problem (P). Moreover we show that the solution is positive. The following initial-boundary conditions on  $u_0$  and g assumptions are used to prove the proposed solution of (P)

$H_1$	$g \in L^{\frac{1}{2}}\left(\Gamma\right)$	
$H_2$	$g \in L^{\frac{3}{2}}\left(\Gamma\right)$	
$H_3$	$u_0 \in L^2\left(\Omega\right)$	
$H_4$	$u_0 \geq 0$ and $g$	$j \ge j$

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If the hypothesis  $H_1$  is satisfies and using the theorem of trace, one can find a lifting of this trace which we denote  $R(g) \in H_0^1(\Omega)$ . Thus by definition it verifies  $\gamma_0(R(g)) = g$ . Now looking for c having the form  $c = \tilde{c} + R(g)$  reduces the problem (P<sub>2</sub>) to  $\tilde{c}$ .

$$\left(\widetilde{P}_{2}\right) \begin{cases} -\Delta \widetilde{c} + \tau \widetilde{c} - \Delta R\left(g\right) + \tau R\left(g\right) = 0 & \text{in } x \in \Omega \\ \widetilde{c} = 0 & \text{on } \Gamma \end{cases}$$

**Definition 1** We say  $(u, \tilde{c}) \in L^2(0, T; H_0^1(\Omega)) \times H_0^1(\Omega)$  with  $u_t \in L^2(0, T; H^{-1}(\Omega))$  is a weak solution of the problem (P) if and only if

$$\langle u_t, v \rangle + B(u, v, t) = 0 \tag{0.2}$$

$$a\left(\tilde{c},q,t\right) = l\left(q\right) \tag{0.3}$$

where

$$\begin{aligned} B\left(u,v,t\right) &= \int_{\Omega} \left(\nabla u \nabla v + \nabla c \nabla u v + \tau c u v\right) dx \\ a\left(\widetilde{c},q\right) &= \int_{\Omega} \left(\nabla \widetilde{c} \nabla q + \tau \widetilde{c}q\right) dx \\ l\left(q\right) &= -\int_{\Omega} \left(\nabla R\left(g\right) \nabla q + \tau R\left(g\right)q\right) dx \end{aligned}$$

for all  $(v,q) \in \left(H_0^1(\Omega)\right)^2$ ,  $0 \le t \le T$ , and

$$u(0,x) = u_0 \in L^2(\Omega)$$
 (0.4)

**Remark 2** Note that  $u \in C([0,T]; L^2(\Omega))$  as  $u \in L^2(0,T; H_0^1(\Omega))$  and  $u_t \in L^2(0,T; H^{-1}(\Omega))$ . Then equality (0.4) makes sense.

#### II. EXISTENCE OF WEAK SOLUTION OF THE PROBLEM (P)

In this section, use the Theorem of Lax- Milgran to study the existence and uniqueness of weak solution of problem (P<sub>2</sub>), which its variational formulat is given by equation (0.3) and use the method of Galerking to study the existence and uniqueness of weak solution of problem (P<sub>1</sub>), which its variational formulat is given by equation (0.2). So we have the existence and uniqueness of weak solution of problem (P).

#### a) Existence of weak solution of the problem (P2)

**Theorem 3** If the hypothesis  $H_1$  holds. Then the problem  $(P_2)$  has only one solution  $c \in H^1(\Omega)$  for any  $q \in H^1(\Omega)$ .

By applying the Theorem of Lax-Milgran, the solution  $\tilde{c}$  of the problem (0.3) exists and it is unique. So (P<sub>2</sub>) has unique solution.

**Remark 4** Elliptic regularity Theorem remains valid provided that the boundary condition g is in the space  $L^{\frac{3}{2}}(\Gamma)$  which is the image by the operator trace space  $H^{2}(\Omega)$ .

**Remark 5** [6] If  $c \in H^2(\Omega)$  and (c is a solution of problem  $(P_2)$ ) this implies that  $c \in W^{1,q}(\Omega)$   $(H^2(\Omega) \hookrightarrow W^{1,q}(\Omega)$  for  $1 \le q \le 2^*$ ).

Using the Maximum Principle one can show that the solution of the problem  $(P_2)$  is positive as follows. Multiplying the first equation of  $(P_2)$  by  $q \in H_0^1(\Omega)$ , we obtain other variational formulat for problem  $(P_2)$ 

$$\left(\widetilde{\mathbf{P}}_{3}\right)\int_{\Omega}\left(\nabla c\nabla q+\tau cq\right)dx=0.$$

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**Proposition 6** [1] If  $g \in L^{\frac{3}{2}}(\Gamma)$  and  $c \in H^1(\Omega) \cap C(\overline{\Omega})$  then the problem  $(\widetilde{P}_3)$  have a positive solution c.

**Proof.** As  $\Gamma$  is smooth enough and  $g \in L^{\frac{3}{2}}(\Gamma)$  then  $c \in H^2(\Omega)$ . And as  $\Omega \subset \mathbb{R}^2$  or  $\mathbb{R}^3$ , by embedding of Sobolev spaces ( $H^2(\overline{\Omega}) \hookrightarrow C(\overline{\Omega})$ ) this implies that  $c \in C(\overline{\Omega})$ . If  $c = g \ge 0$  on  $\Gamma$ , then  $c^- = \min(c, 0) \in H^1_0(\Omega)$ . So, we have

$$\int_{\Omega} cc^{-} dx = \int_{\Omega} (c^{-})^{2} dx$$
$$\int_{\Omega} \nabla c \nabla c^{-} dx = \int_{\Omega} (\nabla c^{-})^{2} dx$$

Since the support of functions  $c^-$  and  $c^+ = \max(c, 0)$  is set  $A(x) = \{x/u(x) = 0\}$ . This implies that  $\nabla u = 0$  on A(x). As  $c = c^+ + c^-$ , thus we have

$$0 = \int (\nabla c^{-})^{2} + \tau (c^{-})^{2} dx \ge \min(1,\tau) \|c^{-}\|_{H_{0}^{1}(\Omega)}^{2}$$

Finally, we find  $c^- = 0$ .

#### b) Existence of a weak solution of the problem (P1)

Before proving the existence and uniqueness of weak solution of problem  $(P_1)$ , we need the following lemma

**Lemma 7** i) For all  $v \in H_0^1(\Omega)$  then B(.,.,t) is continuous in  $H_0^1(\Omega) \times H_0^1(\Omega)$ , there exists a constant positive C such that

$$|B(u, v, t)| \le C ||u||_{H^{1}(\Omega)} ||v||_{H^{1}(\Omega)}$$
(1.1)

ii) For any  $u \in H_0^1(\Omega)$  and  $H_2$  is hold. Then there exists a constant positive  $\beta$  such that

$$\beta \|u\|_{H^{1}_{0}(\Omega)}^{2} \leq B(u, u, t)$$

**Proof.** i) We use the Cauchy-Shwarz inequality and  $c \in H^2(\Omega) \hookrightarrow L^q(\Omega)$  for any  $q \in \left[1, \frac{2n}{n-2}\right[$  with n = 2 or n = 3, we obtain i) as follows

$$\begin{array}{lll} B\left(u,v,t\right) &\leq & \|\nabla u\|_{L^{2}(\Omega)} \|\nabla v\|_{L^{2}(\Omega)} + \|\nabla c\|_{L^{4}(\Omega)} \|u\|_{L^{2}(\Omega)} \|v\|_{L^{4}(\Omega)} \\ & & +\tau \|c\|_{L^{4}(\Omega)} \|u\|_{L^{2}(\Omega)} \|v\|_{L^{4}(\Omega)} \\ &\leq & C \|u\|_{H^{1}(\Omega)} \|v\|_{H^{1}(\Omega)} \end{array}$$

ii) Making use of  $-\Delta c + \tau c = 0$  the expression of B(u, u, t) becomes

$$B(u, u, t) = \int (\nabla u)^2 + \frac{\nabla c}{2} \nabla u^2 + \tau c u^2 dx$$
  
$$= \int (\nabla u)^2 + \left(\tau c - \frac{\Delta c}{2}\right) u^2 dx$$
  
$$= \int (\nabla u)^2 + \frac{1}{2} \tau c u^2 dx \ge \|\nabla u\|_{L^2(\Omega)}^2$$

### Notes

Finally, by Poincarre inequality yields

$$B\left(u, u, t\right) \ge \beta \left\|u\right\|_{H_{0}^{1}(\Omega)}^{2}.$$

To demonstrate the existence of weak solution of  $(P_1)$  via the method of Galerking, we assume  $w_k = w_k(x)$  are smooth functions verifying

$$\{w_k\}_{k=1}^{\infty}$$
 is an orthogonal basis of  $H_0^1(\Omega)$  (1.2)

and

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$$\{w_k\}_{k=1}^{\infty}$$
 is an orthonormal basis of  $L^2(\Omega)$ . (1.3)

Consider a positive integer m. We will look for a function  $u_m : [0 T] \to H_0^1(\Omega)$  of the form

$$u_m(t) := \sum_{k=1}^{m} d_m^k(t) w_k$$
(1.4)

which satisfies

$$d_m^k(0) = (u_0, w_k) \tag{1.5}$$

and

$$\langle u'_m, w_k \rangle + B(u_m, w_k, t) = 0, \quad 0 \le t \le T \text{ and } k = 1, ..., m$$
 (1.6)

where  $u' = u_t$  and here (.,.) denotes the scalar product in  $L^2(\Omega)$ .

**Theorem 8** (construction of the approximate solution) For each integer m, there exists a unique function  $u_m$  of the form (1.4) satisfying (1.5) and (1.6).

**Proof.** Assuming  $u_m$  has the structure (1.4). Substituting (1.4) into (1.5) and using (1.3) we obtained

$$d_m^{\prime k}(t) + \sum_{l=1}^m d_m^l B(w_l, w_k, t) = 0 \qquad 0 \le t \le T \text{ and } k = 1, ..., m$$
(1.7)

According to standard existence theory for ordinary differential equations, there exists a unique absolutely continuous functions  $d_m(t) = (d_m^1, d_m^2, ..., d_m^m)$  satisfying (1.5) and (1.7). So  $u_m$  of the form (1.4) satisfies (1.5) and (1.6) for all  $t \in [0 T]$ .

#### c) Energy estimates

We propose now to send m to infinity and show a subsequence of our solutions  $u_m$  of the approximation problems (1.5) and (1.6) converges to a weak solution of (P<sub>1</sub>). For this we will need some uniform estimates.

**Theorem 9** (Energy estimates) [6]. There exists a constant C, depending only on  $\Omega$ , T and c, such that

$$\max_{0 \le t \le T} \|u_m\|_{L^2(\Omega)} + \|u_m\|_{L^2(0,T; H_0^1(\Omega))} + \|u'_m\|_{L^2(0,T; H^{-1}(\Omega))} \le C \|u_0\|_{L^2(\Omega)}$$
(1.8)

for m = 1, 2, ...

**Proof.** 1. Multiplying equation (1.6) by  $d_m^k(t)$ , summing for k = 1, ..., m, and then recalling (1.4) we find

$$(\acute{u}_m, u_m) + B(u_m, u_m, t) = 0 \tag{1.9}$$

for all  $0 \le t \le T$ . From Lemma 7, there exists constant  $\beta > 0$  such that

$$\beta \|u_m\|_{H^1_0(\Omega)}^2 \le B(u_m, u_m, t)$$
(1.10)

for all  $0 \le t \le T$ . Consequently (1.10) yields the inequality

$$\frac{d}{dt} \left( \left\| u_m \right\|_{L^2(\Omega)}^2 \right) + \beta \left\| u_m \right\|_{H^1_0(\Omega)}^2 \le 0 \text{ for all } 0 \le t \le T.$$
(1.11)

This implies that

$$\|u_m\|_{L^2(\Omega)}^2 \le \|u_m(0)\|_{L^2(\Omega)}^2 \le \|u_0\|_{L^2(\Omega)}^2 \text{ for all } 0 \le t \le T.$$
(1.12)

So we have

$$\max_{0 \le t \le T} \|u_m\|_{L^2(\Omega)} \le \|u_0\|_{L^2(\Omega)} \,. \tag{1.13}$$

Notes

2. Integrate inequality (1.11) from 0 to T and we employ the inequality (1.13) to find

$$\|u_m\|_{L^2(0,T; H^1_0(\Omega))}^2 = \int_0^T \|u_m\|_{H^1_0(\Omega)}^2 dt \le C \|u_0\|_{L^2(\Omega)}^2.$$

3. Fix any  $v \in H_0^1(\Omega)$ , with  $\|v\|_{H_0^1(\Omega)}^2 \leq 1$ , and write  $v = v^1 + v^2$ , where  $v^1 \in span(w_k)_{k=1}^{k=m}$ , and  $(v^2, w_k) = 0$  (k = 1, ..., m). We use (1.6), we deduce for all  $0 \leq t \leq T$  that

$$(u'_m, v^1) + B(u_m, v^1, t) = 0.$$

Then (1.4) implies

$$\langle u'_m, v \rangle = (u'_m, v) = (u'_m, v^1) = -B(u_m, v^1, t),$$

consequently

$$|\langle u'_m, v \rangle| \le C \, \|u_m\|_{H^1_0(\Omega)} \, .$$

Since  $||v^1||^2_{H^1_0(\Omega)} \le ||v||^2_{H^1_0(\Omega)} \le 1$ . Thus

$$||u'_m||_{H^{-1}(\Omega)} \le C ||u_m||_{H^1_0(\Omega)},$$

and therefore

$$\left\|u_{m}'\right\|_{L^{2}(0,T;\ H^{-1}(\Omega))}^{2} = \int_{0}^{T} \left\|u_{m}'\right\|_{H^{-1}(\Omega)}^{2} dt \leq C \int_{0}^{T} \left\|u_{m}\right\|_{H^{1}_{0}(\Omega)}^{2} dt \leq C \left\|u_{0}\right\|_{L^{2}(\Omega)}^{2}.$$

#### d) Existence and uniqueness

Next we pass to limits as  $m \to \infty$ , to build a weak solution of our initial boundary-value problem (P<sub>1</sub>).

**Theorem 10** (Existence of weak solution). Under hypothesis  $H_2$  and  $H_3$ , there exists a weak solution of  $(P_1)$ .

**Proof.** 1. According to the energy estimates (1.8), we see that the sequence  $\{u_m\}_{m=1}^{\infty}$  is bounded in  $L^2(0,T; H_0^1(\Omega))$  and  $\{u'_m\}_{m=1}^{\infty}$  is bounded in  $L^2(0,T; H^{-1}(\Omega))$ . Consequently there exists a subsequence which is also noted by  $\{u_m\}_{m=1}^{\infty}$  and a function  $u \in L^2(0,T; H_0^1(\Omega))$ , with  $u' \in L^2(0,T; H^{-1}(\Omega))$ , such that

$$u_m \rightharpoonup u \qquad \text{weakly in } L^2\left(0, T; H_0^1\left(\Omega\right)\right) \qquad (1.14)$$
  
$$u'_m \rightharpoonup u' \qquad \text{weakly in } L^2\left(0, T; H^{-1}\left(\Omega\right)\right).$$

2. Next fix an integer N and choose a function  $v \in C^1(0,T; H^1_0(\Omega))$  having the form

$$v(t) = \sum_{k=1}^{N} d^{k}(t) w_{k}$$
(1.15)

where  $\{d^k\}_{k=1}^N$  are given smooth functions. We choose  $m \ge N$ , multiply equation (1.6) by  $d^k(t)$ , sum for k = 1, ..., N, and then integrate with respect to t to find

$$\int_{0}^{T} \langle u'_{m}, v \rangle + B(u_{m}, v, t) dt = 0.$$
(1.16)

We recall (1.14) to find upon passing to weak limits that

$$\int_0^T \langle u', v \rangle + B\left(u, v, t\right) dt = 0 \qquad \forall v \in L^2\left(0, T; H_0^1\left(\Omega\right)\right).$$
(1.17)

As functions of the form (1.15) are dense in  $L^{2}(0,T; H_{0}^{1}(\Omega))$ . Hence in particular

$$\langle u', v \rangle + B(u, v, t) dt = 0 \qquad \forall v \in H_0^1(\Omega) \text{ and } \forall t \in [0 T], \qquad (1.18)$$

and from Remark 2 we have  $u \in C(0,T; L^{2}(\Omega))$ .

3. In order to prove  $u(0) = u_0$ , we first note from (1.17) that

$$\int_{0}^{T} -\langle u, v' \rangle + B(u, v, t) dt = (u(0), v(0))$$
(1.19)

for each  $v \in C^1(0,T; H^1_0(\Omega))$  with v(T) = 0. Similarly, from (1.16) we deduce

$$\int_{0}^{T} -\langle u_{m}, v' \rangle + B(u_{m}, v, t) dt = (u_{m}(0), v(0)).$$
 (1.20)

We use again (1.14), we obtain

$$\int_{0}^{T} -\langle u, v' \rangle + B(u, v, t) dt = (u_{0}, v(0)), \qquad (1.21)$$

since  $u_m(0) \to u_0$  in  $L^2(\Omega)$ . Comparing (1.19) and (1.21), we conclude  $u(0) = u_0$ .

**Theorem 11** (Uniqueness of weak solutions) A weak solution of  $(P_1)$  is unique. **Proof.** We suppose there exists two weak solution  $u_1$  and  $u_2$ . We put

$$U = u_2 - u_3$$

then U is also a solution of  $(P_1)$  with  $U_0 = (u_2 - u_1)(0) \equiv 0$ . Setting v = U in identity (1.18) we have

$$\frac{d}{dt}\left(\frac{1}{2} \|U\|_{L^{2}(U)}^{2}\right) + B\left(U, U, t\right) = 0.$$

From Lemma 7 we have  $B(U, U, t) \geq \beta \|U\|_{H_0^1(U)}^2 \geq 0$ , so  $\frac{d}{dt} \left(\frac{1}{2} \|U\|_{L^2(U)}^2\right) \leq 0$ , then integrate with respect to t to find

$$||U||_{L^{2}(\Omega)}^{2} \leq ||U_{0}||_{L^{2}(\Omega)}^{2} = 0,$$

thus  $U \equiv 0$ .

### Notes

#### e) Global solution of problem (P)

Our main results in this paper are stated as follows.

**Theorem 12** i) if  $c \ge c_0 > 0$ . Then la solution (u, c) of problem (P) is global ii) if  $c \ge c_0 > 0$ . Then la solution (u, c) of problem (P) is global. Furthermore there exists  $\tau_0 > 0$  such that  $||u||_{L^2} \le e^{-\tau_0 t} ||u_0||_{L^2}$ .

**Proof.** We put

$$E(t) = \frac{1}{2} \int_{\Omega} u^2 dx \tag{1.22}$$

We derivate the equation (1.22) and we use firsts equations of  $(P_1)$  and  $(P_2)$  to find

i) We have

$$\frac{dE}{dt} = -B\left(u, u, t\right) \le 0,$$

therefore

$$E\left(t\right) \leq E\left(0\right).$$

ii) We have

$$\frac{dE}{dt} = -B(u, u, b, t) = -\int \left( (\nabla u)^2 + \frac{1}{2}\tau c u^2 \right) dx \le \frac{-1}{2}\tau c_0 \|u\|_{L^2(\Omega)}^2 = -\tau_0 E(t)$$

This implies that

$$E(t) \le E(0) e^{-\tau_0 t}.$$

**Proposition 13** [1] Let  $u_0 \in L^2(\Omega)$  and  $u \in C([0,T]; L^2(\Omega)) \cap L^2(0,T; H_0^1(\Omega))$  is the unique weak solution of  $(P_1)$ . If  $u_0 \ge 0$  in  $\Omega$ , then  $u \ge 0$  in  $]0, T[ \times \Omega$ . **Proof.** If  $u_0 \ge 0$  on  $\Gamma$ . Therefore  $u^- = \min(u, 0) \in L^2(]0, T[; H_0^1(\Omega))$ . A reasoning similar to the Proposition 6, we obtain for all  $0 \le t \le T$ 

$$\frac{1}{2}\frac{d}{dt}\int_{\Omega} \left(u^{-}\right)^{2} dx + \int_{\Omega} B\left(u^{-}, u^{-}, t\right) dx = 0$$

Using the Lemma 7 and integrating with respect to  $\tau$  from 0 to t, we get

$$\frac{1}{2} \int_{\Omega} (u^{-})^{2} dx + \beta \int_{0}^{t} \|u(s)\|_{H_{0}^{1}(U)}^{2} ds \leq \frac{1}{2} \int_{\Omega} (u^{-}(0))^{2} dx = 0.$$

Since  $u^{-}(0) = (u_0)^{-} = 0$ . So  $u^{-} = 0$ .

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### Effect of Variable Thermal Conductivity & Heat Source/Sink Near a Stagnation Point on a Linearly Stretching Sheet using HPM

### By Vivek Kumar Sharma & Aisha Rafi

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Abstract- Aim of the paper is to investigate effects of variable thermal conductivity on flow of a viscous incompressible fluid in variable free stream near a stagnation point on a non-conducting stretching sheet. The equations of continuity, momentum and energy are transformed into ordinary differential equations and solved numerically using Similarity transformation and Homotopy Perturbation Method. The velocity and temperature distributions are discussed numerically and presented through graphs. Skin-friction coefficient and the Nusselt number at the sheet are derived, discussed numerically and their numerical values for various values of physical parameter are presented through Tables.

*Keywords:* homotopy perturbation method, similarity transformation method, steady, boundary layer, variable thermal conductivity, stretching sheet, skin-friction coefficient and nusselt number.

GJSFR-F Classification : MSC 2010: 00A69

EFFECTOFVARIABLETHERMALCONDUCTIVITYHEATSOURCESINKNEARASTAGNATIONPOINTONALINEARLYSTRETCHINGSHEETUSINGHPM

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# Effect of Variable Thermal Conductivity & Heat Source/Sink Near a Stagnation Point on a Linearly Stretching Sheet using HPM

Vivek Kumar Sharma <sup>a</sup> & Aisha Rafi <sup>a</sup>

Abstract- Aim of the paper is to investigate effects of variable thermal conductivity on flow of a viscous incompressible fluid in variable free stream near a stagnation point on a non-conducting stretching sheet. The equations of continuity, momentum and energy are transformed into ordinary differential equations and solved numerically using Similarity transformation and Homotopy Perturbation Method. The velocity and temperature distributions are discussed numerically and presented through graphs. Skin-friction coefficient and the Nusselt number at the sheet are derived, discussed numerically and their numerical values for various values of physical parameter are presented through Tables.

Keywords: homotopy perturbation method, similarity transformation method, steady, boundary layer, variable thermal conductivity, stretching sheet, skin-friction coefficient and nusselt number.

### I. INTRODUCTION

Study of heat transfer in boundary layer find applications in extrusion of plastic sheets, polymer, spinning of fibers, cooling of elastic sheets etc. The quality of final product depends on the rate of heat transfer and therefore cooling procedure has to be controlled effectively. Liquid metals have small Prandtl number of order 0.01~ 0.1(e.g. Pr = 0.01 is for Bismuth, Pr = 0.023 for Mercury etc.) and are generally used as coolants because of very large thermal conductivity.

Aim of the present paper is to investigate effects of variable thermal conductivity, heat source/sink and variable free stream on flow of a viscous incompressible electrically conducting fluid and heat transfer on a non-conducting stretching sheet. Linear stretching of the sheet is considered because of its simplicity in modelling of the flow and heat transfer over stretching surface and further it permits the similarity solution, which are useful in understanding the interaction of flow field with temperature field. The heat source and sink is included in the work to understand the effect of internal heat generation and absorption [Chaim (1998)].

The Homotopy Perturbation Method is a combination of the classical perturbation technique and homotopy technique, which has eliminated the limitations of the traditional perturbation methods. This technique can have full advantage of the traditional perturbation techniques. J. H. He, Approximate analytical solution for seepage flow with fractional derivatives in porous media. J. H. He, A coupling method of homotopy technique and perturbation technique for nonlinear problems. To illustrate the basic idea of the Homotopy Perturbation Method for solving nonlinear differential equations, we consider the following nonlinear differential equation:

### A(u) - f(r) = 0

Subject to boundary condition

$$B\left(u,\frac{\partial u}{\partial n}\right) = 0 \tag{2}$$

Notes

Where A is a general differential operator, B is a boundary operator, f(r) is a known analytic function, and  $\Gamma$  is the boundary of the domain  $\Omega$ .

The operator A can, generally speaking, be divided into two parts: a linear part L and a nonlinear part N. Equation can be rewritten as follows:

$$L(u) + N(u) - f(r) = 0$$
 (3)

By the homotopy technique, we construct a homotopy  $V(r,p): \Omega^*(0,1) \rightarrow \mathbb{R}$  which satisfy

$$H(V,p) = (1-p)[L(v) - L(u_0)] + p[A(v) - f(r)] = 0$$
(4)

$$H(V,p) = L(v) - L(u_0) + p L(u_0) + p[N(v) - f(r)] = 0$$
(5)

Where  $p \in [0, 1]$  is an embedding parameter and  $u_0$  is an initial approximation of which satisfies the boundary conditions.

$$H(V,0) = L(v) - L(u_0)$$
  
H(V,1) = A(v) - f(r)] (6)

Thus, the changing process of p from zero to unity is just that of v(r, p) from  $u_0(r)$  to u(r). In Topology, this is called deformation and  $L(v) - L(u_0)$ , A(v) - f(r) are called homotopic. According to the HPM, we can first use the embedding parameter p as a "small parameter," and assume that the solution of can be written as a power series in p:

$$V = V_0 + pV_1 + p^2 V_2 + \dots$$
(7)

Setting p=1 results in the approximate solution of

 $u = \lim V = V_0 + V_1 + \dots$  (8)

p→1

The series is convergent for most cases; however, the convergent rate depends upon the Nonlinear operator A(V). The second derivative of N(V) with respect to V must be small because the parameter may be relatively large; that is,  $p \to 1$ .

In this paper is to investigate effects of variable thermal conductivity on flow of a viscous incompressible fluid in variable free stream near a stagnation point on a non-conducting stretching sheet.

### II. FORMULATION OF THE PROBLEM

Consider steady two-dimensional flow of a viscous incompressible electrically conducting fluid of variable thermal conductivity in the vicinity of a stagnation point on a non-conducting stretching sheet It is assumed that external field is zero, the electric field owing to polarization of charges and Hall Effect are neglected. Stretching sheet is placed in the plane y = 0 and x-axis is taken along the sheet The fluid occupies the upper half plane i.e. y>0. The governing equations are:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 , \qquad (9)$$

$$u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} = -\frac{1}{\rho}\frac{\partial p}{\partial x} + v\frac{\partial^2 u}{\partial y^2},$$
(10)

$$\rho C_p \left( u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = \frac{\partial}{\partial y} \left( k^* \frac{\partial T}{\partial y} \right) + Q \left( T - T_{\infty} \right) , \qquad (11)$$

where  $\varepsilon$  -perturbation parameter,  $\eta$ -similarity parameter { =  $(c/v)^{\frac{1}{2}} y$ },  $\eta_{\infty}$ -value of  $\eta$  at which boundary conditions is achived,  $\kappa$ -uniform thermal conductivity,  $\kappa^*$ -variable thermal conductivity, v-kinematic viscosity,  $\rho$ -density of fluid,  $\psi$ -stream function,  $\sigma$ -electrical conductivity,  $\theta$ -dimensionless temperature{ =  $(T - T_{\infty}) / (T_w - T_{\infty})$ },  $\tau_w$ -shear stress, S-heat source/sink parameter {=  $Q/\rho C_p c$ }, T-fluid temperature.

The second derivatives of u and T with respect to x have been eliminated on the basis of magnitude analysis considering that Reynolds number is high. Hence the Navier-Stokes equation modifies into Prandtl's boundary layer equation. The boundary conditions are.

$$y = 0: \quad u = u_w(x) = c \quad v = 0, \quad T = T_w$$
  
$$y \to \infty: \quad u = U(x) = bx, \quad T = T_\infty$$
(12)

Introducing the stream function  $\psi(x, y)$  as defined by

$$u = \frac{\partial \psi}{\partial y}$$
,  $v = -\frac{\partial \psi}{\partial x}$ , (13)

the similarity variable  $\eta = (c / v)^{1/2} y$  and

$$\Psi(x, y) = (c v)^{1/2} x f(\eta),$$
(14)

into the equations (3) and (5), we get

$$f''' + ff'' - (f')^2 + \lambda^2 = 0, \qquad (15)$$

And

$$(1+\varepsilon) \theta'' + \varepsilon (\theta')^2 + \Pr \theta' f + \Pr S\theta = 0.$$

The governing boundary layer and thermal boundary layer equations (15) and (16) with the boundary conditions (17) are solved using Homotopy Perturbation Method.

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Equations (15) and (16) are non-linear coupled differential equation. To solve these equations, we introduce the following Homotopy.

$$D(f,p) = (1-p)\left[\left(\frac{d^{3}f}{d\eta^{3}}\right) - \left(\frac{d^{3}f_{I}}{d\eta^{3}}\right)\right] + p\left[\frac{d^{3}f}{d\eta^{3}} + f\frac{d^{2}f}{d\eta^{2}} - \left(\frac{df}{d\eta}\right)^{2} + \lambda^{2}\right] = 0 \quad (16)$$

$$D(\theta, p) = (1 - p) \left[ \left( \frac{\partial^2 \theta}{\partial \eta^2} + P_r S \theta \right) - \left( \frac{\partial^2 \theta_I}{\partial \eta^2} + P_r S \theta_I \right) \right]$$
Note

$$+P\left[\left(1+\varepsilon\theta\right)\frac{\partial^{2}\theta}{\partial n^{2}}+\varepsilon\left(\frac{\partial\theta}{\partial n}\right)^{2}+P_{r}\frac{\partial\theta}{\partial n}f+P_{r}S\theta\right]=0$$
(17)

With the following assumption

$$f = f_0 + pf_1 + p^2 f_2 + \cdots$$
 (18)

$$\theta = \theta_0 + p\theta_1 + p^2\theta_2 \tag{19}$$

Using equation (18),(19) into equation (10) and (11) and on comparing the like powers of p, we get the zeoth order equation,

$$\left[ \left( \frac{d^3 f_0}{d\eta^3} \right) - \left( \frac{d^3 f_I}{d\eta^3} \right) \right] + \left[ \frac{d^3 f_0}{d\eta^3} + f_0 \frac{d^2 f_0}{d\eta^2} - \left( \frac{d f_0}{d\eta} \right)^2 + \lambda^2 \right] = 0,$$
(21)

$$\left[ \left( \frac{\partial^2 \theta_0}{\partial \eta^2} + P_r S \theta_0 \right) - \left( \frac{\partial^2 \theta_I}{\partial \eta^2} + P_r S \theta_I \right) \right] + \left[ (1 + \varepsilon \theta_0) \frac{\partial^2 \theta_0}{\partial n^2} + \varepsilon \left( \frac{\partial \theta_0}{\partial n} \right)^2 + P_r \frac{\partial \theta_0}{\partial n} f + P_r S \theta_0 \right]$$
  
= 0, (21)

with the corresponding boundary conditions are of zeroth order equations are:

 $\eta = 0: f_0 = 0, f_0^{'} = 1, \theta_0 = 1;$  (22)

$$\eta = \infty; \quad f_0' = \lambda, \quad \theta_0 = 0 \tag{23}$$

$$\frac{d^{3}f_{1}}{d\eta^{3}} + (e^{-n} - \lambda e^{-\eta}) + (\lambda\eta - e^{-n} + \lambda e^{-n} - \lambda + 1)(\lambda e^{-\eta} - e^{-\eta}) - (\lambda + e^{-n} - \lambda e^{-n})(\lambda + e^{-n} - \lambda e^{-n}) + \lambda^{2} = 0$$
(24)

$$\frac{\partial^2 \theta_I}{\partial \eta^2} + P_r S \theta_I + e^{-\eta} (1 + \varepsilon e^{-\eta}) + \varepsilon e^{-2\eta} - P_r e^{-\eta} (\lambda \eta - e^{-\eta} + \lambda e^{-\eta} - \lambda + 1) - P_r S e^{-\eta}$$
$$= 0$$
(25)

With the corresponding boundary conditions are of first order equations are:

$$\eta = \mathbf{0}; \ f_0 = \mathbf{0}, \ f'_0 = \mathbf{0}, \ \theta_0 = \mathbf{0} \\ \eta = \infty; \quad f'_0 = \mathbf{0}, \quad \theta_0 = \mathbf{0}; \end{cases}$$
(26)

Solving equations with corresponding boundary conditions, the following functions can be obtained successively, by summing up the results, and  $p \to 1$  we write the  $f(\boldsymbol{\eta})$ ,  $\boldsymbol{\theta}(\boldsymbol{\eta})$ , profile as:

$$f(\eta) = \lambda \eta - e^{-\eta} + \lambda e^{-\eta} - \lambda + 1 + (\eta e^{-\eta} + 4e^{-\eta} + 3\eta - 4)$$
(27)

$$\theta(\eta) = e^{-\eta} + \frac{A_2 e^{-\eta}}{(A_1)^2 + 1} + \frac{A_3 e^{-2\eta}}{(A_1)^2 + 4} + \frac{A_4 e^{-\eta}}{(A_1)^2} \left(\eta - \frac{\eta - 2}{(A_1)^2}\right) - \left(\frac{2A_4}{(A_1)^4} + \frac{A_3}{(A_1)^2 + 4} + \frac{A_2}{(A_1)^2 + 1}\right) e^{-\eta}$$
(28)

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where 
$$A_1 = \sqrt{P_r S}$$
:  $A_2 = P_r S + P_r - \lambda P_r - 1$ :  $A_3 = \lambda P_r - P_r - 2$ :  $A_4 = \lambda P_r$  (29)

Skin-Friction: Skin-friction coefficient at the sheet is given by

$$C_f = x f''(0)$$
 (30)

**Nusselt Number:** The rate of heat transfer in terms of the Nusselt number at the sheet is given by

$$Nu = -\theta'(0) \tag{31}$$

### III. Conclusion

It is observed from Table 1 as L increases, the numerical values of f''(0) also increase. It is noted from Table 2 that the numerical values of  $-\theta'(0)$  increase when  $\lambda$ increases and  $-\theta'(0)$  decreases when  $\varepsilon$  increases. The skin-friction coefficient and Nusselt number are presented by equations (30) and (31) and they are directly proportional to f''(0) and  $-\theta'(0)$  respectively. The effects of  $\varepsilon$ , Pr and S on Nusselt number have been presented through Table 3 respectively.

#### Table 1

λ	f"'(0)	λ	f''(0)
0.0	-1	0.1	-1.0800
0.01	-1.0098	1.0	0.0004
0.05	-1.0450	2.0	2.0175

#### Table 2

λ	- heta'(0)	ε	- heta'(0)
0.1	.81235	0.0	0.223558
0.5	.13629	0.05	0.215792
2.0	.24133	0.1	0.204672

#### Table 3

ε	S	$P_{r}$	Nu
0	0	0.5	2.010357
0.1	0	0.5	3.542315
0	-0.1	1	2.565423
0.1	-0.1	1	2.845633





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From figure 1, we observe that as  $\lambda$  increases, value of f' also increases. From figure 2 it is observed that when  $\lambda$  increase simultaneously  $\theta$  also increases. In figure 3,  $\lambda$ , S and Pr are constant but when  $\varepsilon$  increases  $\theta$  will also increased. It is observed in figure 4, s,  $\varepsilon$  and  $\lambda$  are constant, when Pr increases,  $\theta$  will also increase. Figure 5 is a physical model which becomes clearer from figure, 6 and 7.

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# Homeotopy Groups of 2–Dimensional Manifolds with One Boundary Component

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*Introduction*- Let Y be a compact, connected 2—dimensional manifold with boundary. The homeotopy group of Y, denoted H(Y), is defined to be the group of isotopy classes in the space of all homeomorphisms of Yonto Y. This group (also known as the mapping class group) has been studied for various manifolds (see, for example, [2] and [3]). It is also possible to consider "subhomeotopy groups" where there are restrictions placed on the action of the homeomorphisms on the boundary of Y (see, for example, [7] and [8]). In this note we will consider the special case of a compact, connected manifold with exactly on boundary component. For the remainder of this paper we will assume Y represents a compact, connected manifold with exactly one boundary component and we willlet X denote the closed 2—manifold obtained by sewing a disk to the boundary of Y. Let Aut  $\pi_1(X,x_0)$  denote the group of automorphisms of  $\pi_1(X,x_0)$  where  $x_0\varepsilon$ — Bd(Y). In this paper we establish the following result. Theorem. If Y is not aMoebius band or a disk, then H(Y)=Aut  $\pi_1(X, x_0)$ .

GJSFR-F Classification : MSC 2010: 57Q05, 57Q91



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## Homeotopy Groups of 2—Dimensional Manifolds with One Boundary Component

David J. Sprows

### I. INTRODUCTION

Let Y be a compact, connected 2—dimensional manifold with boundary. The homeotopy group of Y, denoted H(Y), is defined to be the group of isotopy classes in the space of all homeomorphisms of Y onto Y. This group (also known as the mapping class group) has been studied for various manifolds (see, for example, [2] and [3]). It is also possible to consider "subhomeotopy groups" where there are restrictions placed on the action of the homeomorphisms on the boundary of Y (see, for example, [7] and [8]). In this note we will consider the special case of a compact, connected manifold with exactly on boundary component. For the remainder of this paper we will assume Y represents a compact, connected manifold with exactly one boundary component and we willlet X denote the closed 2—manifold obtained by sewing a disk to the boundary of Y. Let Aut  $\pi_1(X,x_0)$  denote the group of automorphisms of  $\pi_1(X,x_0)$  where  $x_0 \in X$ — Bd(Y). In this paper we establish the following result.

Theorem. If Y is not a Moebius band or a disk, then  $H(Y) = Aut \pi_1(X, x_0)$ 

### II. Proof of the Theorem

Let  $[(X, x_0), (X, x_0)]$  denote the set of homotopy classes (rel  $x_0$ ) of maps from  $(X, x_0)$  to  $(X, x_0)$  and let [f] denote the homotopy class (rel  $x_0$ ) of a mapping f from  $(X, x_0)$  to  $(X, x_0)$ . Let End  $\pi_1$   $(X, x_0)$  denote the set of endomorphisms of  $\pi_1(X, x_0)$  with the operation of composition.

<u>Lemma 1</u>. If X is a closed 2-manifold with  $\pi_2(X, x_0) = 0$  and  $\phi : [(X, x_0), (X, x_0)] \rightarrow \text{End} \pi_1(X, x_0)$  is given by  $\phi$  ([f]) =f<sub>\*</sub>, then  $\phi$  is a bijection which preserves the operation of composition.

Proof of lemma 1.

- 1. Clearly  $\phi$  is well defined and  $\phi([f \circ g]) = \phi([f]) \circ \phi([g])$
- 2. Claim  $\phi$  is a surjection.

Let F in END  $\pi_1(X,x_0)$ . We define f:  $(X,x_0) \to (X,x_0)$  as follows:

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Take  $x_0$  to be vertex of the triangulation of X and let T be a maximal tree in X. If  $x \in T$ , we let  $f(x) = x_0$ . Now, suppose s is a 1—simplex not in T with h: [0,1] = s. Let  $\Gamma_i$ , be a path in T from  $x_0$  to h(i), i = 0, 1. Define the loop  $\alpha_s$  at  $x_0$  by letting  $\alpha_s$   $(t) = \Gamma_0(t)$  if  $-1 \le t \le 0$ ,  $\alpha_s(t) = h(t)$  if  $0 \le t \le 1$ , and  $\alpha_s(t) = \Gamma_0^{-1}(t)$  if  $1 \le t \le 2$ . If  $F[\alpha_s] = [\beta]$ , we let  $f/s = \beta \circ h^{-1}s$ . This defines f on the 1-skeleton of X.

Finally, if  $\Delta$  is a 2-simplex with edges s1, s2, and s3, then  $[\alpha_{S1} * \alpha_{S2} * \alpha_{S3}] = 1$ . This means  $f/\delta\Delta$  is null homotopic, i.e., f extends to  $\Delta$ . Hence, the mapping of f defined on the 1-skeleton as above, extends to a mapping defined on all of the 2-dimensional manifold X. Note that by construction,  $f_*[\alpha_s] = F[\alpha_s]$  and since  $\{[\alpha_s]: S \text{ is a 1-simplex of } X'' \text{ generates } \pi_1(X, x_0) \text{ we have } f_* = F.$ 

Notes

### III. Claim $\phi$ is an Injection

Suppose  $f_* = g_*$ . As in Part 2), let T be a maximal tree. f/T is homotopic (rel  $x_0$ ) to a map which sends T to  $x_0$  (just use the retraction of T to  $x_0$ ). Hence, by the homotopy extension property, f is homotopic (rel  $x_0$ ) to a map f' with f'(T) =  $x_0$ . Therefore, we can assume  $f(T) = g(T) = x_0$ . In particular, for each 1-simplex s, f/s and g/s are loops at  $x_0$ . Since  $f_*=g_*$  this means f/s is homotopic to g/s (rel  $x_0$ ). Thus, for each 2—simplex  $\Delta$  we have a map H:  $\delta(\Delta xI) \rightarrow X$  where  $H/\Delta x0 = f/\Delta$ ,  $H/\Delta xI = f/\Delta$  and for each 1-simplex s in  $\delta\Delta$ , H/SxI is a homotopy from f/s to g/s. Since  $\pi 2(X, x_0) = 0$  and  $\delta(\Delta xI) = S^2$ , this map H can be extended to all of  $\Delta xI$ . Fitting together each of these H's we get a homotopy (rel  $x_0$ ) from f to g, i.e., [f] = [g].

Lemma 2 : If X is a closed 2—manifold and  $h:(X,x_0) \to (X,x_0)$  is a homeomorphism which is homotopic to the identity (rel  $x_0$ ) then h is isotopic to the identity (rel  $x_0$ ).

*Proof.* This is a special case of Theorem 6.3 of [1].

Lemma 3. If X is a closed 2—manifold with  $x_0 \in X$  and G is an automorphism of  $\pi 1(X,x_0)$ , then there exists a homeomorphism h:  $(X,x_0) + (X,x_0)$  with  $h_* = G$ .

*Proof.* This result is proved in [4].

Let  $H(X,x_0)$  denote the group of isotopy classes (rel  $x_0$ ) in the space of all homeomorphisms of  $(X,x_0)$  onto  $(X,x_0)$ .

By Lemma 2 the function from  $H(X,x_0)$  to  $[(X,x_0)]$  which sends the isotopy class (rel  $x_0$ ) of a homeomorphism to its homotopy class (rel  $x_0$ ) is an injection. By Lemma 1 the composition  $H(X,x_0) \rightarrow [(X,x_0), (X,x_0)] \rightarrow \text{End } \pi 1(X,x_0)$  is a monomorphism of the group  $H(X,x_0)$  onto a subgroup of Aut  $\pi 1(X,x_0)$ . Finally Lemma 3 shows that this monomorphism is an isomorphism onto Aut  $\pi_1(X,x_0)$ . Thus the proof of the theorem reduces to showing that H(Y) is isomorphic to  $H(X,x_0)$ . This is done by the following lemma. <u>Lemma 4</u>. If Y is a compact, connected 2—dimensional manifold with one boundary component and X is the closed 2—manifold obtained by sewing a disk to the boundary of Y and  $x_0 \in X$ -Bd(Y)then  $H(X,x_0) = H(Y)$ .

<u>Proof.</u> This result is a special case of Theorem 6 of [6].

Notes

<u>Remark 3</u>. If Y is a Moebius band or disk, then H(Y) is not isomorphic to Aut  $\pi 1(X,x0)$ . In the case Y is a disk, so that  $X = S^2$ , we have Aut  $\pi 1(X,x_0) = 1$  while  $H(Y) = Z_2$  (see Theorem 4.2 of [4]). In the case Y is a Moebius band, so that  $X = P^2$ , we have Aut  $\pi 1$ ,  $(X,x_0) = 1$  while  $H(Y) = Z_2$  (see Theorem 8.1 of [4]).

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## Improved Class of Ratio -Cum- Product Estimators of Finite Population Mean in two Phase Sampling

By Waikhom Warseen Chanu & B. K. Singh

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Abstract- In the present study, we have proposed a class of ratio-cum-product estimators for estimating finite population mean of the study variable in two phase sampling. The bias and mean square error of the proposed estimator have been obtained. The asymptotically optimum estimator (AOE) in this class has also been identified along with its approximate bias and mean square error. Comparison of the proposed class of estimators with other estimators is also worked out theoretically to demonstrate the superiority of the proposed estimator over the other estimators.

Keywords: ratio-cum-product, mean, two phase sampling, asymptotically optimum estimator, bias, mean square error.

GJSFR-F Classification : 62D05

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# Improved Class of Ratio -Cum- Product Estimators of Finite Population Mean in two Phase Sampling

Waikhom Warseen Chanu  $^{\alpha}$  & B. K. Singh  $^{\sigma}$ 

Abstract- In the present study, we have proposed a class of ratio-cum-product estimators for estimating finite population mean  $\overline{Y}$  of the study variable y in two phase sampling. The bias and mean square error of the proposed estimator have been obtained. The asymtotically optimum estimator (AOE) in this class has also been identified along with its approximate bias and mean square error. Comparison of the proposed class of estimators with other estimators is also worked out theoretically to demonstrate the superiority of the proposed estimator over the other estimators. *Keywords: ratio-cum-product, mean, two phase sampling, asymptotically optimum estimator, bias, mean square error.* 

### I. INTRODUCTION

The literature on survey sampling describes a great variety of techniques for using auxiliary information in order to obtain improved estimators for estimating some most common population parameter such as population total, population mean, population proportion, population ratio etc. More often we are interested in the estimation of the mean of a certain characteristic of a finite population on the basis of a sample taken from the population following a specified sampling procedure.

Use of auxiliary information has shown its significance in improving the efficiency of estimators of unknown population parameters. Cochran (1940) used auxiliary information in the form of population mean of auxiliary variate at estimation stage for the estimation of population parameters when study and auxiliary variate are positively correlated. In case of negative correlation between study variate and auxiliary variate, Robson (1957) defined product estimator for the estimation of population mean which was revisited by Murthy (1967). Ratio estimator performs better than simple mean estimator in case of positive correlation between study variate and auxiliary variate.

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For further discussion on ratio cum product estimator, the reader is referred to Singh (1967), Shah and Shah (1978), Singh and Tailor (2005), Tailor and Sharma (2009), Tailor and Sharma (2009), Sharma and Tailor (2010), Choudhury and Singh (2011) etc.

When the population mean  $\bar{X}$  of the auxiliary variable x is unknown before start of the survey, it is estimated from a preliminary large sample on which only the auxiliary characteristic x is observed. The value of X in the estimator is then replaced by its estimate. After then a smaller second-phase sample of the variate of interest (study variate) y is then taken. This technique is known as double sampling or two-phase sampling. Neyman (1938) was the first to give the concept of double sampling in connection with collecting information on the strata sizes in a stratified sampling Consider a finite population  $U = (u_1, u_2, u_3, ..., u_N)$  of size N units, y and x are the study and auxiliary variate respectively. When the population mean  $\bar{X}$  of x is not known, a first phase sample of size  $n_1$  is drawn from the population on which only the x characteristic is measured in order to furnish a good estimate of  $\bar{X}$ . After then a second-phase sample of size n ( $n < n_1$ ) is drawn on which both the variates y and xare measured.

The usual ratio and product estimators in double sampling are:

$$\bar{Y_R}^d = \bar{y}\frac{\bar{x_1}}{\bar{x}}$$

and

$$\bar{Y_P}^d = \bar{y}\frac{\bar{z}}{\bar{z_1}}$$

where  $\bar{x}, \bar{y}$  and  $\bar{z}$  are the sample mean of x, y and z respectively based on the sample of size n out of the population N units and  $\bar{y} = \frac{1}{n} \sum_{i=1}^{n} y_i$ ,  $\bar{x_1} = \frac{1}{n_1} \sum_{i=1}^{n_1} x_i$  and  $\bar{z_1} = \frac{1}{n_1} \sum_{i=1}^{n_1} z_i$  denote the sample mean of x and z based on the first- phase sample of the size  $n_1$ .

Singh (1967) improved the ratio and the product methods of estimation by studying the ratio cum product estimator for estimating  $\bar{Y}$  as

Notes

$$\bar{Y}_{RP} = \bar{y} \left( \frac{\bar{X}}{\bar{x}} \frac{\bar{z}}{\bar{Z}} \right)$$

Motivated by Singh (1967), Choudhury and Singh (2011) proposed a modified class of ratio cum product type of estimator for estimating population mean  $\bar{Y}$  as

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$$\bar{Y}_{RP}^{(\alpha)} = \bar{y} \left( \frac{\bar{X}}{\bar{x}} \frac{\bar{z}}{\bar{Z}} \right)^{\alpha}$$

Motivated by ?) and as an extension to the work of Choudhury and Singh (2011), we have developed an improved class of ratio-cum-product estimators in double sampling to estimate the population mean  $\bar{Y}$  theoretically and studied the properties of the proposed estimator.

#### II. The Proposed Estimator

The proposed improved class of ratio-cum-product estimators of population mean  $\bar{Y}$  in two-phase sampling is given as

$$\bar{Y}_{RP}^{w(d)} = \bar{y} \left( \frac{\bar{x_1}}{\bar{x}} \cdot \frac{\bar{z}}{\bar{z_1}} \right)^{\alpha}.$$
(1)

where  $\alpha$  is a suitably chosen constant.

To obtain the bias and MSE of  $\bar{Y}_{RP}^{w(d)}$  to the first degree of approximation, we write

$$e_{0} = (\bar{y} - \bar{Y}) / \bar{Y}, \quad e_{1} = (\bar{x} - \bar{Y}) / \bar{X}, \quad e_{2} = (\bar{x}_{1} - \bar{X}) / \bar{X}, \quad e_{3} = (\bar{z} - \bar{Z}) / \bar{Z},$$
$$e_{4} = (\bar{z}_{1} - \bar{Z}) / \bar{Z}.$$

Expressing  $Y_{RP}^{\overline{w}(d)}$  in terms of e's and neglecting higher power of e's, we have

$$Y_{RP}^{\bar{w}(d)} = \bar{Y} \left(1 + e_0\right) \left\{ \left(1 + e_2\right) \left(1 + e_1\right)^{-1} \left(1 + e_3\right) \left(1 + e_4\right)^{-1} \right\}^{\alpha}.$$

Assuming the sample size to be large enough so that  $|e_1| < 1$ ,  $|e_4| < 1$  and expanding  $(1 + e_1)^{-1}$ ,  $(1 + e_4)^{-1}$  in powers of  $e_1$ ,  $e_4$ , multiplying out and neglecting higher powers of e's, we have

$$Y_{RP}^{\bar{w}(d)} = \bar{Y} (1 + e_0) \left[ 1 - \left\{ e_1 - e_2 - e_3 + e_4 - e_1^2 - e_4^2 + e_1 e_2 - e_1 e_4 + e_2 e_4 + e_1 e_3 - e_2 e_3 + e_3 e_4 \right\} \right]^{\alpha} + e_2 e_4 + e_1 e_3 - e_2 e_3 + e_3 e_4 - \frac{e_1^2}{2} + \frac{e_2^2}{2} + \frac{e_3^2}{2} - \frac{e_4^2}{2} + e_0 e_1 - e_0 e_2 + e_0 e_4 - e_0 e_3 \right) + \alpha^2 \left( \frac{e_1^2}{2} + \frac{e_2^2}{2} + \frac{e_3^2}{2} + \frac{e_4^2}{2} - e_1 e_2 - e_1 e_3 + e_1 e_4 + e_2 e_3 - e_2 e_4 - e_3 e_4 \right) \right]$$
(2)

The following two cases will be considered separately.

**Case I**: When the second phase sample of size n is subsample of the first phase of size  $n_1$ .

Case II: when the second phase sample of size n is drawn independently of the first phase sample of size  $n_1$ .

#### CASE I

III. Bias, mse and Optimum Value of  $\, ar{Y} {}^{wd}_{RP} \,$  in Case I

In this case, we have

$$E(e_{0}) = E(e_{1}) = E(e_{2}) = E(e_{3}) = E(e_{4}) = 0;$$

$$E(e_{0}^{2}) = \left(\frac{1-f}{n}\right)C_{Y}^{2}; \quad E(e_{1}^{2}) = \left(\frac{1-f}{n}\right)C_{X}^{2}; \quad E(e_{3}^{2}) = \left(\frac{1-f}{n}\right)C_{Z}^{2};$$

$$E(e_{2}^{2}) = E(e_{1}e_{2}) = \left(\frac{1-f^{*}}{n}\right)C_{X}^{2}; \quad E(e_{4}^{2}) = E(e_{3}e_{4}) = \left(\frac{1-f^{*}}{n}\right)C_{Z}^{2};$$

$$E(e_{0}e_{1}) = \left(\frac{1-f}{n}\right)\rho_{YX}C_{Y}C_{X}; \quad E(e_{0}e_{2}) = \left(\frac{1-f^{*}}{n}\right)\rho_{YX}C_{Y}C_{X};$$

$$E(e_{0}e_{3}) = \left(\frac{1-f}{n}\right)\rho_{YZ}C_{Y}C_{Z}; \quad E(e_{0}e_{4}) = \left(\frac{1-f^{*}}{n}\right)\rho_{YZ}C_{Y}C_{Z};$$

$$E(e_{1}e_{3}) = \left(\frac{1-f}{n}\right)\rho_{XZ}C_{X}C_{Z};$$

$$E(e_{1}e_{4}) = E(e_{2}e_{3}) = E(e_{2}e_{4}) = \left(\frac{1-f^{*}}{n}\right)\rho_{XZ}C_{X}C_{Z} \qquad (3)$$

where  $f = \frac{n}{N}$  is the sampling fraction,  $f^* = \frac{n_1}{N}$ ,  $C_X^2 = \frac{S_X^2}{X^2}$ ,  $C_Y^2 = \frac{S_Y^2}{Y^2}$ ,  $C_Z^2 = \frac{S_Z^2}{Z^2}$ .

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Taking expectations on both the sides and using the results of (3) in (2), we get the bias of  $\bar{Y}_{RP}^{w(d)}$  as

$$B\left(\bar{Y}_{RP}^{w(d)}\right)_{I} = \bar{Y}\left(\frac{1-f_{1}}{2n}\right) \left[\alpha^{2}\left\{C_{X}^{2}+C_{Z}^{2}\left(1-2K_{XZ}\right)\right\}+\alpha\left\{C_{X}^{2}\left(1-2K_{YX}\right)\right\}\right]$$
(4)

$$-C_Z^2 \left(1 - 2K_{YZ}\right) \Big\} \Big]$$
$$= \bar{Y} \left(\frac{1 - f_1}{2n}\right) \left[\alpha^2 K_3 + \alpha \left(K_1 - K_2\right)\right]$$
(5)

where  $f_1 = \frac{n}{n_1}$ ,  $K_{YX} = \rho_{YX} \frac{C_Y}{C_X}$ ,  $K_1 = C_X^2 (1 - 2K_{YX})$ ,  $K_2 = C_Z^2 (1 - 2K_{YZ})$ ,  $K_3 = C_X^2 + C_Z^2 (1 - 2K_{XZ})$ .

Now from equation (2), we have

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$$\bar{Y}_{RP}^{w(d)} - \bar{Y} = \bar{Y} \left[ e_0 - \alpha \left( e_1 - e_2 - e_3 + e_4 \right) \right].$$

Squaring both the sides and taking expectations in the above equation and using the results of (3), we get the mean square error of  $\bar{Y}_{RP}^{w(d)}$  to the first degree of approximation as

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{I} = \bar{Y}^{2}\left(\frac{1-f}{n}\right)C_{Y}^{2} + \bar{Y}^{2}\left(\frac{1-f_{1}}{n}\right)\left[\alpha^{2}\left\{C_{X}^{2} + C_{Z}^{2}\left(1-2K_{XZ}\right)\right\} - 2\alpha\left(C_{YX} - C_{YZ}\right)\right]$$
$$= \bar{Y}^{2}\left(\frac{1-f}{n}\right)C_{Y}^{2} + \bar{Y}^{2}\left(\frac{1-f_{1}}{n}\right)\left[\alpha^{2}K_{3} - 2\alpha S_{1}\right]$$
(6)

where  $S_1 = C_{YX} - C_{YZ}, C_{YZ} = \rho_{YZ}C_YC_Z, C_{YX} = \rho_{YX}C_YC_X$ 

Differentiating  $M\left(\bar{Y}_{RP}^{w(d)}\right)$  w.r.t  $\alpha$  and equating to zero, we get

$$\alpha = \frac{S_1}{K_3}.\tag{7}$$

Now putting the optimum value of  $\alpha$  from (7) in the proposed estimator (1), we get the asymptotically optimum estimator(AOE) as

$$\left(\bar{Y}_{RP}^{w(d)}\right)_{I(opt)} = \bar{y} \left(\frac{\bar{x}_1}{\bar{x}} \frac{\bar{z}}{\bar{z}_1}\right)^{I\alpha(opt)}$$

Therefore, after putting the value of  $\alpha$  in (4) and (6), we obtain the optimum bias and MSE of  $\bar{Y}_{RP}^{w(d)}$  respectively as

$$B\left(\bar{Y}_{RP}^{w(d)}\right)_{I\alpha(opt)} = \bar{Y}\left(\frac{1-f_1}{2n}\right)\frac{S_1}{K_3}\left[K_1' - K_2'\right]$$
Notes

where

$$K'_1 = C_X^2 (1 - K_{YX}), \quad K'_2 = C_Z^2 (1 - K_{YZ}).$$

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{I\alpha(opt)} = \bar{Y}^2\left(\frac{1-f}{n}\right)C_Y^2 - \bar{Y}^2\left(\frac{1-f_1}{n}\right)\frac{S_1^2}{K_3}.$$
(8)

**Remark 1** For  $\alpha = 1$ , the estimator reduces to ratio cum product estimator in double sampling. The bias and MSE of  $\bar{Y}_{RP}^{(d)}$  are obtained by putting  $\alpha = 1$  in relation (4) and (6) as follows

$$B\left(\bar{Y}_{RP}^{(d)}\right)_{I} = \bar{Y}\left(\frac{1-f}{n}\right)\left[K_{1}' - C_{XZ} + C_{YZ}\right]$$

$$\tag{9}$$

where

$$K_1' = C_X^2 \left( 1 - K_{YX} \right)$$

and 
$$MSE\left(\bar{Y}_{RP}^{(d)}\right)_{I} = \left(\frac{1-f}{n}\right)S_{Y}^{2} + \bar{Y}^{2}\left(\frac{1-f_{1}}{n}\right)\left(K_{1}+K_{4}\right)$$
 (10)

where

$$K_4 = C_Z^2 \left( 1 - 2K_{XZ} + 2K_{YZ} \right).$$

**Remark 2** For  $\alpha = 1$  and when the auxiliary variate z is not used, i.e if z is nonzero constant, the proposed estimator reduces to the usual ratio estimator in two phase sampling. The bias and mean square error of  $\bar{Y}_R^d$  can be obtained by putting  $\alpha = 1$  and omitting the terms of z in equation (4) and (6), respectively as

$$B\left(\bar{Y}_{R}^{d}\right)_{I} = \bar{Y}\left(\frac{1-f_{1}}{n}\right)K_{1}^{\prime}$$

and

$$MSE\left(\bar{Y}_{R}^{d}\right)_{I} = \bar{Y}\left(\frac{1-f}{n}\right)C_{Y}^{2} + \bar{Y}\left(\frac{1-f_{1}}{n}\right)K_{1}.$$
(11)

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**Remark 3** For  $\alpha = 1$  and when the auxiliary variate x is not used, i.e if x is non-zero, the proposed estimator reduces to the usual product estimator in two phase sampling. The bias and mean square error of  $\bar{Y}_P^d$  can be obtained by putting  $\alpha = 1$  and omitting the terms of x in equation (4) and (6), respectively as

Notes

$$B\left(\bar{Y}_{P}^{d}\right)_{I} = \bar{Y}\left(\frac{1-f_{1}}{n}\right)C_{YZ}$$

and 
$$M\left(\bar{Y}_P^d\right)_I = \bar{Y}^2\left(\frac{1-f}{n}\right)C_Y^2 + \bar{Y}^2\left(\frac{1-f_1}{n}\right)K_5$$
 (12)

where  $K_5 = C_X^2 (1 + 2K_{YX})$ .

#### IV. EFFICIENCY COMPARISONS

Compairison of the optimum proposed estimator  $\left(\bar{Y}_{RP}^{w(d)}\right)_{I\alpha(opt)}$ 

a) with sample mean per unit estimator  $\bar{y}$ 

The MSE of sample mean  $\bar{y}$  under SRSWOR sampling scheme is given by

$$V\left(\bar{y}\right) = \left(\frac{1-f}{n}\right)S_Y^2.$$
(13)

From equation (13) and (8), we observed

$$V(\bar{y}) - M\left(\bar{Y}_{RP}^{w(d)}\right)_{I\alpha I(opt)} = \frac{S_1^2}{K_3} > 0.$$
 (14)

if  $K_3 > 0$  i.e  $K_{XZ} < 1/2$ .

b) with ratio estimator in double sampling

From equation (11) and (8), we observed

$$M\left(\bar{Y}_{R}^{d}\right) - M\left[\bar{Y}_{RP}^{w(d)}\right]_{I\alpha(opt)} = \bar{Y}^{2}\left(\frac{1-f_{1}}{n}\right)\left(K_{1} + \frac{S_{1}^{2}}{K_{3}}\right) > 0.$$
(15)

if  $K_1 > 0, K_3 > 0$  i.e.  $K_{YX} < 1/2, K_{XZ} < 1/2$ .

#### c) with product estimator in double sampling

From (12) and (8), we observed

$$M\left(\bar{Y}_{P}^{d}\right) - MSE\left[\bar{Y}_{RP}^{w(d)}\right]_{I\alpha(opt)} = \bar{Y}^{2}\left(\frac{1-f_{1}}{n}\right)\left[K_{5} + \frac{S_{1}^{2}}{K_{3}}\right] > 0$$
(16)  
if  $K_{3} > 0, K_{5} > 0$  i.e.  $K_{XZ} < 1/2.$ 

#### d) with ratio cum product estimators in double sampling

From (10) and (8), we observed

$$M\left(\bar{Y}_{RP}^{d}\right) - M\left[\bar{Y}_{RP}^{w(d)}\right]_{I\alpha(opt)} = \bar{Y}^{2}\left[K_{1} + K_{5} + \frac{S_{1}^{2}}{K_{3}}\right] > 0$$
(17)

if  $K_1 > 0, K_3 > 0, K_5 > 0$  i.e.  $K_{YX} < 1/2, K_{XZ} < 1/2, K_{XZ} - K_{YZ} < 1/2$ .

Now we state the theorem

**Theorem 4** To the first degree of approximation, the proposed class of estimators  $\bar{Y}_{Rd}^{w(d)}$  under the optimality (7) is consider to be more efficient than  $\bar{Y}_{R}^{d}$ ,  $\bar{Y}_{P}^{d}$ ,  $\bar{Y}_{RP}^{d}$ ,  $\bar{Y}_{RQ}^{d}$ ,  $\bar{Y}_{RQ}^{$ 

V. Bias, mse and Optimum Value of  $\ \bar{Y}_{RP}^{w(d)}$  in Case II In this case, we have

 $E(e_0) = E(e_1) = E(e_2) = E(e_3) = E(e_4) = 0;$  $E(e_2^2) = \left(\frac{1-f}{2}\right)C_2^2; \quad E(e_2^2) = \left(\frac{1-f}{2}\right)C_2^2;$ 

$$E\left(e_{0}^{2}\right) = \left(\frac{1-f}{n}\right)C_{Y}^{2}; \quad E\left(e_{1}^{2}\right) = \left(\frac{1-f}{n}\right)C_{X}^{2}; \quad E\left(e_{3}^{2}\right) = \left(\frac{1-f}{n}\right)C_{Z}^{2};$$

$$E\left(e_{2}^{2}\right) = E\left(e_{1}e_{2}\right) = \left(\frac{1-f^{*}}{n}\right)C_{X}^{2}; \quad E\left(e_{4}^{2}\right) = E\left(e_{3}e_{4}\right) = \left(\frac{1-f^{*}}{n}\right)C_{Z}^{2};$$

$$E\left(e_{0}e_{1}\right) = \left(\frac{1-f}{n}\right)\rho_{YX}C_{Y}C_{X}; \quad E\left(e_{0}e_{3}\right) = \left(\frac{1-f}{n}\right)\rho_{YZ}C_{Y}C_{Z};$$

$$E\left(e_{1}e_{3}\right) = \left(\frac{1-f}{n}\right)\rho_{XZ}C_{X}C_{Z}; \quad E\left(e_{2}e_{4}\right) = \left(\frac{1-f^{*}}{n}\right)\rho_{XZ}C_{X}C_{Z};$$

$$E\left(e_{0}e_{2}\right) = E\left(e_{0}e_{4}\right) = E\left(e_{1}e_{2}\right) = E\left(e_{1}e_{4}\right) = E\left(e_{3}e_{2}\right) = E\left(e_{3}e_{4}\right) = 0. \quad (18)$$

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Notes

Taking expectations in (2) and using the results of (18), we get the bias of  $\bar{Y}_{RP}^{w(d)}$  to the first degree of approximation as

$$B\left(\bar{Y}_{RP}^{w(d)}\right)_{II} = \bar{Y}\left[\alpha^{2}N_{1}K_{3} + \alpha\left\{f''\left(C_{X}^{2} - C_{Z}^{2}\right) - f'S_{1}\right\}\right].$$

Notes

where

$$f' = \frac{1-f}{n}, \quad f'' = \frac{1-f_1}{2n}, \quad N_1 = \frac{1}{2} \left( \frac{1}{n} + \frac{1}{n_1} - \frac{2}{N} \right).$$

Squaring and taking expectations in both the sides of (2) and using the results of (18), we obtain the MSE of  $\left(\bar{Y}_{RP}^{w(d)}\right)$  to the first degree of approximation as

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{II} = \bar{Y}^2 f' C_Y^2 + \bar{Y}^2 2 \left[\alpha^2 N_1 K_3 - \alpha f' S_1\right].$$
 (19)

Minimization of (19) is obtained with optimum value of  $\alpha$  as

$$\alpha = \frac{f'S_1}{2N_1K_3} = \alpha_{II(opt)}.$$
(20)

Substituting the value of  $\alpha$  from (20) in (1) gives the AOE of (1) as

$$\left\{\bar{Y}_{RP}^{w(d)}\right\}_{II(opt)} = \bar{y} \left(\frac{\bar{x_1}}{\bar{x}} \cdot \frac{\bar{z}}{\bar{z_1}}\right)^{\alpha_{II(opt)}}.$$
(21)

Thus, the resulting bias and MSE of (21) are, respectively given as

$$B\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)} = \bar{Y}\frac{f'S_1}{2N_1K_3}\left[f''C_1 - \frac{f'S_1}{2N_1K_1}\right]$$

where

$$C_1 = C_X^2 - C_Z^2$$
 and

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)} = \bar{Y}^2 f' C_Y^2 - \bar{Y}^2 \frac{f'^2 S_1^2}{2N_1 K_3}$$

For simplicity, we assume that the population size N is large enough as compared to the sample sizes n and  $n_1$  so that the finite population correction (FPC) terms 1/Nand 2/N are ignored. Ignoring the FPC in (19), the MSE of  $\left(\bar{Y}_{RP}^{w(d)}\right)_{II}$  reduces to

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{II} = \bar{Y}\frac{C_{Y}^{2}}{n} + \bar{Y}^{2}\left[\alpha^{2}\frac{1}{2}\left(\frac{1}{n} + \frac{1}{n_{1}}\right)K_{3} - \alpha\frac{S_{1}}{n}\right]$$

which is minimized for

$$\alpha = \frac{n_1 S_1}{(n+n_1) K_3} = \alpha^*_{II(opt)} \quad (\text{say})$$
(22)

Notes

Substituting the value of  $\alpha$  from (22) in (1), we obtained AOE of (1) as

$$\left(\bar{Y}_{RP}^{w(d)}\right)^* = \bar{Y}\left[\frac{x}{x_1}, \frac{z_1}{z}\right]^{\alpha^*_{II(opt)}}$$

Therefore, the resulting MSE of  $\left(\bar{Y}_{RP}^{w(d)}\right)^*$  is

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*} = \bar{Y}^{2} \frac{C_{Y}^{2}}{n} - \bar{Y}^{2} \frac{n_{1}S_{1}^{2}}{n\left(n+n_{1}\right)K_{3}}.$$
(23)

**Remark 5** For  $\alpha = 1$ , the proposed estimator reduces to ratio cum product estimator in double sampling and MSE is given as

$$M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*} = \bar{Y}^{2}\frac{C_{Y}^{2}}{n} + \bar{Y}^{2}2\left[\frac{1}{2}\left(\frac{1}{n} + \frac{1}{n_{1}}\right)K_{3} - \frac{S_{1}}{n}\right].$$
 (24)

Ignoring the FPC, the variance of  $\bar{y}$  under SRSWOR is given by

$$V(y)_{II} = \bar{Y}^2 \frac{C_Y^2}{n}.$$
 (25)

and the MSE of  $(\bar{Y}_{Rd})_{II}$  and  $(\bar{Y}_{Pd})_{II}$  are given by

$$M\left(\bar{Y}_{Rd}\right)_{II} = \bar{Y}^2 \frac{C_Y^2}{n} + \bar{Y}^2 \left[ C_X^2 \left( \frac{2}{n} - \frac{1}{n_1} - \frac{2}{n} K_{YX} \right) \right]$$
(26)

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and

$$\left(\bar{Y}_{Pd}\right)_{II} = \bar{Y}^2 \frac{C_Y^2}{n} + \bar{Y}^2 \left[\frac{C_X^2}{n} + \frac{1}{2}\left(\frac{1}{n} - \frac{1}{n_1}\right)C_Z^2 + \frac{2C_{YZ}}{n}\right]$$
(27)

respectively.

 $N_{\mathrm{otes}}$ 

#### VI. EFFICIENCY COPMPARISONS

Compairison of the optimum proposed estimator  $\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*}$ 

a) with sample mean per unit estimator

From (25) and (23), we observed that

$$V(\bar{y})_{II} - M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*} = \bar{Y}^{2} \frac{n_{1}S_{1}^{2}}{n(n+n_{1})K_{3}} > 0$$
(28)

if

## $K_3 > 0$ i.e $K_{XZ} < 1/2$ .

#### b) with ratio estimator in double sampling

From (26) and (23), we observed that

$$M\left(\bar{Y}_{Rd}\right)_{II} - M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*} = \bar{Y}^{2}\left[\frac{n_{1}S_{1}^{2}}{n\left(n+n_{1}\right)K_{3}} + C_{X}^{2}P\right] > 0$$
(29)

if

if

$$K_3 > 0, P > 0$$
 i.e  $K_{XZ} < 1/2, K_{YX} < 1 - \frac{n}{2n_1}$ 

where

$$P = \frac{2}{n} \left( 1 - K_{YX} \right) - \frac{1}{n_1}$$

c) with product estimator in double sampling

From (27) and (23), we observed that

$$M\left(\bar{Y}_{Pd}\right)_{II} - M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*} = \bar{Y}^{2}\left[Q - \frac{C_{Z}^{2}}{2n_{1}} + \frac{n_{1}S_{1}^{2}}{n\left(n+n_{1}\right)K_{3}}\right] > 0 \qquad (30)$$

$$K_3 > 0$$
 i.e  $K_{XZ} < 1/2$ , where  $Q = C_X^2 + \frac{C_Z^2}{2} + 3C_{YX}$ .

#### d) with ratio cum product estimator in double sampling

From (24) and (23), we observed that

$$M\left(\bar{Y}_{RPd}\right)_{II} - M\left(\bar{Y}_{RP}^{w(d)}\right)_{II\alpha(opt)}^{*} = \bar{Y}^{2}\left[\left(\frac{1}{n} + \frac{1}{n_{1}}\right)K_{3} - \frac{S_{1}}{2n} + \frac{n_{1}S_{1}^{2}}{n\left(n+n_{1}\right)K_{3}}\right] > 0$$

$$K_{3} > 0 \text{ i.e } K_{XZ} < 1/2.$$
(31)

Notes

#### VII. CONCLUSION

We have developed an efficient class of ratio-cum-product estimators in two phase sampling. The comparative study shows that the proposed estimator  $\bar{Y}_{RP}^{w(d)}$  established their superiority over sample mean  $\bar{Y}$ , ratio estimator  $\bar{Y}_{R}^{d}$ , product estimator  $\bar{Y}_{P}^{d}$  and ratio-cum-product estimator  $\bar{Y}_{RP}^{d}$  in two-phase sampling under the given conditions. Hence from the resulting equation (14), (15), (16) and (17), we conclude that under the given conditions the proposed estimator is consider to be the best estimator.

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#### TECHNIQUES FOR WRITING A GOOD QUALITY RESEARCH PAPER:

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24. Never copy others' work: Never copy others' work and give it your name because if evaluator has seen it anywhere you will be in trouble.

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**28. Make colleagues:** Always try to make colleagues. No matter how sharper or intelligent you are, if you make colleagues you can have several ideas, which will be helpful for your research.

29. Think technically: Always think technically. If anything happens, then search its reasons, its benefits, and demerits.

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#### INFORMAL GUIDELINES OF RESEARCH PAPER WRITING

Key points to remember:

- Submit all work in its final form.
- Write your paper in the form, which is presented in the guidelines using the template.
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The introduction will be compiled from reference matter and will reflect the design processes or outline of basis that direct you to make study. As you will carry out the process of study, the method and process section will be constructed as like that. The result segment will show related statistics in nearly sequential order and will direct the reviewers next to the similar intellectual paths throughout the data that you took to carry out your study. The discussion section will provide understanding of the data and projections as to the implication of the results. The use of good quality references all through the paper will give the effort trustworthiness by representing an alertness of prior workings.

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To make a paper clear

· Adhere to recommended page limits

#### Mistakes to evade

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#### In every sections of your document

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- Fundamental goal
- To the point depiction of the research
- Consequences, including <u>definite statistics</u> if the consequences are quantitative in nature, account quantitative data; results of any numerical analysis should be reported
- Significant conclusions or questions that track from the research(es)

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- Present a justification. Status your particular theory (es) or aim(s), and describe the logic that led you to choose them.
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#### Approach:

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- If use of a definite type of tools.
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- Simplify details how procedures were completed not how they were exclusively performed on a particular day.
- If well known procedures were used, account the procedure by name, possibly with reference, and that's all.

#### Approach:

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#### What to keep away from

- Resources and methods are not a set of information.
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The principle of a results segment is to present and demonstrate your conclusion. Create this part a entirely objective details of the outcome, and save all understanding for the discussion.

The page length of this segment is set by the sum and types of data to be reported. Carry on to be to the point, by means of statistics and tables, if suitable, to present consequences most efficiently. You must obviously differentiate material that would usually be incorporated in a study editorial from any unprocessed data or additional appendix matter that would not be available. In fact, such matter should not be submitted at all except requested by the instructor.


Content

- Sum up your conclusion in text and demonstrate them, if suitable, with figures and tables.
- In manuscript, explain each of your consequences, point the reader to remarks that are most appropriate.
- Present a background, such as by describing the question that was addressed by creation an exacting study.
- Explain results of control experiments and comprise remarks that are not accessible in a prescribed figure or table, if appropriate.

• Examine your data, then prepare the analyzed (transformed) data in the form of a figure (graph), table, or in manuscript form. What to stay away from

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#### Approach

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- Give details all of your remarks as much as possible, focus on mechanisms.
- Make a decision if the tentative design sufficiently addressed the theory, and whether or not it was correctly restricted.
- Try to present substitute explanations if sensible alternatives be present.
- One research will not counter an overall question, so maintain the large picture in mind, where do you go next? The best studies unlock new avenues of study. What questions remain?
- Recommendations for detailed papers will offer supplementary suggestions.

#### Approach:

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- Submit to work done by specific persons (including you) in past tense.
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Introduction	Containing all background details with clear goal and appropriate details, flow specification, no grammar and spelling mistake, well organized sentence and paragraph, reference cited	Unclear and confusing data, appropriate format, grammar and spelling errors with unorganized matter	Out of place depth and content, hazy format
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Result	Well organized, Clear and specific, Correct units with precision, correct data, well structuring of paragraph, no grammar and spelling mistake	Complete and embarrassed text, difficult to comprehend	Irregular format with wrong facts and figures
Discussion	Well organized, meaningful specification, sound conclusion, logical and concise explanation, highly structured paragraph reference cited	Wordy, unclear conclusion, spurious	Conclusion is not cited, unorganized, difficult to comprehend
References	Complete and correct format, well organized	Beside the point, Incomplete	Wrong format and structuring

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