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# Mathematics and Decision Sciences 

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Dissemination Sinusoidal Waves

Discovering Thoughts, Inventing Future

Global Journal of Science Frontier Research: F mathematics \& Decision Sciences

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# The Axisymmetric Slow Viscous Flow About A Shear Stress Free Sphere 

By S. K. Sen, M. Kamran Chowdhury \& M. Jalal Ahammad<br>University of Chittagong, Bangladesh

Abstract- Harper's sphere theorem for the axisymmetric slow viscous flow exterior to a shear stress-free sphere is established in an alternative way and then given an extension of the theorem for the flow interior to the same sphere.

Keywords: harper's theorem, viscous flow, shear stress, circle theorem.
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# The Axisymmetric Slow Viscous Flow About A Shear Stress Free Sphere 

S. K. Sen ${ }^{\alpha}$, M. Kamran Chowdhury ${ }^{\circ}$ \& M. Jalal Ahammad ${ }^{\rho}$

Abstract- Harper's sphere theorem for the axisymmetric slow viscous flow exterior to a shear stress-free sphere is established in an alternative way and then given an extension of the theorem for the flow interior to the same sphere. Keywords: harper's theorem, viscous flow, shear stress, circle theorem.

## I. Introduction

In Harper [1] it is stated that sphere on which there is no shear stress are found as boundaries in slow viscous fluid flow in two important contexts. The earth's core is, to a good approximation, such a boundary for the convection in its mantle, and the surface of a gas bubble is such a boundary for the flow outside it. In the literature corresponding to Harper's sphere theorem [1] for the axisymmetrical slow viscous flow past a shear stress-free sphere, there are the sphere theorems for axisymmetrical potential flows outside or inside a rigid sphere due to Butler [2] in terms of Stokes stream function [3]. Again there are exterior sphere theorem due to Weiss [4] and the interior sphere theorem of Ludford et al. [5] each for a general irrotational motion of inviscid fluid, both being expressed in terms of the potential function. Furthermore, for axisymmetrical slow viscous fluid motion outside or inside a rigid sphere there are sphere theorems in terms of the Stokes stream function, which are due to Collins $[6,7]$.

The two dimensional analogue of Harper's theorem [1] referred above is the circle theorem due to Usha et al. [8] for the slow viscous flow past a shear free circular boundary. Relevantly, there is a circle theorem for potential flow past a circular boundary, which is due to Milne - Thomson [3, 9]. Further, in the two-dimensional viscous flow theory similar theorems are found in Avudainayagon et al. [10] and Sen [11] for solving the problems of slow viscous flow past a rigid circular boundary with shear stress.

Following Batchelor [12], we may that when a body of small size moves through fluid, it generates a flow problem which is important in a variety of physical contexts, such as setting of sediment in liquid and fall of mist droplets in air. The matter of great practical interest is the drag force exerted by the fluid on the body. Except in a few simple bodies, such as spherical ones exact solutions for arbitrary body shapes in viscous fluid motions are, in general, not found in the literature.

Our main interest lies in studying the viscous flow about arbitrary rigid bodies which are shear stress-free. With this object in mind first we derive Harper's theorem for a shear stress-free sphere by an analytic technique; and this is done in section 3. For this purpose we need some relevant mathematical results, which are established in the following section.

[^0]
## II. Mathematical Theory

In this section, we first derive Stokes' equation in terms of the Stokes stream function $\psi=\psi(r, \theta)$ for the axisymmetrical motion about axisymmetrical bodies, such as a rigid sphere and then the condition of no shear stress on the sphere, due to an axisymmetrical fluid motion.

When the inertia force in a steady viscous flow field is negligibly small, the Navier-Stokes equations, governing of the flow become

$$
\begin{equation*}
\operatorname{grad} p=\mu \nabla^{2} q \tag{1}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{div} q=0 \tag{2}
\end{equation*}
$$

where $q$ is the fluid velocity, $p$ the pressure, and $\mu$ the coefficient of viscosity of the fluid.

In the present paper it is convenient to derive the scalar expression of the vector equations (1) and (2) in spherical polar coordinates ( $r, \theta, \phi$ ) and then to express them in terms of Stokes stream function $\psi=\psi(r, \theta)$ as a dependent variable for the differential equations for an axis-symmetrical slow fluid motion in a viscous fluid.

The scalar expressions of equations (1) and (2) can be derived with the help of the relevant results in Batchelor [12, appendix 2] as

$$
\begin{align*}
& -\frac{\partial p}{\partial r}=\mu\left\{\nabla^{2} q_{r}-\frac{2 q_{r}}{r^{2}}-\frac{2}{r^{2} \sin \theta} \frac{\partial}{\partial \theta}\left(q_{\theta} \sin \theta\right)-\frac{2}{r^{2} \sin \theta} \frac{\partial q_{\phi}}{\partial \phi}\right\}  \tag{3}\\
& \quad \frac{1}{r} \frac{\partial p}{\partial \theta}=\mu\left\{\nabla^{2} q_{\theta}+\frac{2}{r^{2}} \frac{\partial q_{r}}{\partial \theta}-\frac{q_{\theta}}{r^{2} \sin ^{2} \theta}-\frac{2 \cos \theta}{r^{2} \sin ^{2} \theta} \frac{\partial q_{\phi}}{\partial \phi}\right\}  \tag{4}\\
& \frac{1}{r \sin \theta} \frac{\partial p}{\partial \phi}=\mu\left\{\nabla^{2} q_{\phi}+\frac{2}{r^{2} \sin \theta} \frac{\partial q_{r}}{\partial \phi}+\frac{2 \cos \theta}{r^{2} \sin ^{2} \theta} \frac{\partial q_{\theta}}{\partial \phi}-\frac{q_{\phi}}{r^{2} \sin ^{2} \theta}\right\} \tag{5}
\end{align*}
$$

and

$$
\begin{equation*}
\frac{\partial\left(r^{2} q_{r}\right)}{\partial r}+\frac{r}{\sin \theta} \frac{\partial}{\partial \theta}\left(q_{\theta} \sin \theta\right)+\frac{r}{\sin \theta} \frac{\partial q_{\phi}}{\partial \phi}=0 \tag{6}
\end{equation*}
$$

If the fluid motion is axisymmetrical about the z -axis, the fluid velocity everywhere in the flow field becomes independent of the azimuthal coordinate $\phi$ and the azimuthal velocity component $q_{\phi}=0$. Thus equations (3) to (5) appear as

$$
\begin{gather*}
-\frac{\partial p}{\partial r}=\mu\left\{\nabla^{2} q_{r}-\frac{2 q_{r}}{r^{2}}-\frac{2}{r^{2} \sin \theta} \frac{\partial}{\partial \theta}\left(q_{\theta} \sin \theta\right)\right\}  \tag{7}\\
\frac{1}{r} \frac{\partial p}{\partial \theta}=\mu\left\{\nabla^{2} q_{\theta}+\frac{2}{r^{2}} \frac{\partial q_{r}}{\partial \theta}-\frac{q_{\theta}}{r^{2} \sin ^{2} \theta}\right\} \tag{8}
\end{gather*}
$$

And

$$
\begin{equation*}
\frac{\partial}{\partial r}\left(r^{2} q_{r}\right)+\frac{r}{\sin \theta} \frac{\partial}{\partial \theta}\left(q_{\theta} \sin \theta\right)=0 \tag{9}
\end{equation*}
$$

where $\nabla^{2}$ is the three - dimensional Laplace's operator.

These equations can be further simplified with the help of the formulae for velocity components $q_{\mathrm{r}}$ and $q_{\theta}$, defined in terms of Stokes stream function $\psi=\psi(r, \theta)$ for the axisymmetrical fluid motion; and these formulae are

$$
\begin{equation*}
q_{r}=-\frac{1}{r^{2} \sin \theta} \frac{\partial \psi}{\partial \theta} \text { and } q_{\theta}=\frac{1}{r \sin \theta} \frac{\partial \psi}{\partial r}, \tag{10}
\end{equation*}
$$

which clearly satisfy the mass conservation equation (9).
Now eliminating $q_{\mathrm{r}}$ and $q_{\theta}$ from equations (7) and (8) by using the relations (10), yields

$$
\begin{gather*}
\frac{\partial p}{\partial r}=\mu \frac{1}{r^{2} \sin \theta}\left[\frac{\partial^{2} \psi}{\partial \theta \partial r}+\frac{1}{r^{2}} \frac{\partial}{\partial \theta}\left(-\cot \theta \frac{\partial \psi}{\partial \theta}+\frac{\partial^{2} \psi}{\partial \theta^{2}}\right)\right]  \tag{11}\\
\frac{1}{r} \frac{\partial p}{\partial \theta}=\frac{1}{r^{2} \sin \theta}\left[r \frac{\partial^{3} \psi}{\partial r^{3}}-\frac{\sin \theta}{r}\left(\frac{\partial^{2}}{\partial r \partial \theta}\left(\operatorname{cosec} \theta \frac{\partial \psi}{\partial \theta}\right) \frac{2}{r^{2}} \frac{\partial}{\partial \theta}\left(\operatorname{cosec} \theta \frac{\partial \psi}{\partial \theta}\right)\right]\right. \tag{12}
\end{gather*}
$$

On using the operator, $E^{2}=\frac{\partial^{2}}{\partial r^{2}}+\frac{\sin \theta}{r^{2}} \frac{\partial}{\partial \theta}\left(\frac{1}{\sin \theta} \frac{\partial}{\partial \theta}\right)$ for treating the axisymmetrical fluid motion [3], two concise forms of equations (11) and (12) are easily obtained as

$$
\begin{align*}
\frac{\partial p}{\partial r} & =\mu \frac{1}{r^{2} \sin \theta} \frac{\partial}{\partial \theta}\left(E^{2} \psi\right)  \tag{13}\\
\frac{\partial p}{\partial \theta} & =-\mu \frac{1}{\sin \theta} \frac{\partial}{\partial r}\left(E^{2} \psi\right) \tag{14}
\end{align*}
$$

Next eliminating the pressure $p$ from these equations results in

$$
\begin{equation*}
E^{4} \psi=0, \tag{15}
\end{equation*}
$$

which is obtained by different methods in Milne-Thomson [3]. The last equation may be called Stokes equation for the stream function of axisymmetrical and slow viscous fluid motion. We note that the stream function $\psi=\psi(r, \theta)$ for a slow viscous fluid motion past a axisymmetrical rigid body must satisfy the differential equation (15).

A general solution for the stream function $\psi$ in spherical polar coordinates is given in [13]. For the convenience of the reference in our present study for a viscous flow past a shear stress-free sphere, we only quote here the relevant part of the general solution and this is

$$
\begin{equation*}
\psi(r, \theta)=\sum_{n=2}^{\infty}\left(a_{n} r^{n}+b_{n} r^{n+1}+c_{n} r^{n+2}+d_{n} r^{-n+3}\right) \phi_{n}(\xi) \tag{16}
\end{equation*}
$$

where $a_{\mathrm{n}}, b_{\mathrm{n}}, c_{\mathrm{n}}$ and $d_{\mathrm{n}}$ are arbitrary real constants, $\xi=\cos \theta$ and $\phi_{\mathrm{n}}(\xi)$ is the Gegenbauer function of the first kind defined by

$$
\begin{equation*}
\phi_{n}(\xi)=\frac{P_{n-2}(\xi)-P_{n}(\xi)}{2 n-1}, n \geq 2 \tag{17}
\end{equation*}
$$

where $P_{\mathrm{n}}(\xi)$ is the Legendre function of the first kind.

Now we are interested in deriving the condition for no shear stress on a rigid sphere in axisymmetrical viscous fluid motion. Here on the surface of a sphere $r=a$, out of the six components of the stress tensor [12, appendix 2] only three exist, which are

$$
\begin{gathered}
\sigma_{r r}=2 \mu\left(\frac{\partial q_{r}}{\partial r}\right)_{r=a} \\
\sigma_{\phi r}=2 \mu\left(\frac{1}{r \sin \theta} \frac{\partial q_{r}}{\partial \phi}+\frac{r}{2} \frac{\partial}{\partial r}\left(\frac{q_{\phi}}{r}\right)\right)_{r=a}
\end{gathered}
$$

We first evaluate the velocity components $q_{r}$ and $q_{\theta}$ by substituting Stokes' stream function (16) in the formulae (10) for the axisymmetrical flow past the shear stress- free sphere $r=a$.

On using the stress components (18), we then find that, on $r=a$, the normal stress $\sigma_{r r} \neq 0$, the shearing stress $\sigma_{\phi r}=0$ in the $\phi$ - direction, since $q_{\phi}=0$; and $q_{r}$ is independent of $\phi$ and the shear stress in the $\theta$ - direction $\sigma_{\theta r} \neq 0$. Finally, it is easy to calculate that

$$
\begin{equation*}
\sigma_{\theta r}=0 \text { on } r=a \quad \text { when } \frac{\partial}{\partial r}\left(\frac{1}{r^{2}} \frac{\partial \psi}{\partial r}\right)=0 \tag{19}
\end{equation*}
$$

and $q_{r}=0$ on $r=a$ when $\psi=0$.
Therefore, the results (19) and (20) are the required conditions for shear stress free-sphere $r=a$ for axisymmetrical fluid motion past the same sphere.

Our future aim is to solve the problems of axisymmetrical flows past arbitrary symmetrical body shapes which are shear stress free e.g. oblate and prolate spheroids, etc.

With this object in mind, we now present a relatively different analysis to establish the Harper's theorem [1] for the slow axisymmetrical viscous flow exterior to a shear stress-free sphere, and finally we also add an extension of the same theorem for the flow interior to the same sphere.

## iII. Harper's Theorem

In an unlimited incompressible viscous fluid there is a steady and slow axisymmetrical motion and the motion is characterized by the Stokes stream function $\psi_{0}=\psi_{0}(r, \theta)$, whose singularities are all at a distance greater than 'a' from the origin and $\psi_{0}(r, \theta) \sim O\left(r^{2}\right)$ near the origin. Then if a shear stress-free sphere is introduced into the flow, Stokes stream function for the new flow outside the same sphere become

$$
\begin{equation*}
\psi(r, \theta)=\psi_{0}(r, \theta)-\frac{r^{3}}{a^{3}} \psi_{0}\left(\frac{a^{2}}{r}, \theta\right) . \tag{21}
\end{equation*}
$$

Proof: Since the singularities of $\psi_{0}=\psi_{0}(r, \theta)$ are at a distance greater than ' $a$ ' from the origin, $\psi_{o}(r, \theta)$ is regular at the origin. Then we suppose $\psi_{o}$ in the absence of any boundary, has an expression of the from

$$
\begin{equation*}
\psi_{0}(r, \theta)=\sum_{n=2}^{\infty}\left(A_{n} r^{n}+B_{n} r^{n+2}\right) \phi(\xi) \tag{22}
\end{equation*}
$$

where $A_{n}, B_{n}$ are all known constants and $\phi_{n}(\xi)$ is the Gegenbauer function of, $\xi=\cos \theta$, of the first kind.

If a shear stress free sphere $r=a$ is now introduced into the viscous flow, the Stokes stream function for a possible new fluid motion must be obtained from the general expression (16), that is,

$$
\begin{equation*}
\psi(r, \theta)=\sum_{n=2}^{\infty}\left(A_{n} r^{n}+B_{n} r^{n+2}+C_{n} r^{-n+1}+D_{n} r^{-n+3}\right) \phi_{n}(\xi) \tag{23}
\end{equation*}
$$

where the last two terms constitute the perturbation stream function of the flow due to the presence of the sphere, and where $C_{n}$ and $D_{n}$ are the constants to be determined.
Here the conditions for the flow to be possible are on $r=a, \psi(r, \theta)=0$,
and on $r=a, \frac{\partial}{\partial r}\left(\frac{1}{r^{2}} \frac{\partial \psi}{\partial r}\right)=0$.
Thus, by using these conditions, we have

$$
\begin{align*}
& A_{n} a^{n}+B_{n} a^{n+2}+C_{n} a^{-n+1}+D_{n} a^{-n+3}=0  \tag{24}\\
& A_{n} n(n-3) a^{n-4}+B_{n}(n-1)(n+2) a^{n-2}+C_{n}(-n+1)(-n-2) a^{-n-3}+D_{n}(-n+3)(-n) \\
& a^{-n-1}=0 . \tag{25}
\end{align*}
$$

Solving (24) and (25) for $C_{n}$ and $D_{n}$, we obtain

$$
\begin{gather*}
C_{n}=-B_{n} a^{2 n+1},  \tag{26}\\
D_{n}=-A_{n} a^{2 n-3} . \tag{27}
\end{gather*}
$$

Next, we adopt the following analysis to obtain the result (21). Using the basic stream function (22) in (23) gives

$$
\begin{equation*}
\psi(r, \theta)=\psi_{o}(r, \theta)+\sum_{n=2}^{\infty}\left(C_{n} r^{-n+1}+D_{n} r^{-n+3}\right) \phi_{n}(\xi) . \tag{28}
\end{equation*}
$$

Substituting (26) and (27) in this expression yields

$$
\begin{equation*}
\psi(r, \theta)=\psi_{o}(r, \theta)-\sum_{n=2}^{\infty}\left(B_{n} a^{2 n+1} r^{-n+1}+A_{n} a^{2 n-3} r^{-n+3}\right) \phi_{n}(\xi) \tag{29}
\end{equation*}
$$

Now our object is to give the expression (29) a closed form and this is done as follows. From the expression (22) one gets

$$
\sum_{n=2}^{\infty}\left(A_{n} a^{2 n} r^{-n}+B_{n} a^{2 n+4} r^{-n-2}\right) \phi_{n}(\xi)=\psi_{0}\left(\frac{a^{2}}{r}, \theta\right)
$$

Multiplying both sides by $\frac{r^{3}}{a^{3}}$, gives

$$
\begin{equation*}
\sum_{n=2}^{\infty}\left(A_{n} a^{2 n-3} r^{-n+3}+B_{n} a^{2 n+1} r^{-n+1}\right) \phi_{n}(\xi)=\frac{a^{3}}{r^{3}} \psi_{0}\left(\frac{a^{2}}{r}, \theta\right) \tag{30}
\end{equation*}
$$

Finally, substituting (30) in (29) yields the Stokes stream function for the slow viscous fluid motion exterior to the shear stress-free sphere $r=a$ as

$$
\psi(r, \theta)=\psi_{o}(r, \theta)-\left(\frac{r^{3}}{a^{3}}\right) \psi_{o}\left(\frac{a^{2}}{r}, \theta\right)
$$

which is in agreement with Harper's result [1]. We then show that the perturbation velocity due to the last term in (3.14) vanishes at infinity. Since $\psi_{o}(r, \theta)$ is $O\left(r^{2}\right)$ near the origin the perturbation stream function $\left(\frac{r^{3}}{a^{3}}\right) \psi_{\rho}\left(\frac{a^{2}}{r}, \theta\right)$ is clearly $O(r)$ at infinity which implies a vanishing velocity at infinity. Hence the theorem is established.

## IV. Extension of Harper's Theorem

We now extend Harper's sphere theorem for the viscous flow exterior to a shear stress-free sphere, to case of the flow interior to the same sphere. This extension corresponds to the Butler's interior sphere theorem [2] for the axi-symmetric and irrotational inviscid fluid flow within a sphere.

## a) An Extension of Harper's Sphere Theorem

Let an axi-symmetric slow flow in an incompressible viscous fluid in the absence of rigid boundaries be characterized by Stokes steam $\psi_{0}=\psi_{o}(r, \theta)$, whose singularities are all at a distance less ' $a$ ' from the origin. Let $\psi_{0} \sim O\left(\frac{1}{r^{k}}\right), k \geq 1$ as $r \rightarrow \infty$. Now if a shear stress-free rigid sphere be introduced into the flow, the resultant flow interior to the sphere becomes

$$
\begin{equation*}
\psi=\psi_{o}(r, \theta)-\frac{r^{3}}{a^{3}} \psi_{o}\left(\frac{a^{2}}{r}, \theta\right) \tag{31}
\end{equation*}
$$

Proof : Since the singularities of the Stokes stream function $\psi_{o}(r, \theta)$ are all at a distance less than ' $a$ ' from the origin, the function is regular everywhere in the region outside the sphere $r=a$, i.e., the region $r \geq a$.
Therefore a relevant expansion of $\psi_{o}(r, \theta)$ must be an expansion of the from

$$
\begin{equation*}
\psi_{o}(r, \theta)=\sum_{n=2}^{\infty}\left(A_{n} \frac{1}{r^{n-1}}+B_{n} \frac{1}{r^{-n+3}}\right) \phi_{n}(\xi) \tag{32}
\end{equation*}
$$

where $A_{n}$ and $B_{n}$ are all known constants, and $\phi_{n}(\xi)$ is the Gegenbauer function of the first kind already referred above.

If the shear free rigid sphere $r=a$ now is introduced into the basic flow characterized by the stream function (32), the Stokes stream function for the disturbed fluid motion may be given by

$$
\begin{equation*}
\psi(r, \theta)=\sum_{n=2}^{\infty}\left(A_{n} \frac{1}{r^{n-1}}+B_{n} \frac{1}{r^{n-3}}+\mathrm{C}_{\mathrm{n}} \mathrm{r}^{\mathrm{n}}+\mathrm{D}_{\mathrm{n}} \mathrm{r}^{\mathrm{n}+2}\right) \phi_{n}(\xi) \tag{33}
\end{equation*}
$$

where the last two terms constitute the perturbation stream function with the undetermined constants $C_{n}$ and $D_{n}$.

First we determine the constants $C_{n}$ and $D_{n}$ as follows. On the shear stress-free sphere $r=a$, the Stokes stream function (33) must satisfy the boundary conditions on $r=a, \psi=0$,
and

$$
r=a, \frac{\partial}{\partial r}\left(\frac{1}{r^{2}} \frac{\partial \psi}{\partial r}\right)=0
$$

Using these boundary conditions one obtains

$$
\begin{gather*}
A_{n} a^{-n+1}+B_{n} a^{-n+3}+C_{n} a^{n}+D_{n} a^{n+2}=0  \tag{34}\\
A_{n}(n-1)(n+2) a^{-n-3}+B_{n} n(n-3) a^{-n-1}+C_{n} n(n-3) a^{n-4}+D_{n}(n-1)(n+2) a^{n-2}=0 \tag{35}
\end{gather*}
$$

Solving (34) and (35) for the values of $C_{n}$ and $D_{n}$, we get very simple results as

$$
\begin{equation*}
C_{n}=-B_{n} a^{-2 n+3} \text { and } D_{n}=-A_{n} a^{-2 n-1} . \tag{36}
\end{equation*}
$$

We now give a closed form of the stream function (33) in the following way. At the outset we note that the first two terms of (33) may be replaced by the stream function $\psi_{o}(r, \theta)$ referred to the expansion (32). Thus we have

$$
\begin{equation*}
\psi(r, \theta)=\psi_{o}(r, \theta)-\sum_{n=2}^{\infty}\left(B_{n} a^{-2 n+3} r^{n}+A_{n} a^{-2 n-1} r^{n+2}\right) \phi_{n}(\xi) . \tag{37}
\end{equation*}
$$

By using the relation (32) we at once have

$$
\begin{equation*}
\sum_{n=2}^{\infty}\left(A_{n} a^{-2 n-1} r^{n+2}+B_{n} a^{-2 n+3} r^{n}\right)=\frac{r^{3}}{a^{3}} \psi_{0}\left(\frac{a^{2}}{r}, \theta\right) \tag{38}
\end{equation*}
$$

Finally, the use of this relation in (37), yields the Stokes stream function in closed form for the flow within a shear stress-free sphere as

$$
\begin{equation*}
\psi(r, \theta)=\psi_{o}(r, \theta)-\left(\frac{r^{3}}{a^{3}}\right) \psi_{o}\left(\frac{a^{2}}{r}, \theta\right) . \tag{39}
\end{equation*}
$$

Next we show that the stream function $\psi(r, \theta)$ gives a finite value at the origin.
Since $\psi_{0}(r, \theta) \sim O\left(\frac{1}{r}\right)$ for large r , the last term on the right hand side of the stream function (39) is $O\left(r^{4}\right)$ near the origin so that the same term gives the vanishing velocity at the origin. Thus the theorem is established.
Example : A source and sink interior to a shear stress free sphere.
Let us consider, there be a source of strength $m$ at the point $A_{1}(-c, 0,0)$ and a sink of strength $-m$ at the point $\mathrm{A}_{2}(\mathrm{c}, 0,0)$ on the axis of symmetry, z-axis. The Stokes stream due to their combination is given by

$$
\begin{equation*}
\psi_{0}(r, \theta)=m \cos \theta_{1}-m \cos \theta_{2} . \tag{40}
\end{equation*}
$$

To find out the Stokes stream function for the flow within the sphere $r=a$, first we are to show that $\psi_{0}(r, \theta) \approx O\left(\frac{1}{r}\right)$ for large $r$.

Since the source and the sink lie within the sphere $r=a$, we see that $c$ is less than a, i.e., $c<a$, then we can rewrite the stream function (40) as

$$
\begin{equation*}
\psi_{0}(r, \theta)=\frac{m(r \cos \theta+c)}{\sqrt{r^{2}+c^{2}+2 r c \cos \theta}}-\frac{m(r \cos \theta-c)}{\sqrt{r^{2}+c^{2}-2 r c \cos \theta}}, \tag{41}
\end{equation*}
$$

which can be expended as

$$
\begin{gather*}
\psi_{0}(r, \theta)=m(r \cos \theta+c) \sum_{n=0}^{\infty} \frac{(-1)^{n} c^{n}}{r^{n+1}} P_{n}(\cos \theta) \\
-m(r \cos \theta-c) \sum_{n=0}^{\infty} \frac{c^{n}}{r^{n+1}} P_{n}(\cos \theta) \tag{42}
\end{gather*}
$$

After the reduction we see that $\psi_{0}(r, \theta) \approx O\left(\frac{1}{r}\right)$ for large $r$.
Therefore here the extension of Harper's theorem applies and yields the Stokes stream function for the flow within the sphere as

$$
\begin{equation*}
\psi(r, \theta)=m \cos \theta_{1}-m \cos \theta_{2}+\frac{m}{a c} \cos \theta_{3},-\frac{m}{a c} \cos \theta_{4}+\frac{m}{a c} r^{2}\left(R_{01}^{2}-R_{02}^{2}\right) \tag{43}
\end{equation*}
$$

where the last four terms constitute the image system outside the sphere $r=a$, and where

$$
\begin{gathered}
R_{01}^{2}=r^{2}+\frac{a^{4}}{c^{2}}+2\left(\frac{a^{2}}{c}\right) r \cos \theta, \text { and } R_{02}^{2}=r^{2}+\frac{a^{4}}{c^{2}}-2\left(\frac{a^{2}}{c}\right) r \cos \theta . \\
\text { V. CONCLUSION }
\end{gathered}
$$

We have shown an alternative way of the proof of Harper's theorem. In addition, we have extended the theorem for the flow interior to the same sphere and illustrate with an example. Numerical solutions of the problem can be useful for bubble rising research. This type of the problem has a great interest in geophysical applications.

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# A Unified Integral Associated with the Aleph Function 

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Abstract-In this note we obtain a unified new integral whose integrand contains product of Aleph function and generalized multivariable polynomials having general arguments. Several integrals containing many simpler functions follow as special cases of this integral.

Keywords: aleph function, generalized polynomials, hypergeometric function, fox's H- function. GJSFR-F Classification : FOR Code : MSC 2010: 31A10


Strictly as per the compliance and regulations of :


[^1]

# A Unified Integral Associated with the Aleph Function 

Harshita Garg ${ }^{\alpha}$ \& Ashok Singh Shekhawat ${ }^{\circ}$

Abstract- In this note we obtain a unified new integral whose integrand contains product of Aleph function and generalized multivariable polynomials having general arguments. Several integrals containing many simpler functions follow as special cases of this integral.
Keywords: aleph function, generalized polynomials, hypergeometric function, fox's H-function.

## I. Introduction

The Aleph function introduced by Südland et al [10] is defined as Mellin-Barnes type contour integrals as following:

$$
\begin{align*}
& =\frac{1}{2 \pi \mathrm{i}} \int_{\mathrm{L}} \Omega_{\mathrm{p}_{\mathrm{i}}, \mathrm{q}_{\mathrm{i}}, \mathrm{c}_{\mathrm{i}} ; r}^{\mathrm{e},}(\xi) \mathrm{x}^{-\xi} \mathrm{d} \xi \tag{1.1}
\end{align*}
$$

For all $\mathrm{x} \neq 0$, where $\mathrm{i}=\sqrt{-1}$ and

$$
\begin{equation*}
\Omega_{p_{i}, q_{i}, c_{i} ; \mathrm{r}}^{e, f}(\xi)=\frac{\prod_{\mathrm{j}=1}^{\mathrm{e}} \Gamma\left(\mathrm{~b}_{\mathrm{j}}+\mathrm{B}_{\mathrm{j}} \xi\right) \prod_{\mathrm{j}=1}^{\mathrm{f}} \Gamma\left(1-\mathrm{a}_{\mathrm{j}}-\mathrm{A}_{\mathrm{j}} \xi\right)}{\sum_{\mathrm{i}=1}^{\mathrm{r}} \mathrm{c}_{\mathrm{i}} \prod_{\mathrm{j}=\mathrm{f}+1}^{\mathrm{p}_{\mathrm{i}}} \Gamma\left(\mathrm{a}_{\mathrm{ji}}+\mathrm{A}_{\mathrm{ji}} \xi\right) \prod_{\mathrm{j}=\mathrm{e}+1}^{q_{i}} \Gamma\left(1-\mathrm{b}_{\mathrm{j}}-\mathrm{B}_{\mathrm{ji}} \xi\right)} \tag{1.2}
\end{equation*}
$$

The $\mathrm{L}=\mathrm{L}_{\mathrm{i} p o}$ is a suitable contour of the Mellin-Barnes type which runs from $\gamma-\mathrm{i} \infty$ to $\gamma+\mathrm{i} \infty$ with $\gamma \in \mathrm{R}$, the integers e , $\mathrm{f}, \mathrm{p}_{\mathrm{i}}$, $\mathrm{q}_{\mathrm{i}}$ satisfy the inequality $0 \leq \mathrm{f} \leq \mathrm{p}_{\mathrm{i}}, 1 \leq \mathrm{e} \leq$ $q_{i}, c_{i} \& 0 ; i=1, \ldots, r$. The parameters $A_{j}, B_{j}, A_{j i}, B_{j i}$ are positive real numbers and $a_{j}, b_{j}$, $\mathrm{a}_{\mathrm{ji}}, \mathrm{b}_{\mathrm{ji}}$ are complex numbers, such that the poles of $\Gamma\left(\mathrm{b}_{\mathrm{j}}+\mathrm{B}_{\mathrm{j}} \xi\right), \mathrm{j}=1,2, \ldots ., \mathrm{e}$ separating from those of $\Gamma\left(1-\mathrm{a}_{\mathrm{j}}-\mathrm{A}_{\mathrm{j}} \xi\right), \mathrm{j}=1, \ldots, \mathrm{f}$. All the poles of the integrand (1.2) are supposed to be easy and empty products are considered as unity. The existence conditions [4] for the Aleph function (1.2) are given below:

$$
\begin{gather*}
\psi_{\mathrm{k}}>0,|\arg (\mathrm{x})|<\frac{\pi}{2} \psi_{\mathrm{k}} ; \mathrm{k}=1, \ldots, \mathrm{r},  \tag{1.3}\\
\psi_{\mathrm{k}} \geq 0,|\arg (\mathrm{x})|<\frac{\pi}{2} \psi_{\mathrm{k}} \text { and } \mathrm{R}\left\{\Lambda_{\mathrm{k}}\right\}+1<0 \tag{1.4}
\end{gather*}
$$

[^2]Where

$$
\begin{align*}
& \psi_{k}=\sum_{j=1}^{f} A_{j}+\sum_{j=1}^{e} B_{j}-C_{k}\left(\sum_{j=f+1}^{p_{k}} A_{j k}+\sum_{j=e+1}^{q_{k}} B_{j k}\right)  \tag{1.5}\\
& \Lambda_{k}=\sum_{j=1}^{e} b_{j}-\sum_{j=1}^{f} a_{k}+C_{k}\left(\sum_{j=1}^{q_{k}} b_{j k}-\sum_{j=f+1}^{p_{k}} a_{j k}\right)+\frac{1}{2}\left(p_{k}-q_{k}\right) \tag{1.6}
\end{align*}
$$

The generalized polynomial defined by Srivastava [5] is as follows:

$$
\begin{equation*}
\mathrm{S}_{\mathrm{t}_{1}, \ldots, f_{\mathrm{s}}}^{\mathrm{e}_{1}, \ldots \mathrm{e}_{\mathrm{s}}}\left[z_{1}, \ldots, z_{\mathrm{s}}\right]=\sum_{\beta_{1}=0}^{\left[\mathrm{f}_{1} / \mathrm{e}_{1}\right]} \ldots \sum_{\beta_{\mathrm{s}}=0}^{\left[\mathrm{f}_{\mathrm{s}} / \mathrm{e}_{\mathrm{s}}\right]} \frac{\left(-\mathrm{f}_{1}\right)_{\mathrm{e}_{1} \beta_{1}}}{\beta_{1}!} \ldots \frac{\left(-\mathrm{f}_{\mathrm{s}}\right)_{\mathrm{e}_{\mathrm{s}} \beta_{\mathrm{s}}}}{\beta_{\mathrm{s}}!} . \mathrm{A}\left[\mathrm{f}_{1}, \beta_{1} ; \ldots ; \mathrm{f}_{\mathrm{s}}, \beta_{\mathrm{s}}\right] \mathrm{z}_{1}^{\beta_{1}} \ldots \mathrm{Z}_{\mathrm{s}}^{\beta_{\mathrm{s}}} \tag{1.7}
\end{equation*}
$$

Where $\mathrm{f}_{\mathrm{i}}=0,1,2 \ldots \quad \forall \mathrm{i}=(1, \ldots, \mathrm{~s}), \mathrm{e}_{1}, \ldots, \mathrm{e}_{\mathrm{s}}$ are arbitrary positive integers and the coefficients $\left[\mathrm{f}_{1}, \beta_{1} ; \ldots ; \mathrm{f}_{\mathrm{s}}, \beta_{\mathrm{s}}\right.$ ] are arbitrary constants, real or complex.

## iI. The Main Integral

We derive the following result:

$$
\begin{aligned}
& \int_{0}^{\infty} z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu} . \aleph_{\mathrm{p}_{\mathrm{i}}, \mathrm{q}_{\mathrm{i}}, \mathrm{c}_{\mathrm{i}} ; \mathrm{r}}^{\mathrm{e}, \mathrm{r}}\left[t\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] \\
& \mathrm{S}_{\mathrm{f}_{1}, \ldots, \ldots, \mathrm{f}_{\mathrm{s}}}^{\mathrm{e}_{1}, \ldots \mathrm{e}_{5}}\left[x_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}} \ldots . . x_{2}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{s}}\right] d z \\
& =2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\delta\left[f_{1} / e_{1}\right]} \sum_{\beta_{1}=0}^{\left[\mathrm{f}_{s} / e_{\mathrm{s}}\right]} \sum_{\beta_{s}=0} \frac{\left(-\mathrm{f}_{1}\right)_{\mathrm{e}_{1} \beta_{1}}}{\beta_{1}!} \ldots \frac{\left(-\mathrm{f}_{\mathrm{s}}\right)_{\mathrm{e}_{\mathrm{s}} \beta_{\mathrm{s}}}}{\beta_{\mathrm{s}}!} \mathrm{A}\left[\mathrm{f}_{1}, \beta_{1} ; \ldots ; \mathrm{f}_{\mathrm{s}}, \beta_{\mathrm{s}}\right] \mathrm{x}_{1}^{\beta_{1}} \ldots \mathrm{x}_{\mathrm{s}}^{\beta_{\mathrm{s}}} \alpha^{\left(-\mu-\sum_{\mathrm{i}=1}^{\mathrm{s}} \mathrm{a}_{\mathrm{i}} \beta_{\mathrm{i}}\right)}
\end{aligned}
$$

Where
(i) Where $\lambda>0, \operatorname{Re}(\delta, \mu, a)>0$
(ii) $\operatorname{Re}(\delta)-\operatorname{Re}(\mu)-\lambda \min _{1 \leq j \leq e} \operatorname{Re}\left(\frac{\mathrm{~b}_{\mathrm{j}}}{\beta_{\mathrm{j}}}\right)<0$ and
(iii) $\mathrm{e}_{1}, \ldots, \mathrm{e}_{\mathrm{s}}$ are arbitrary positive integers and the coefficients $\left[\mathrm{f}_{1}, \beta_{1} ; \ldots ; \mathrm{f}_{\mathrm{s}}, \beta_{\mathrm{s}}\right]$ are arbitrary constants, real or complex.
PROOF: The integral in (2.1) can be obtained by using the Aleph function in terms of Mellin-Barnes contour integral given by (1.1) and the definition of a generalized polynomials given by (1.7), then interchanging the order of summation and integration (which is permissible under the conditions stated with (2.1)) and evaluating the inner integral by using a result given by Oberthettinger F. [3] and we get the desired result.

## III. Special Cases

(1) Taking general class of polynomials in our main integral(2.1), we have

$$
\begin{aligned}
& \int_{0}^{\infty} Z z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu} \cdot \aleph_{\mathrm{p}_{\mathrm{i}}, \mathrm{q}_{\mathrm{i}}, \mathrm{c}_{\mathrm{i}}, \mathrm{r}}^{\mathrm{e}, \mathrm{r}}\left[t\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] \\
& \mathrm{S}_{f^{\prime}}^{\mathrm{e}_{1}, \ldots, e_{s}} \cdot\left[X_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}} \ldots . x_{2}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{s}}\right] d z
\end{aligned}
$$

$$
\begin{aligned}
& =2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\delta} \sum_{\beta_{1}, \ldots \ldots, \beta_{s}=0}^{\mathrm{e}_{1} \beta_{1}+\ldots+e_{s} \beta_{s} \leq f} \frac{\left(-\mathrm{f}^{\prime}\right)_{\mathrm{e}_{1} \beta_{1}+\ldots . .+e_{s} \beta_{\mathrm{s}}}}{\beta_{1}!\ldots \ldots . . \beta_{s}!} \mathrm{A}\left[\mathrm{f}^{\prime} ; \beta_{1} ; \ldots ; \beta_{\mathrm{s}}\right] \mathrm{x}_{1}^{\beta_{1}} \ldots \mathrm{x}_{\mathrm{s}}^{\beta_{\mathrm{s}}} \alpha^{\left(-\mu-\sum_{\mathrm{i}=1}^{s} \mathrm{a}_{\mathrm{i}} \beta_{\mathrm{i}}\right)}
\end{aligned}
$$

Where $\mathrm{e}_{1}, \ldots, \mathrm{e}_{\mathrm{s}}$ are arbitrary positive integers and the coefficients [ $\mathrm{f}^{\prime}, \beta_{1} ; \ldots ; \beta_{\mathrm{s}}$ ] are arbitrary constants, real or complex and valid sufficient conditions (i), (ii) and of (2.1). (2) If we take $s \rightarrow 1, e_{1}=2, A_{f_{1}, \beta_{1}}=(-1)^{\beta_{1}}$ then by applying our results given in (2.1) to the case of Hermite polynomial [7] and [12] and by taking

$$
\mathrm{S}_{\mathrm{f}_{1}}^{2}(\mathrm{x}) \rightarrow \mathrm{x}^{\mathrm{f}_{1} / 2} \mathrm{H}_{\mathrm{f}_{1}}\left[\frac{1}{2 \sqrt{\mathrm{x}}}\right]
$$

We have the following result

$$
\begin{aligned}
& \int_{0}^{\infty} Z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu}\left[X_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}}\right]^{\frac{\mathrm{f}_{1}}{2}} \\
& \left.\left.. \mathrm{H}_{\mathrm{f}_{1}}\left[\frac{1}{2 \sqrt{x_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}}}}\right]{\stackrel{N}{\mathrm{p}_{\mathrm{i}}, \mathrm{q}_{\mathrm{i}}, \mathrm{c}_{\mathrm{i}}, \mathrm{r}}}_{\mathrm{e,f}}\right] t\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] d z \\
& =2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\delta} \sum_{\beta_{1}=0}^{\left[\mathrm{f}_{1} / 2\right]} \frac{\left(-\mathrm{f}_{1}\right)_{2 \beta_{1}}}{\beta_{1}!}(-1)^{\beta_{1}} \mathrm{x}_{1}^{\beta_{1}} \alpha^{\left(-\mu-\mathrm{a}_{1} \beta_{1}\right)}
\end{aligned}
$$

Valid under the set of sufficient conditions (i) and (ii) of (2.1)
(3) For the Laguerre polynomials ([7] and [12]) setting $\mathrm{s} \rightarrow 1, \mathrm{~S}_{\mathrm{f}_{1}}^{\prime}(\mathrm{x}) \rightarrow \mathrm{L}_{\mathrm{f}_{1}}^{\left(\alpha^{\prime}\right)}(\mathrm{x})$ in which case $e_{1}=1, \quad A_{f_{1}, \beta_{1}}=\binom{f_{1}+\alpha^{\prime}}{f_{1}} \frac{1}{\left(\alpha^{\prime}+1\right)_{\beta_{1}^{\prime}}}$ the results (2.1) reduce to the following formulae:

$$
\begin{aligned}
& \int_{0}^{\infty} Z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu} L_{\mathrm{f}_{1}}^{(\alpha)}\left(x_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}}\right) \\
& . \aleph_{\mathrm{p}_{\mathrm{i}}, \mathrm{q}_{\mathrm{i}}, \mathrm{c}_{\mathrm{i}} ;}^{\mathrm{e}, \mathrm{r}}\left[t\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] d z \\
& =2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\delta} \sum_{\beta_{1}=0}^{\left[\mathrm{f}_{1} / 2\right]} \frac{\left(-\mathrm{f}_{1}\right)_{2 \beta_{1}}}{\beta_{1}!}\binom{\mathrm{f}_{1}+\alpha^{\prime}}{\mathrm{f}_{1}} \frac{1}{\left(\alpha^{\prime}+1\right)_{\beta_{1}^{\prime}}} \mathrm{x}_{1}^{\beta_{1}} \alpha^{\left(-\mu-\mathrm{a}_{1} \beta_{1}\right)}
\end{aligned}
$$

Valid under the set of sufficient conditions (i) and (ii) of (2.1)
(4) Taking $c_{i} \rightarrow 1$, Aleph function reduces to I-function given by Saxena [5], then our main integral (2.1) reduces to the following form:

$$
\begin{aligned}
& \int_{0}^{\infty} z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu} . I_{\mathrm{p}_{\mathrm{i}}, \mathrm{q}_{\mathrm{i}} \mathrm{r}}^{\mathrm{e}, \mathrm{r}}\left[t\left[z+\alpha+\left(\mathrm{z}^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] \\
& . S_{f_{1}, \ldots, \mathrm{f}_{\mathrm{s}}}^{\mathrm{e}_{1}, \ldots, \mathrm{e}_{\mathrm{s}}}\left[X_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}} \ldots . . x_{2}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{s}}\right] d z \\
& =2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\delta} \sum_{\beta_{1}=0}^{\left[\mathrm{f}_{1} / \mathrm{e}_{\mathrm{e}}\right]} \ldots \sum_{\beta_{\mathrm{s}}=0}^{\left[\mathrm{f}_{\mathrm{s}} / \mathrm{e}_{\mathrm{e}}\right]} \frac{\left(-\mathrm{f}_{1}\right)_{\mathrm{e}_{1} \beta_{1}}}{\beta_{1}!} \ldots \frac{\left(-\mathrm{f}_{\mathrm{s}}\right)_{\mathrm{e}_{\mathrm{s}} \beta_{\mathrm{s}}}}{\beta_{\mathrm{s}}!} \mathrm{A}\left[\mathrm{f}_{1}, \beta_{1} ; \ldots ; \mathrm{f}_{\mathrm{s}}, \beta_{\mathrm{s}}\right] \mathrm{x}_{1}^{\beta_{1}} \ldots \mathrm{x}_{\mathrm{s}}^{\beta_{\mathrm{s}}}
\end{aligned}
$$

Valid under the set of sufficient conditions (i), (ii) and (iii) of (2.1)
(5) Taking $\mathrm{c}_{\mathrm{i}} \rightarrow 1$ and $\mathrm{r}=1$ Aleph function reduces to Fox's H-function[1], then our main integral (2.1) reduces to the following form:

$$
\begin{aligned}
& \int_{0}^{\infty} z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu} \cdot H_{\mathrm{p}, \mathrm{q}}^{\mathrm{e}, \mathrm{f}}\left[t\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] \\
& . \mathrm{S}_{\mathrm{f}_{1}, \ldots, \mathrm{~s}_{\mathrm{s}}}^{\mathrm{e}_{1}, \ldots \mathrm{e}_{\mathrm{s}}}\left[X_{1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{1}} \ldots . X_{2}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a_{s}}\right] d z \\
& =2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\varepsilon^{[ }\left[\sum_{\beta_{1}=0}\left[\mathrm{f}_{\mathrm{e}}\right]\right.} \ldots \sum_{\beta_{\mathrm{s}}=0}^{\left[\mathrm{f}_{\mathrm{s}} / \mathrm{e}_{\mathrm{s}}\right]} \frac{\left(-\mathrm{f}_{1}\right)_{\mathrm{e}_{1} \beta_{1}}}{\beta_{1}!} \ldots \frac{\left(-\mathrm{f}_{\mathrm{s}}\right)_{\mathrm{e}_{\mathrm{s}} \beta_{s}}}{\beta_{\mathrm{s}}!} \mathrm{A}\left[\mathrm{f}_{1}, \beta_{1} ; \ldots ; \mathrm{f}_{\mathrm{s}}, \beta_{\mathrm{s}}\right] \mathrm{x}_{1}^{\beta_{1}} \ldots \mathrm{x}_{\mathrm{s}}^{\beta_{\mathrm{s}}}
\end{aligned}
$$

Valid under the set of sufficient conditions (i), (ii) and (iii) of (2.1)
(6) If we take $c_{i} \rightarrow 1, r=1$ and $e_{1}, \ldots, e_{s} \rightarrow e$ and $f_{1}, \ldots, f_{s} \rightarrow f$ i.s. $(1, . ., s \rightarrow 1)$ in the integral (2.1), we arrive at the following result which is obtained by Garg and Mittal [2].

$$
\begin{gathered}
\int_{0}^{\infty} z^{\delta-1}\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\mu} \cdot H_{\mathrm{p}, \mathrm{q}}^{\mathrm{e}, \mathrm{f}}\left[t\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-\lambda}\right] \\
\mathrm{S}_{\mathrm{f}}^{\mathrm{e}}\left[x\left[z+\alpha+\left(z^{2}+2 \alpha z\right)^{\frac{1}{2}}\right]^{-a}\right] d z \\
=2 \Gamma(2 \delta)\left(\frac{\alpha}{2}\right)^{\delta} \sum_{\beta=0}^{[\mathrm{f} / \mathrm{e}]} \frac{(-\mathrm{f})_{\mathrm{e} \beta}}{\beta!} \mathrm{A}[\mathrm{f}, \beta] \mathrm{x}^{\beta} \alpha^{(-\mu-\mathrm{a} \beta)} \\
\quad . H_{\mathrm{p}+2, \mathrm{q}+2}^{\mathrm{e}, \mathrm{f}+2}\left[\left.t \alpha^{-\lambda}\right|_{\left(\mathrm{b}_{1}, \mathrm{~B}_{\mathrm{p}}\right),\left(\mathrm{b}_{\mathrm{q}}, \mathrm{~B}_{\mathrm{q}}\right),(-\mu-\mathrm{a} \beta \beta-\delta ;),(1-\mu-\mathrm{a} \beta ; \lambda)} ^{(-\mu-\alpha ; \lambda),(1+\delta-\mu-\mathrm{a} \beta ; \lambda),\left(\mathrm{a}_{1}, \mathrm{~A}_{1}\right),\left(\mathrm{a}_{\mathrm{p}}, \mathrm{~A}_{\mathrm{p}}\right)}\right]
\end{gathered}
$$

## IV. Conclusion

The result so established may be found useful in several interesting situation appearing in the literature on mathematical analysis. The result (3.1) not only gives the value of the integral but also 'augments' the coefficients in the series in the integrand to give a ${ }_{4} F_{3}$ series as the integrated series.

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# Integrated Decision Making for Ground Handling Management 

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Summary- In this paper a hierarchical structure for the management of airport ground handling activities is proposed. The main decision making processes in charge of the managerial units composing a proposed ground handling management organization are considered. The global objective is to turn available the ground handling resources so that arriving and departing flight are serviced with as little delay as possible. Two operational situations are considered: a normal one where small delays are coped with when arriving and departing traffic is globally on schedule, and a disrupted situation where arriving or departing traffic suffer very large delays.

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# Integrated Decision Making for Ground Handling Management 

Salma Fitouri-Trabelsi ${ }^{\alpha}$, Felix Mora-Camino ${ }^{\circ}$, Carlos Alberto N. Cosenza ${ }^{\rho}$ \& Li Weigang ${ }^{\omega}$

Summary-In this paper a hierarchical structure for the management of airport ground handling activities is proposed. The main decision making processes in charge of the managerial units composing a proposed ground handling management organization are considered. The global objective is to turn available the ground handling resources so that arriving and departing flight are serviced with as little delay as possible. Two operational situations are considered: a normal one where small delays are coped with when arriving and departing traffic is globally on schedule, and a disrupted situation where arriving or departing traffic suffer very large delays.

## I. Introduction

The sustained global economic growth of the last decades has been possible with the development of improved means of communication and of transportation of people and goods. It has been particularly the case with air transportation where, during the last forty years, the number of passengers has been multiplied by seven. This increase of passenger volume has generated a permanent challenge for civil aviation authorities, airlines and airports to supply sufficient capacity to provide a safe transportation service with acceptable quality standards (Santos et al., 2010). In the last decade, new traffic management practices, such as Airport Collaborative Decision Making (A-CDM) (Eurocontrol, 2011), based on multi-agent and collaborative decision making concepts have been introduced at airports. Among the many activities which contribute to the safety and efficiency of air transportation, airport ground handling plays an important role even if it has remained in the shadow of other traffic activities in the Operations Research literature. While among the overall airport operations costs, ground handling costs represent a rather small portion, their dysfunction can generate huge extra costs for airlines and airports as well as high discomfort for passengers (Pestana, 2008).

In this study a hierarchical structure for the management of airport ground handling activities is considered. The global objective is to turn available the ground handling resources so that arriving and departing flight are serviced with as little delay as possible. Two operational situations are considered: a normal one where small delays are coped with when arriving and departing traffic is globally on schedule, and a disrupted situation where arriving or departing traffic suffer very large delays.

In the first situation a ground handling coordinator produces an estimate of the necessary resources from each ground handling service provider while these service providers assign the available resources to the scheduled ground handling activities. At both levels, the formulation of corresponding optimization problems leads to NPcomplete problems while a new solution should be at hand whenever new operations conditions appear. So, heuristic approaches have been developed to generate working solutions to this overall problem. While in the case of normal operations these heuristics consider the flights according to their nominal schedule, in the disrupted operations, flights are treated in accordance with an estimated degree of criticity computed by the

[^3]ground handling coordinator. The proposed approach is illustrated with traffic data from a large European airport.

## II. Hierarchical Structure for the Management of Ground Handilng at Airports

When considering ground handling organization in different airports, it appears that this organization depends strongly on the size and the physical organization of the airside as well as on the volume and composition of traffic. Then, a large diversity of actual ground handling organizations is found in major and medium size airports. Then it does not appear desirable to propose a general paradigm to organize airport ground handling since the resulting efficiency can be quite unequal from an airport to the next.
However, when some key characteristics are met, delimiting a specific class of ground handling situations, common organizing principles can be of interest.

Here some assumptions with respect to airport ground handling characteristics, which are frequently encountered in medium to large airports, are adopted. They are the following:

Here is considered the case of airports in which ground handling is performed by a set of specialized operators working in parallel under the management of the airport authorities.

The ground handling process is supposed to follow pre-established sequencings and to be performed at the parking stands. It is supposed that the parking stands are assigned to arriving flights by the airport and communicated through ATC, while the status of the parking stands is monitored by ATC which is in charge of driving the aircraft out of the parking position. It is also supposed that the arriving parking position is its departure parking position for the next flight. This last assumption introduces constraints on the ground handling activities.

From the considerations developed in the previous paragraph, it appears interesting to consider that the airport ground handling operators do not interact directly within the A-CDM framework (Eurocontrol,2011) , but through a ground handling coordinator.

The introduction of the GHC led to a hierarchical structure for the ground handling management as it showed in Fig.1.


## a) Ground Handling Coordinator

This coordinator will be a communication interface between the other A-CDM partners and the ground handling managers.
The principal functions of the GHC are:

- To provide to the other airport partners:
- predictions of ground handling delays
- Generation of milestones
- To provide to the ground handling managers:
- Predictions about activity levels
- required ground handling resources per period
i. Ground handling milestones monitoring

The ground handling activities around an aircraft can be divided in two set of operation:

- The set of arrival ground handling operations, $A_{i}^{g h}$, which includes all the ground handling activities which must be performed to conclude properly the current commercial flight. The main arrival ground handling activities are de-boarding passengers, unloading baggage, performing cleaning and sanitation.
- The set of departure ground handling operations, $D_{i}^{g h}$, which gathers the ground handling activities which must be performed to prepare the next commercial flight. The main departure activities are passengers boarding, baggage loading, fuelling, catering.
The possible milestones monitored by the ground handling coordinator are:
- time of start of arrival ground handling activities :

$$
\begin{equation*}
T_{i}^{a g h}=\min _{k \in A_{i}^{j k}}\left\{t_{i k}^{a g h}\right\} \tag{1}
\end{equation*}
$$

- time of completion of arrival ground handling activities :

$$
\begin{equation*}
\tau_{i}^{a g h}=\max _{k \in A_{i i}^{\prime!}}\left\{t_{i k}^{a g h}+d_{i k}^{a g h}\right\} \tag{2}
\end{equation*}
$$

- time of start of departure ground handling activities :

$$
\begin{equation*}
T_{i}^{d g h}=\min _{k \in D_{i j}^{s i j}}\left\{t_{i k}^{d g h}\right\} \tag{3}
\end{equation*}
$$

- time of completion of departure ground handling activities :

$$
\begin{equation*}
\tau_{i}^{d g h}=\max _{k \in D_{i i}^{g h}}\left\{t_{i k}^{d g h}+d_{i k}^{d g h}\right\} \tag{4}
\end{equation*}
$$ $d_{i k}^{d g h}$ is the duration of the ground handling activity k on aircraft i. All these time related variables and parameter adopt two values: their estimated value which can evolve and their effective value at completion.

ii. Global planning of ground handling resources

The planning of ground handling resources should be performed at start for a whole day of operation by considering as basic input information:

- the time schedule of arriving and departure flight,
- the operational characteristics of these flights.

The prediction of the necessary GH resources (vehicles and work force) over the operations period is performed in three steps:

- a global ground handling assignment (GGHA) problem is solved for a nominal schedule of flights. A fast heuristic solution is proposed ( greedy approach)
- totalization of necessary resources is performed for each time interval. Here a time interval within the operating period is chosen for the resources used by task $t$ :

$$
\begin{equation*}
u_{t}=\max \left\{\text { Timing, } \min _{j \in K} s_{j}^{t}\right\} \tag{5}
\end{equation*}
$$

- margins are added to the estimation of necessary resources:

For arrival ground handling activities:

$$
\begin{equation*}
r_{i}^{k}=n_{i}^{k}+p_{A}^{k} A_{i}^{k} \tag{6}
\end{equation*}
$$

For departure ground handling activities:

$$
\begin{equation*}
r_{i}^{k}=n_{i}^{k}+p_{D}^{k} D_{i}^{k} \tag{7}
\end{equation*}
$$

where: $n_{i}^{k}$ is the nominal number of teams (vehicle and staff) of type i necessary at period $k$ to process scheduled arrivals/departures, $r_{i}^{k}$ is the computed required number of teams of type $i$ necessary at period $k$, to process schedules arrivals/departures, included reserve, $A_{i}^{k}$ is the number of teams of type $i$ necessary to handle flight arrivals at parking stands during the previous half an hour which are supposed to be processed before period $k, D_{i}^{k}$ is the number of teams of type $i$ necessary to handle flight departures at parking stands during the previous half an hour which are supposed to be
processed before period k and, $p_{A}^{k}$ is the probability that an arrival scheduled within half an hour before period k is delayed and should be processed at period k and $p_{D}^{k}$ is the probability that a departure scheduled within half an hour before period $k$ is delayed and should be processed at period $k$.

## b) Ground Handling Manager

The ground handling manager has two principal functions:

- Planning operations
- Managing operations


## i. Planning operations

To achieve this function the ground handling manager has to:

- Solve its pairing problem to cover all planned demands for its services: during the current operations period. Result: list of duties which will be performed by its GHU's.
- Create the ground handling units by assigning its resources to its duties (a resource roastering problem).


## ii. Managing operations

Managing operations consists in the first time to update the assignment of his ground handling resources to aircraft considering the information received from the GHC has in case of:

- perturbation at the level the aircraft's arrival times
- perturbation at the level of the duration of performing of the tasks
- weather conditions (strong rain, snow, strong wind, etc.)

It consists also in monitoring the GHUs. A ground handling unit can be in the following states:

- deactivated: either the equipment is not ready (under repair or maintenance) or the operators are not available,
- waiting for assignment: the unit is enabled but has not been assigned to flights,
- assigned: the unit has been assigned to one or more flights, but the realization of the activity on the first of these flights is planned far in the time horizon,
- made ready to perform its next activity: this happens when the planned time to perform a ground handling activity is near. This corresponds either to the time necessary to adapt the resource to the flight to be served or to a minimum time delay to inform the operators of the next operation,
- operating: the unit is performing the activity (transfer operations and processing at aircraft or terminal).


## iil. Nominal Decision Making Processes with the Proposed Approach

a) The ground handling coordinator level

The decision making considered at this level is to solve the global ground handling assignment which is the first step of the global planning of ground handling resources.

A fast heuristic solution is proposed ( greedy approach) which consists in. this approach will ensure the feasibility of all ground handling operations. The idea of the porposed heuristic is to rank arriving and departing aircraft according to their planned start time of the corresponding ground operations (either arrival ground handling tasks or departure grand handling tasks). Then the GHC will process in this order each aircraft ground handling activity by linking each task to a route to build a ground handling duty:

- To cover task j at aircraft k it will search between the already created routes of type $j$, which one can cope with it, within the planned interval and at lower transportation cost.
- If none of the existing route provides a feasible solution
- and there are remaining capacity of type $j$ at the corresponding base, a new route of type $j$ starting at this base is created with first stop at aircraft k.
- and there are no remaining transport capacity at base of type $j$, add this task at the route of type j which minimizes the mix of resulting delay for aircraft k and of distance travelled to reach it with the weight $\lambda$.
Then repeat with all the expected ground handling tasks j at an arriving or departing aircraft.

This will produce feasible sets of duties (routes) to be performed by the different ground handling fleets and workforce. Then this data will be used by the ground handling coordinator to compute, according to the process proposed in the previous chapter, the level of resources that each ground handling manager must provide at each time period. These resources will be afterwards either effectively used to process aircraft and passengers or will remain as a warm reserve to face perturbations and incidents.

## b) The ground handling manager level

In a nominal situation, the ground handler fleet managers will assign a vehicle and a work team to each route. This vehicle may be changed by another to pursue the duty in accordance with operational considerations (refueling need, mechanical failure, etc) while work teams will be shifted according to labor and safety regulations.

Here it is supposed that there are enough spare vehicles and work teams to meet operational perturbations.
The proposed heuristic consists in:

- For each ground handling manager:
- Order the aircraft in accordance with their arrival/ departure time,
depending on the type of the ground handing fleet service.
depending on the type of the ground handing fleet service.
- Availability of all vehicles of the fleet.
- The distance from its current position to the considered aircraft

This is a rather simple greedy heuristic which provides for each fleet facing the current service demand a complete solution through a reduced computational effort. So there is no limitation in calling back this solution process any time a significant perturbation occurs.

In the case of ground handling fleets involved in unloading/loading activities at parked aircraft, aircraft will be duplicated considering their current scheduled arrival time at the parking position and their current scheduled departure time from the same parking position. Then each duplicate will be ordered according to increasing time.

From the solutions of the assignment problems solved by each ground handling manager, the ground handling coordinator forward the milestones corresponding to the completion of ground handling activities to the airlines and the ATC to produce if necessary new estimates for the departure schedule of the aircraft.

## c) Case of study

To validate the proposed ground handling organization and the associated decision making processes real traffic data from Palma de Mallorca Airport was considered. Palma de Mallorca Airport is, with respect to aircraft and passengers traffic, the third largest Spanish airport. During the summer period it is one of the busiest airports in Europe, and was used by 22.7 million passengers in 2011. The airport is the main base for the Spanish carrier Air Europa and also a focus airport for German carrier Air Berlin. It occupies an area of $6.3 \mathrm{~km} 2(2.4 \mathrm{sq} \mathrm{mi})$. Due to rapid growth of aircraft traffic and passenger numbers, additional infrastructure has been added to the

- Oor
two first terminals A (1965) and B (1972). It is composed now of two runways, four terminals and 180 parking stand ( 27 of them at aprons) (PDM, 2012). It can handle up to 25 million passengers per year, with a capacity to dispatch 12,000 passengers per hour.

To evaluate the proposed approach, we tested it using aircraft traffic for a 24 h period ( $01 / 08 / 2007$ ) with 690 arrivals and departures distributed between the four parking areas related with the four terminals of Palma de Mallorca Airport. Except for aircraft staying at night at the airport, all ground handling operations are done in the context of fast turnaround operations. Different sizes of ground handling fleets have been considered. The resulting earliest departure time for aircraft have been compared with the real time departure data, showing that with rather reduced ground handling fleets at each terminal, the proposed heuristic, coded in Java, does not generate additional delays. Fig. 2 displays the hourly traffic of arriving and departing aircraft on a typical summer day at this airport. It appears that aircraft traffic remains intense from early morning until the beginning of night hours.


Figure 2:01/08/2007 PDM Airport aircraft hourly traffic
The proposed heuristic approach has been tested for the aircraft traffic with the ground handling fleets of Fig.3.


Figure 3 : Nominal composition of ground handling fleets

## i. Implementing the global planning of ground handling resources

This approach is proposed to calculate the nominal number of resources required for each ground handling manager during a day of traffic. The solution of this approach is given in the Table 1. It represents the number of the aircraft which will be performed by each ground handling unit of each ground handling service provider.

| Ground <br> handling <br> activity | GHU1 | GHU2 | GHU3 | GHU4 | GHU5 | GHU6 | GHU7 | GHU8 | GHU9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| De-- <br> boarding/ <br> Boarding | 71 | 58 | 43 | 38 | 32 | 25 | 19 | 12 | 6 |


| passengers |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Unloading/ <br> Loading <br> baggage | 133 | 95 | 93 | 85 | 66 | 79 | 60 | 51 | 28 |
| Catering | 86 | 80 | 66 | 58 | 55 |  |  |  |  |
| Cleaning | 97 | 77 | 60 | 61 | 50 |  |  |  |  |
| Refuelling | 103 | 92 | 84 | 66 |  |  |  |  |  |
| Sanitation | 144 | 94 | 59 | 34 | 14 |  |  |  |  |
| Water <br> Supply | 103 | 82 | 66 | 53 | 41 |  |  |  |  |
| Push back | 118 | 112 | 84 | 37 | 31 |  |  |  |  |

Table 1 : Solution of hierarchical approach
Using this solution, only 14 aircraft will have a delay at the level of the departure times with a maximum delay of 14 minutes. The 14 aircraft that would leave their parking stand later that which it had been predicted their departure times match with busiest flight traffic period.

This global planning of ground handling resources as it has been described is composed of three steps:

For the first step, it has been supposed that the nominal number of each ground handling resources is presented in the figure.

In the second step, the unit time period which has been considered has been taken equal to the maximum between 5 minutes and the smallest duration of a ground handling operation, including transfer time according to the formula (5).

| Ground handling activity | Duration (min) |
| :---: | :---: |
| De-boarding passengers | 5 |
| Catering | 5 |
| Cleaning | 5 |
| Boarding passengers | 5 |
| Unloading baggage | 5 |
| Fuelling | 5 |
| Loading baggage | 5 |
| Sanitation | 5 |
| Potable water supply | 5 |
| Push-back | 5 |

Table 2: The unit time period of each ground handling operation results
The third step of the estimation of the necessary resources at a given time for all ground handling managers is performed by adding margins to the nominal level of demand of scheduled arrival and departure flights. This is done according to formula (6) and (7).

The figures presented below provide the size of the resources required for each ground handling manager to perform their corresponding ground handling tasks in case of perturbations that can occur during the day. As it can be seen, the number of reserved resources increases in the busiest flight traffic period (arrival/departure aircraft) according to the Fig-4.


Figure 4: Number of the resources required for each ground handling activities each of period of time
ii. Implementing the heuristics for on-line GHFA

To test the efficiency of this approach, the accurate arrival times of each considered flights are supposed to be communicated to the ground handling managers thirty minutes before the effective landing. Here, this allows the ground handling
managers to reassign the ground handling resources by considering the updated arrival times at the parking stands of the flights announced to land within the next half hour. Aircraft within five minutes to land have been supposed to maintain the previous assignment solution. No flight directed towards the considered airport has duration less than forty minutes. Then the real departure times where compared with the ones obtained through the proposed heuristic approach. The considered ground handling resources were the ones effectively existing at that airport.

The application of the proposed heuristic approach to the nominal schedule of arrivals during the considered reference day provided a feasible assignment for each ground handling manager in at most 0.3 seconds. These solutions led to delays with respect to scheduled departure schedule involving only 36 aircraft, with a maximum delay of 16 minutes. The average delay among delayed aircraft has been of 7 minutes. Fig. 5 displays the hourly distribution of delayed aircraft at departure resulting from the application of the proposed decentralized approach. Clearly, the occurrence of these delays corresponds to the busiest aircraft traffic periods at the airport where ground handling resources become short. The proposed heuristic could be restarted using higher ground handling resource levels provided by the ground handling coordinator to improve the expected delay performance of the system.


Figure 5: Hourly delays distribution resulting from the proposed heuristic
Historical data from 01/08/2007 at Palma de Mallorca Airport indicate that about 244 aircraft departures where delayed for multiple reasons, including one of the main reasons, ground handling delays. The maximum observed delay is about 520 minutes and the average delay among delayed aircraft has been of 30 minutes. There is information about the use of a particular system to manage ground handling at that airport.

It is clear, that in theory, the proposed heuristic approach provide significantly improved results with respect to departure delays. Then it can be expected for this particular airport that, even if the implementation of the proposed heuristic approach is not perfectly performed, some noticeable improvement with respect to the current practice will take effect. This is quite noteworthy since the proposed heuristic has not been particularly improved with respect to a basic greedy approach. IV. Ground Handling Management Under Disruption

To our knowledge there exists no specific definition for airport disruption while
recent works refer to this situation (Ploog, 2005) and (Tanger and al, 2013)
ut providing any definition. According to the British Standards Institute (Business
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tasks is sensitive to unexpected events such as additional travel time due to traffic congestion on airside service ways or machine breakdowns. Then it could be considered that ground handling management tackles in permanence disrupted situations.
a) fuzzy heuristic for on-line ground handling management problem.

The problem for each ground handling fleet is here to assign ground handling vehicles to arriving or departing aircraft so that each aircraft is serviced by a vehicle while, according to the current operational situation, no delay or a minimum delay is produced. For that, the airline ground station managers generate resources requests to the ground handling fleet managers. The produced schedules are based on the predicted arrival times as well as the scheduled departure times. These schedules take not only into consideration the possible variation of the ground handling tasks durations by using a fuzzy dual formalism (Cosenza, 2011; Cosenza, 2012), but consider also the criticality of the flight. This criticality depends on the current predicted delay as well as the operational consequences on other flights. Then more critical flights may get their ground handling solution treated before earlier less critical scheduled flights. The following notations are adopted: Each task of the turnaround process $t \in\{1, \ldots, T\}$ is carried out on an aircraft a(i) associated to a flight $\mathrm{i}, \mathrm{i} \in \mathrm{I}$, ( $\mathrm{I}=\mathrm{IA} \cup \mathrm{ID}$, IA is the set of arriving flights and ID is the set of departing flights) by a specific service provider $k \in\{1, \ldots, K\}$.

## b) Fuzzy-based ranking of flights

The first step of the proposed heuristic consists in performing an initial ordering of the flights in accordance with their current predicted arrival time $\hat{t}_{i}^{a}$ at their assigned parking amended by considering their criticality. To each arriving flight $\mathrm{i} \in \mathrm{I}$, can be assigned the difference $\Delta t_{i}^{a}=\hat{t}_{i}^{a}-\bar{t}_{i}^{a}$ between the predicted arrival time $\hat{t}_{i}^{a}$ and the scheduled arrival time $\bar{t}_{i}^{a}$. Here $\hat{t}_{i}^{a}$ and $\bar{t}_{i}^{a}$ can be either real numbers or fuzzy dual numbers, where $\hat{t}_{i}^{a}$ is provided by the ATC. Each arriving flight must cope with two types of operational constraints:

Connection constraints when arriving passengers must reach without delay another departing flight.

Departure schedule when the arriving aircraft must be ready to start a new flight with a tight schedule..

When considering connection constraints, let $C_{i}$ be the set of departing flights connected to arriving flight i . The time margin between fight i and each flight j in $C_{i}$ is given by:

$$
\begin{equation*}
\tilde{m}_{i j}^{a}=\bar{t}_{j}^{d}-\hat{t}_{i}^{a}-\max \left\{\tilde{d}_{d b}^{i}+\tilde{T}_{i j}, \tilde{d}_{u l}^{i}+\tilde{\theta}_{i j}\right\} \quad j \in C_{i} \tag{8}
\end{equation*}
$$

Here $\widetilde{T}_{i j}$ and $\tilde{\theta}_{i j}$ are respectively the connecting delay for passengers and luggage between flights i and j . The margin between arrival flight i and departure flight j serviced in immediate succession by the same aircraft is:

$$
\begin{equation*}
\tilde{m}_{i j}^{a}=\bar{t}_{j}^{d}-\hat{t}_{i}^{a}-\tilde{D}_{i j} \text { with } j=\sigma(i) \tag{9}
\end{equation*}
$$

where $\tilde{D}_{i j}$ is the minimum fuzzy dual duration of ground handling around arrival of flight i and departure of flight j . Here ${ }_{\sigma(i)}$ provides the number of the next flight serviced by the aircraft operating flight i. Then:

$$
\tilde{D}_{i j}=\max \left\{\begin{array}{c}
\tilde{d}_{u l}+\tilde{d}_{d u}+\tilde{d}_{u}  \tag{10}\\
\tilde{d}_{d b}+\tilde{d}_{c a}+\tilde{d}_{b d} \\
\tilde{d}_{d b}+\tilde{d}_{c l}+\tilde{d}_{b d} \\
\tilde{d}_{s a}+\tilde{d}_{w a}
\end{array}\right\}+\tilde{d}_{p b}
$$

Then, the fuzzy margin of arriving aircraft $i$ is given by:

$$
\begin{equation*}
\tilde{m}_{i}^{a}=\min _{j \in C_{i} \cup \sigma(i)} \tilde{m}_{i j}^{a} \tag{11}
\end{equation*}
$$

The amended arrival time for flight i is then given by:

$$
\begin{equation*}
\tilde{t}_{i}^{a}=\hat{t}_{i}^{a}+\tilde{m}_{i}^{a} \tag{12}
\end{equation*}
$$

To each departing flight $\mathrm{I} \in \mathrm{ID}$, can be assigned the difference $\Delta t_{i}^{d}=\hat{t}_{i}^{d}-\bar{t}_{i}^{d}$ between the predicted departure time $\hat{t}_{i}^{d}$ and the scheduled departure time $\bar{t}_{i}^{d}$. Here also, $\hat{t}_{i}^{d}$ and $\bar{t}_{i}^{d}$ can be either real numbers or fuzzy dual numbers. Symmetrically, each departing flight must cope with operational constraints related with successive flights by the same aircraft and flight connections for passengers and cargo.

In the case in which the ground handling tasks are relative to a departing flight j , the amended predicted time to start grand handling activities at the corresponding parking position is now given by:

$$
\begin{equation*}
\tilde{\tilde{t}}_{j}^{d}=\bar{t}_{j}^{d}-\min _{i \mid j \in C_{i} \text { and } i=\sigma^{-1}(j)} \tilde{m}_{i j}^{a} \tag{13}
\end{equation*}
$$

with

$$
\tilde{m}_{i \sigma(i)}^{a}=\max \left\{\begin{array}{c}
\tilde{d}_{f u}+\tilde{d}_{u}  \tag{14}\\
\tilde{d}_{c a}+\tilde{d}_{b d} \\
\tilde{d}_{w a}
\end{array}\right\}+\tilde{d}_{p b}
$$

Then, to each flight $i$, either arriving or departing, is assigned a time parameter $\tau_{i}$ such as:
For arriving flights:

$$
\begin{equation*}
\tau_{i}=\left\|\tilde{\tau}_{i}^{a}\right\| \tag{15}
\end{equation*}
$$

For departing flights:

$$
\begin{equation*}
\tau_{i}=\left\|\widetilde{\tau}_{i}^{d}\right\| \tag{16}
\end{equation*}
$$

where $\|\|$ is the fuzzy dual pseudo norm. Then the flights, either arriving or departing, present in the considered period of operation can be ranked according to an increasing $\tau_{i}$ index. Let the integer ra (i) be the amended rank of flight i.

## c) Ground Handling Fleets assignment to flights

Then flights are processed in the produced order ra(i) where ground handling vehicles are assigned to the corresponding aircraft. In the case of an arriving flight, ground handling arrival tasks (unloading luggage, de-boarding, cleaning and sanitation) are coped with by assigning the corresponding vehicles in accordance to their previous assigned tasks with other aircraft, their current availability, and their current distance to the considered aircraft. Here the common reference time schedule for the ground handling arrival tasks is $\hat{t}_{i}^{a}, i \in I_{A}$. In the case of a departing flight, ground handling departure tasks (fuelling, catering, luggage loading, boarding, water and push back) are also coped with by assigning the corresponding vehicles in accordance to their previous assigned tasks with other aircraft, their current availability, and their current distance to the considered aircraft. Here the common reference time schedule for the ground handling departure tasks is $B^{\text {low }}\left(\tilde{\tilde{t}}_{i}^{d}\right), i \in I_{D}$.

In both cases it is considered that the whole set of different ground handling vehicles necessary at arrival or departure is assigned by considering the common reference time schedule. This assignment of vehicles to flights either arriving or departing is performed on a greedy base by considering the closest vehicle available to
perform the required task. This will make that at the start of ground handling activities for an arrival or departure flight, all necessary resources will be nearby the parking place and that scheduling constraints between elementary ground handling tasks will be coped with locally without need of communication between the different ground handling fleet managers. This is a rather simple greedy heuristic which provides for each fleet facing the current service demand a complete solution through a reduced computational effort. So there is no limitation in calling back this solution process any time a significant perturbation occurs.

## d) Illustration of the proposed approach

To evaluate the proposed approach, the data used on the case of study of the previous part has been modified to create artificially a disruption situation. Here it has been considered that for any external reason, for exemple some severe weather conditions, a part of earlier scheduled arriving flights in the morning have been delayed and the airport operates under a concentrated arriving traffic at capacity between 11a.m. and 1 p.m.. Then, the effective arrivals and scheduled departures are those of Table.3.

It is considered that during and after this period the airside capacity of the airport is insufficient, including taxiing capacity with the appearence of queues of taxiing aircraft, parking positions with apron congestion and saturated ground handling capacity. In that conditions, transfer times for aircraft and ground handling units activities durations are subject to large uncertainties. Here it has been considered two scenarios for the uncertainty: in the first one additional delays are between $0 \%$ and $40 \%$ of the original duration between 11a.m. and 2 p.m. with return to nominal situation afterwards, in the second scenario additional delays are between $0 \%$ and $40 \%$ of the original duration between 11a.m. and noon, between $20 \%$ and $60 \%$ of the original duration between noon and 1:30 p.m., between $0 \%$ and $40 \%$ of the original duration between 1:30 p.m. and 2:30 p.m. with return to nominal situation afterwards.

$$
10 \mathrm{~h} \rightarrow 11 \mathrm{~h} \quad 11 \mathrm{~h} \rightarrow 12 \mathrm{~h} \quad 12 \mathrm{~h} \rightarrow 13 \mathrm{~h} \quad 13 \mathrm{~h} \rightarrow 14 \mathrm{~h} \quad 14 \mathrm{~h} \rightarrow 15 \mathrm{~h} \quad 15 \mathrm{~h} \rightarrow 16 \mathrm{~h}
$$

| Arrival <br> traffic | $20+30$ | $34+15$ | 25 | 7 | 15 | 15 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Scheduled <br> departures | 17 | 19 | $28+15$ | $17+20$ | $17+10$ | 17 |
|  |  |  |  |  |  |  |

## Table 3 : Effective arrivals and scheduled departures

In the case of this airport, there are no connections between the flights since in general this airport is a final destination for most of the passengers, so the arrival and the departure priority lists coincide. The priority list is calculated here by taking into account the predicted departure date of the flight $j$, which is the flight serviced by the same aircraft than for flight i. Here $\quad \tilde{D}_{i j}$ is the minimum fuzzy dual duration of ground handling around arrival of flight $i$ and departure of flight $j$ and the real arrival date of the flight i respecting the considering degree of uncertainty. This duration $\tilde{\Delta}_{i j}$, which is a fuzzy dual number, can be expressed by:

$$
\begin{equation*}
\tilde{\Delta}_{i j}=\left(\tilde{D}_{i j}+\hat{t}_{i}^{a}-\bar{t}_{j}^{d}\right) \tag{17}
\end{equation*}
$$

This application provided a feasible assignment for each ground handling manager in at most 0.4 seconds each updating of the priority lists.

The numerical results show that the delayed aircraft get in general the highest priority on the list. During the period of time between 11a.m and 2:30 p.m. ground handling achieves to serve 200 flights (arrival and departure of aircraft). The main numerical results are displayed in Table.4.

Scenario 1

| Mean delay for GH processing at arrival | 7.36 min | 8.86 min |
| :--- | :--- | :--- |
| Maximum delay for GH processing at arrival | 27 min | 30 min |
| Mean delay for GH processing at departure | 45.1 min | 59.4 min |
| Maximum delay for GH processing at departure | 195 min | 197 min |
|  |  |  |

## Table 4 : Statistical results for disruption scenarios

Fig. 6 displays the hourly distribution of delayed aircraft at departure resulting from the application of the proposed approach for the two scenarios. It appears that the impact of arriving traffic delays has resulted in an airport disruption situation which has extended in the afternoon. In the first scenario it can be considered that the disruption situation ends around 5 p.m. and in the other case it ends around 8 p.m. It appears then, that the more uncertainty about airside operations delays, the less the available ground handling capacity is able to cope with this disruption situation. Then insuring predictability of airside delays through fluidity of operations even in heavy activity levels situations emerge as an important objective.


Figure 6 : The hourly distribution of delayed aircraft at departure for the scenario 1 and the scenario 2

## V. Conclusion

In this paper, an organization for the ground handling management has been proposed. This proposed organization is based on the introduction of a ground handling coordination which has the role of a communication interface between the ground handling manager and the other airport partners. The solution of the different assignment problems solved by the ground handling coordinator and ground handling managers has been considered. A heuristic approaches has been developed in that case. In the case of the pairing problems faced by the ground handling managers, a heuristic approach has been developed. The whole process has been illustrated by considering a case study with real traffic where it has been assumed that flight arrival times are perfectly known half an hour in advance. Even if scheduled and effective arrival times are different, the adopted traffic situation can be considered as normal. Also the ground handling management
has been considered in the case of a huge traffic perturbation characterizing an airport traffic situation can be considered as normal. Also the ground handling management
has been considered in the case of a huge traffic perturbation characterizing an airport disruption. The operations planning procedures performed within the proposed management structure of ground handling have been revised by adopting temporary new objectives and taking into account the uncertainty with respect to activity delays in this situation. During the disruption period, the ground handling coordinator takes over the direction of the ground handling management by imposing to the ground over the direction of the ground handling management by imposing to the ground
handling managers, priority lists of flights to be processed. The computation of these priority lists makes use of fuzzy dual calculus to take into account delays uncertainty.
The feasibility of the proposed approach is displayed by considering the case of a priority lists makes use of fuzzy dual calculus to take into account delays uncertainty.
The feasibility of the proposed approach is displayed by considering the case of a disruption at Palma de Mallorca airport.

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# Special Pairs of Pythagorean Triangles and Dhuruva Number 

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Abstract- We present pairs of Pythagorean triangles, such that in each pair, the difference between their perimeters is two times the Dhuruva number. Also we present the number of pairs of primitive and non-primitive Pythagorean triangles.

Keywords: pairs of pythagorean triangles, dhuruva number, primitive and non-primitive pythagorean triangles.

GJSFR-F Classification : FOR Code : MSC 2010: 12D15

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# Special Pairs of Pythagorean Triangles and Dhuruva Number 

M. A. Gopalan ${ }^{\alpha}$, S. Vidhyalakshmi ${ }^{\circ}$, E. Premalatha ${ }^{\text { }}$ \& R. Presenna ${ }^{\omega}$

Abstract- We present pairs of Pythagorean triangles, such that in each pair, the difference between their perimeters is two times the Dhuruva number. Also we present the number of pairs of primitive and non-primitive Pythagorean triangles.
Keywords: pairs of pythagorean triangles, dhuruva number, primitive and non-primitive pythagorean triangles.

## I. Introduction

The fascinating branch of mathematics is the theory of numbers where in Pythagorean triangles have been a matter of interest to various mathematicians and to the lovers of mathematics, because it is a treasure house in which the search for many hidden connection is a treasure hunt. For a rich variety of fascinating problems one may refer [1-17].A careful observer of patterns may note that there is a one to one correspondence between the polygonal numbers and the number of sides of the polygon. Apart from the above patterns we have some more fascinating patterns of numbers namely Jarasandha numbers, Nasty numbers and Dhuruva numbers. These numbers have been presented in [18-21].

In [22-24], special Pythagorean triangles connected with polygonal numbers and Nasty numbers are obtained. Recently in [25], special Pythagorean triangles in connection with Hardy Ramanujan number 1729 are exhibited. In [26], Pythagorean triangles in connections with 5 -digit Dhuruva numbers are presented.

In this communication, we search for pairs of Pythagorean triangles, such that in each pair, the difference between their perimeters is two times the Dhuruva number.

## iI. Basic Definitons

## Definition 2.1

The ternary quadratic Diophantine equation given by $x^{2}+y^{2}=z^{2}$ is known as Pythagorean equation where $\mathrm{x}, \mathrm{y}, \mathrm{z}$ are natural numbers. The above equation is also referred to as Pythagorean triangle and denote it by $\mathrm{T}(\mathrm{x}, \mathrm{y}, \mathrm{z})$.

Also, in Pythagorean triangle $\mathrm{T}(\mathrm{x}, \mathrm{y}, \mathrm{z}): x^{2}+y^{2}=z^{2}, \mathrm{x}$ and y are called its legs and z its hypotenuse.

## Definition 2.2

Most cited solution of the Pythagorean equation is $x=m^{2}-n^{2}, y=2 m n, z=m^{2}+n^{2}$, where $m>n>0$. This solution is called primitive, if $m, n$ are of opposite parity and $\operatorname{gcd}(\mathrm{m}, \mathrm{n})=1$.

[^4]
## Definition 2.3: Dhuruva numbers

The numbers which do not change when we perform a single operation or a sequence of operations are known as Dhuruva numbers.

## iII. Method of Analysis

Let $\mathrm{PT}_{1}, \mathrm{PT}_{2}$ be two distinct Pythagorean triangles with generators $\mathrm{m}, \mathrm{q}(\mathrm{m}>\mathrm{q}$ $>0)$, and $\mathrm{p}, \mathrm{q}(\mathrm{p}>\mathrm{q}>0)$ respectively.Let $\mathrm{P}_{1}, \mathrm{P}_{2}$ be the perimeters of $\mathrm{PT}_{1}, \mathrm{PT}_{2}$ such that $\mathrm{P}_{1}-\mathrm{P}_{2}=2$ times the 3-digit Dhuruva numer 495.
The above relation leads to the equation

After performing numerical computations, it is noted that there are 82 distinct values for $m, p$ and $q$ satisfying (1). For simplicity and clear understanding, we have presented below in table1 the values of $m, p, q, P_{1}$ and $P_{2}$.

| S.no | m | q | p | $\mathrm{P}_{1}$ | $\mathrm{P}_{2}$ | $\left(\mathrm{P}_{1}-\mathrm{P}_{2}\right) / 2$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 166 | 164 | 165 | 109560 | 108570 | 495 |
| 2 | 167 | 162 | 166 | 109886 | 108896 | 495 |
| 3 | 168 | 160 | 167 | 110208 | 109218 | 495 |
| 4 | 169 | 158 | 168 | 110526 | 109536 | 495 |
| 5 | 170 | 156 | 169 | 110840 | 109850 | 495 |
| 6 | 171 | 154 | 170 | 111150 | 110160 | 495 |
| 7 | 172 | 152 | 171 | 111456 | 110466 | 495 |
| 8 | 173 | 150 | 172 | 111758 | 110768 | 495 |
| 9 | 174 | 148 | 173 | 112056 | 111066 | 495 |
| 10 | 175 | 146 | 174 | 112350 | 111360 | 495 |
| 11 | 176 | 144 | 175 | 112640 | 111650 | 495 |
| 12 | 177 | 142 | 176 | 112926 | 111936 | 495 |
| 13 | 178 | 140 | 177 | 113208 | 112218 | 495 |
| 14 | 179 | 138 | 178 | 113486 | 112496 | 495 |
| 15 | 180 | 136 | 179 | 113760 | 112770 | 495 |
| 16 | 181 | 134 | 180 | 114030 | 113040 | 495 |
| 17 | 182 | 132 | 181 | 114296 | 113306 | 495 |
| 18 | 183 | 130 | 182 | 114558 | 113568 | 495 |
| 19 | 184 | 128 | 183 | 114816 | 113826 | 495 |
| 20 | 185 | 126 | 184 | 115070 | 114080 | 495 |
| 21 | 186 | 124 | 185 | 115320 | 114330 | 495 |
| 22 | 187 | 122 | 186 | 115566 | 114576 | 495 |
| 23 | 188 | 120 | 187 | 115808 | 114818 | 495 |
| 24 | 189 | 118 | 188 | 116046 | 115056 | 495 |
| 25 | 190 | 116 | 189 | 116280 | 115290 | 495 |
| 26 | 191 | 114 | 190 | 116510 | 115520 | 495 |
| 27 | 192 | 112 | 191 | 116736 | 115746 | 495 |
| 28 | 193 | 110 | 192 | 116958 | 115968 | 495 |
| 29 | 194 | 108 | 193 | 117176 | 116186 | 495 |
| 30 | 195 | 106 | 194 | 117390 | 116400 | 495 |


| 31 | 196 | 104 | 195 | 117600 | 116610 | 495 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 32 | 197 | 102 | 196 | 117806 | 116816 | 495 |
| 33 | 198 | 100 | 197 | 118008 | 117018 | 495 |
| 34 | 199 | 98 | 198 | 118206 | 117216 | 495 |
| 35 | 200 | 96 | 199 | 118400 | 117410 | 495 |
| 36 | 201 | 94 | 200 | 118590 | 117600 | 495 |
| 37 | 202 | 92 | 201 | 118776 | 117786 | 495 |
| 38 | 203 | 90 | 202 | 118958 | 117968 | 495 |
| 39 | 204 | 88 | 203 | 119136 | 118146 | 495 |
| 40 | 205 | 86 | 204 | 119310 | 118320 | 495 |
| 41 | 206 | 84 | 205 | 119480 | 118490 | 495 |
| 42 | 207 | 82 | 206 | 119646 | 118656 | 495 |
| 43 | 208 | 80 | 207 | 119808 | 118818 | 495 |
| 44 | 209 | 78 | 208 | 119966 | 118976 | 495 |
| 45 | 210 | 76 | 209 | 120120 | 119130 | 495 |
| 46 | 211 | 74 | 210 | 120270 | 119280 | 495 |
| 47 | 212 | 72 | 211 | 120416 | 119426 | 495 |
| 48 | 213 | 70 | 212 | 120558 | 119568 | 495 |
| 49 | 214 | 68 | 213 | 120696 | 119706 | 495 |
| 50 | 215 | 66 | 214 | 120830 | 119840 | 495 |
| 51 | 216 | 64 | 215 | 120960 | 119970 | 495 |
| 52 | 217 | 62 | 216 | 121086 | 120096 | 495 |
| 53 | 218 | 60 | 217 | 121208 | 120218 | 495 |
| 54 | 219 | 58 | 218 | 121326 | 120336 | 495 |
| 55 | 220 | 56 | 219 | 121440 | 120450 | 495 |
| 56 | 221 | 54 | 220 | 121550 | 120560 | 495 |
| 57 | 222 | 52 | 221 | 121656 | 120666 | 495 |
| 58 | 223 | 50 | 222 | 121758 | 120768 | 495 |
| 59 | 224 | 48 | 223 | 121856 | 120866 | 495 |
| 60 | 225 | 46 | 224 | 121950 | 120960 | 495 |
| 61 | 226 | 44 | 225 | 122040 | 121050 | 495 |
| 62 | 227 | 42 | 226 | 122126 | 121136 | 495 |
| 63 | 228 | 40 | 227 | 122208 | 121218 | 495 |
| 64 | 229 | 38 | 228 | 122286 | 121296 | 495 |
| 65 | 230 | 36 | 229 | 122360 | 121370 | 495 |
| 66 | 231 | 34 | 230 | 122430 | 121440 | 495 |
| 67 | 232 | 32 | 231 | 122496 | 121506 | 495 |
| 68 | 233 | 30 | 232 | 122558 | 121568 | 495 |
| 69 | 234 | 28 | 233 | 122616 | 121626 | 495 |
| 70 | 235 | 26 | 234 | 122670 | 121680 | 495 |
| 71 | 236 | 24 | 235 | 122720 | 121730 | 495 |
| 72 | 237 | 22 | 236 | 122766 | 121776 | 495 |
| 73 | 238 | 20 | 237 | 122808 | 121818 | 495 |
| 74 | 239 | 18 | 238 | 122846 | 121856 | 495 |


| 75 | 240 | 16 | 239 | 122880 | 121890 | 495 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| 76 | 241 | 14 | 240 | 122910 | 121920 | 495 |
| 77 | 242 | 12 | 241 | 122936 | 121946 | 495 |
| 78 | 243 | 10 | 242 | 122958 | 121968 | 495 |
| 79 | 244 | 8 | 243 | 122976 | 121986 | 495 |
| 80 | 245 | 6 | 244 | 122990 | 122000 | 495 |
| 81 | 246 | 4 | 245 | 123000 | 122010 | 495 |
| 82 | 247 | 2 | 246 | 123006 | 122016 | 495 |

Thus it is seen that there are 82 pairs of Pythagorean triangles such that for each pair the difference in the perimeters is twice the 3 - digit Dhuruva number 495 .

Out of these 82 pairs of Pythagorean triangles 6 -pairs are non-primitive and in each of the remaining pairs, one of the triangles is primitive and the other is nonprimitive triangle.

A similar observation, regarding 5- digit and 6- digit dhuruva numbers are exhibited in the table 2 below.

| Dhuruva number | pairs of <br> Pythagorean <br> triangles | pairs of non- <br> primitve <br> Pythagorean <br> triangles | pairs of primitve <br> and non-primitve <br> Pythagorean <br> triangles |
| :---: | :---: | :---: | :---: |
| 53955 | 8992 | 908 | 8084 |
| 59995 | 9998 | 2111 | 7887 |
| 549945 | 91657 | 1 | 91656 |

IV. Conclusion

One may search for the connections between the pairs of Pythagorean triangles and other Dhuruva numbers of higher order.

## V. Acknowledgement

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# Dissemination Sinusoidal Waves in of A Viscoelastic Strip 

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Abstract-In this paper we consider the spectral problem for the wave propagation in extended plates of variable thickness. Describes how to solve problems and numerical results of wave propagation in infinitely large plates of variable thickness. Viscous properties of the material are taken into account by means of an integral operator Voltaire. The study is part of the spatial theory of visco elastic. The technique is based on the separation of spatial variables and formulating boundary eigenvalues problem to be solved by the method of orthogonal sweep Godunov. Numerical values obtained for the real and imaginary parts of phase velocity as a function of wave number. When this coincidence numerical results obtained with the known data.

Keywords: plate, spectral problem, frequency, variable thickness, orthogonal sweep.
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# Dissemination Sinusoidal Waves in of A Viscoelastic Strip 

Safarov Ismail Ibrahimovich ${ }^{\alpha}$, Akhmedov Maqsud Sharipovich ${ }^{\circ}$ \& Boltayev Zafar Ihterovich ${ }^{\circ}$


#### Abstract

Annotation- In this paper we consider the spectral problem for the wave propagation in extended plates of variable thickness. Describes how to solve problems and numerical results of wave propagation in infinitely large plates of variable thickness. Viscous properties of the material are taken into account by means of an integral operator Voltaire. The study is part of the spatial theory of visco elastic. The technique is based on the separation of spatial variables and formulating boundary eigenvalues problem to be solved by the method of orthogonal sweep Godunov. Numerical values obtained for the real and imaginary parts of phase velocity as a function of wave number. When this coincidence numerical results obtained with the known data.


Keywords: plate, spectral problem, frequency, variable thickness, orthogonal sweep.

## I. Introduction

Known [1,2] that in normal wave deformable layer (Lamb wave) is not orthogonal thickness, i.e. the integral of the scalar product of vectors of displacements of two different waves, considered as functions of the coordinate perpendicular to the surface layer is not zero. They also are not orthogonal conjugate waves is obtained by considering the dual problem. This introduces additional difficulties in solving practical problems $[3,4,8]$. In this paper, we present spectral problem formulation and methods of its tasks.

## II. Statement of the Wave Problem and the Basic Relations for the Plate Kirchhoff - Love Variable Thickness

Derive the fundamental relationships of the classical theory of plates with variable thickness on the basis of the principle of virtual displacements. In the threedimensional formulation of the elasticity problem reduces to the solution of the variation equation, which has the form:

$$
\begin{equation*}
\delta A_{F}+\delta A_{I}=0 \tag{1}
\end{equation*}
$$

For virtual work $\left(\delta A_{F}\right)$ internal forces, we have:

$$
\begin{equation*}
\delta A_{F}=-\delta \Pi=-\int_{V} \sigma_{i j} \delta \varepsilon_{i j} d V \tag{2}
\end{equation*}
$$

where $\Pi$ - potential energy; $\sigma_{i j}-$ components of the stress tensor; $\varepsilon_{i j}$ - components the deformation tensor; $V$ - the volume occupied by the body.

[^5]The physical properties of the plastic material describes the relationship

$$
\begin{equation*}
\sigma_{i j}=\bar{\lambda} \varepsilon_{k k} \delta_{i j}+2 \bar{\mu} \varepsilon_{i j} \quad(i, j, k=1,2,3) \tag{3}
\end{equation*}
$$

Where $\sigma_{i j}, \varepsilon_{i j}$ - components of the stress and strain tensors.

$$
\begin{equation*}
\lambda=\frac{E v(1+i \eta)}{(1+v)(1-2 v)} ; \quad \mu=\frac{E(1+i \eta)}{2(1+v)} \tag{4}
\end{equation*}
$$

Integrated in the case the Young's modulus of the viscoelastic material $E^{*}=E^{\prime}+i E^{\prime \prime}=E^{\prime}\left(1+i \eta_{e}\right)$ an analogue of the classical Young modulus [85]. Using a complex representation for the elastic modulus (Young's modulus) for the polymeric material can be written as

$$
\begin{equation*}
E^{*}(\omega)=E(\omega)[1+i \eta \omega] \tag{5}
\end{equation*}
$$

Where two functions of vibration frequency $E(\omega)$ and $\eta(\omega)$ may be represented by analytical variety of ways $[1,2]$.
For virtual work of inertial forces $\left(\delta A_{I}\right)$ we can write the following relation:

$$
\begin{equation*}
\delta A_{I}=-\int_{V} \rho \ddot{u}_{i} \delta u_{i} d V, \tag{3}
\end{equation*}
$$

where $\rho$ - body density; $u_{i}$ - displacement components; $\ddot{u}_{i}=\partial^{2} u i / \partial t^{2} ; t$ - time. Here and below, summation over repeated indices. Consider the wedge plate shown in Fig. 1, along the axis of an infinite $\boldsymbol{x}_{2}$. In accordance with the hypotheses of Kirchhoff - Love have:

$$
\begin{align*}
& \sigma_{13}=\sigma_{23}=\sigma_{33}=0 ; \\
& u_{i}=-x_{3} \frac{\ddot{a} W}{\ddot{a} x_{i}} ;  \tag{4}\\
& W\left(x_{3}\right) \equiv W,
\end{align*}
$$

where $W$ - deflection of the middle plane of the plate.
Neglecting in (3) members to take account of the inertia of rotation normal to the median plane we obtain:

$$
\begin{align*}
& -\int_{s} d s \int_{-h / 2}^{h / 2}\left(\sigma_{11} \delta \varepsilon_{11}+2 \sigma_{12} \delta \varepsilon_{12}+\sigma_{22} \delta \varepsilon_{22}\right) d x_{3}- \\
& -\int_{s} d s \int_{-h / 2}^{h / 2} \rho \frac{\partial^{2} W}{\partial t^{2}} \delta W d z=0 \tag{5}
\end{align*}
$$

The expressions for the components of strain and stress tensors are determined from the geometric relationships and relations generalized Hooke's law, which, taking into account the kinematic hypotheses (4) takes the form:

$$
\left\{\begin{array}{l}
\quad \varepsilon_{i j}=\frac{1}{2}\left(\frac{\partial u_{i}}{\partial x_{j}}+\frac{\partial u_{j}}{\partial x_{i}}\right)-x_{3} \frac{\partial^{2} W}{\partial x_{i} \partial x_{j}} ; \quad i, j=1,2  \tag{6}\\
\sigma_{11}=\frac{E}{1-v}\left(\varepsilon_{11}+v \varepsilon_{22}\right) \Gamma_{\kappa} ; \\
\sigma_{22}=\frac{E}{1-v}\left(\varepsilon_{22}+v \varepsilon_{11}\right) \Gamma_{\kappa} \\
\sigma_{12}=\frac{E}{1-v} \varepsilon_{12} \Gamma_{\kappa}
\end{array}\right.
$$

Where $E$ - Young's modulus; $v$ - Poisson's ratio of the plate material. Introducing the following notation:

$$
\begin{align*}
& M_{11}=-D\left(\frac{\partial^{2} W}{\partial x_{1}^{2}}+v \frac{\partial^{2} W}{\partial x_{2}^{2}}\right) \\
& M_{22}=-D\left(\frac{\partial^{2} W}{\partial x_{2}^{2}}+v \frac{\partial^{2} W}{\partial x_{1}^{2}}\right)  \tag{7}\\
& M_{12}=-D(1-v) \frac{\partial^{2} W}{\partial x_{1} \partial x_{2}} ; \\
& \bar{D}=\frac{\bar{E} h^{3}}{12\left(1-v^{2}\right)}=D_{1} \Gamma_{k} ; \quad D_{1}=\frac{E h^{3}}{12\left(1-v^{2}\right)} ; \quad \Gamma_{\kappa}=1+i \eta(\omega)
\end{align*}
$$

and integrating over the thickness of the plate, let (2.5) to the following form

$$
\begin{align*}
\int_{s}\left(M_{11} \frac{\partial^{2} \delta W}{\partial x_{1}^{2}}\right. & \left.+2 M_{12} \frac{\partial^{2} \delta W}{\partial x_{1} \partial x_{2}}+M_{22} \frac{\partial^{2} \delta W}{\partial x_{2}^{2}}\right) d S-  \tag{8}\\
& -\int_{s} \rho h \frac{\partial^{2} W}{\partial t^{2}} \delta W d s=0
\end{align*}
$$

Converting the first integral (8) twice by parts and equating to zero the coefficients of variation $\delta W$ inside the body and on its borders obtain the following differential equation:

$$
\begin{equation*}
\frac{\partial^{2} M_{11}}{\partial x_{1}^{2}}+2 \frac{\partial^{2} M_{12}}{\partial x_{1} \partial x_{2}}+\frac{\partial^{2} M_{22}}{\partial x_{2}^{2}}-\rho h \frac{\partial^{2} W}{\partial t^{2}}=0 \tag{9}
\end{equation*}
$$

with natural boundary conditions

$$
\begin{aligned}
M_{11}\left(0, I_{1}\right) & =0 \\
\frac{\partial M_{11}}{\partial x_{1}}+2 \frac{\partial M_{12}}{\partial x_{2}} & =0, \quad x_{1}=0, I_{1}
\end{aligned}
$$

the main alternative, which will be the following:

$$
\left\{\begin{array}{l}
\frac{\partial W}{\partial x_{1}}=0 \\
W=0, \quad x_{1}=0, l_{1}
\end{array}\right.
$$

Introducing new variables

$$
W, \varphi_{1}=\frac{\partial W}{\partial x_{1}}, M_{11}, Q_{1}=\frac{\partial M_{11}}{\partial x_{1}}+2 \frac{\partial M_{12}}{\partial x_{2}}
$$

and express through them $M_{22}$ with the help of (2.7). Then

$$
M_{22}=-D \frac{\partial^{2} W}{\partial x_{2}^{2}}+v M_{11}+v^{2} D \frac{\partial^{2} W}{\partial x_{2}^{2}}
$$

or

$$
\begin{equation*}
M_{22}=-\frac{E h^{3}}{12} \frac{\partial^{2} W}{\partial x_{2}^{2}}+v M_{11} \tag{10}
\end{equation*}
$$

We note that $M_{11}$ and $M_{22}$ are bending moments, at M12 the torque.
Thus, we arrive at the following system of equations:

$$
\left\{\begin{array}{l}
\frac{\partial W}{\partial x_{1}}=\varphi_{1} ; \\
\quad \frac{\partial \varphi}{\partial x_{1}}=-\frac{M_{11}}{D}-v \frac{\partial^{2} W}{\partial x_{2}^{2}} ;  \tag{11}\\
\frac{\partial M_{11}}{\partial x_{1}}=Q_{1}+\frac{\bar{E} h^{3}}{6(1+v)} \frac{\partial^{2} \varphi}{\partial x_{2}^{2}} ; \\
\frac{\partial Q_{1}}{\partial x_{1}}=-v \frac{\partial^{2} M_{11}}{\partial x_{2}^{2}}+\frac{\bar{E} h^{3}}{12} \frac{\partial^{2} W}{\partial x_{2}^{2}}+\rho h \frac{\partial^{2} W}{\partial t^{2}},
\end{array}\right.
$$

Or

$$
\left\{\begin{array}{l}
\frac{\partial W}{\partial x_{1}}=\varphi_{1} ; \\
\quad \frac{\partial \varphi}{\partial x_{1}}=-\frac{6(1-v)}{h^{3}} \frac{M_{11} \cdot 2(1-v)}{\bar{E}}-v \frac{\partial^{2} W}{\partial x_{2}^{2}} ; \\
\quad \frac{2(1+v)}{\bar{E}} \frac{\partial M_{11}}{\partial x_{1}}=\frac{2(1+v)}{\bar{E}} Q_{1}+\frac{h^{3}}{3} \frac{\partial \varphi}{\partial x_{2}} ; \\
\frac{2(1+v)}{\bar{E}} \frac{\partial Q_{1}}{\partial x_{1}}=-v \frac{2(1+v)}{\bar{E}} \frac{\partial^{2} M_{11}}{\partial x_{2}^{2}}+\frac{(1+v) h^{3}}{6} \frac{\partial^{4} W}{\partial x_{2}^{4}}+\frac{2(1+v)}{\bar{E}} \rho h \frac{\partial^{2} W}{\partial t^{2}},
\end{array}\right.
$$

Or

$$
\left\{\begin{align*}
& \frac{\partial y_{1}}{\partial x_{1}}=y_{2} ;  \tag{12}\\
& \frac{\partial y_{2}}{\partial x_{1}}=-\frac{6(1-v)}{h^{3}} y_{3}-v \frac{\partial^{2} y_{1}}{\partial x_{2}^{2}} \\
& \frac{\partial y_{3}}{\partial x_{1}}=y_{4}+\frac{h^{3}}{3} \frac{\partial y_{2}}{\partial x_{2}} \\
& \frac{\partial y_{4}}{\partial x_{1}}=-v \frac{\partial^{2} y_{3}}{\partial x_{2}^{2}}+\frac{(1+v) h^{3}}{6} \frac{\partial^{4} y_{1}}{\partial x_{2}^{4}}+\frac{h}{C_{s}^{2}} \frac{\partial^{2} y_{1}}{\partial t^{2}}
\end{align*}\right.
$$

Where $y_{1}=W, y_{2}=\varphi_{1}, y_{3}=\frac{2(1+v)}{E} M_{11}, y_{4}=\frac{2(1+v)}{E} Q, \tilde{N}_{s}^{2}=\frac{E}{2(1+v) \rho}, C_{s}-$ shear wave Velocity

Among the many solutions of (12) we choose those that describe harmonic plane waves propagating along the axis $x_{2}$

$$
\begin{equation*}
y_{i}=z_{i}\left(x_{1}\right) e^{i\left(\hat{e} \tilde{o}_{2}-\omega t\right)} \tag{13}
\end{equation*}
$$

Substituting the solution (13) in the system of differential equations (12), we obtain a system of ordinary differential equations of the first order, solved for the derivative:

$$
\left\{\begin{array}{l}
z_{1}^{\prime}=z_{2}  \tag{14}\\
z_{2}^{\prime}=-\frac{6(1-v)}{h^{3}} z_{3}+v \kappa^{2} z_{1} \\
z_{3}^{\prime}=z_{4}-\frac{h^{3} \Gamma_{k}}{3} \kappa^{2} z_{2} \\
z_{4}^{\prime}=v \kappa^{2} z_{3}+\frac{(1+v) h}{6} \kappa^{4} z_{1}-h\left(\frac{\omega}{C_{s}}\right)^{2} \Gamma_{k} z_{1}
\end{array}\right.
$$

The boundary conditions for this system can be written as follows:
a) free left edge of the plate:

$$
\begin{equation*}
z_{3}(0)=z_{4}(0)=0 \tag{15}
\end{equation*}
$$

б) free right edge of the plate:

$$
\begin{equation*}
z_{3}\left(l_{1}\right)=z_{4}\left(l_{1}\right)=0 \tag{16,a}
\end{equation*}
$$

в) pinched right edge of the plate:

$$
\begin{equation*}
z_{1}\left(l_{1}\right)=z_{2}\left(l_{1}\right)=0 \tag{16,б}
\end{equation*}
$$

T Thus formed the spectral problem (14-16) in the parameter $\omega$, describing the propagation of flexural waves in a flat edge plate Kirchhoff-Love.

## iii. Basic Relations for Timoshenko Plates of Variable Thickness. Statement of the Wave Problem

Applying the principle of virtual displacements (1-3), replacing the KirchhoffLove hypotheses (2.4) on the hypothesis Timoshenko:

$$
\begin{align*}
& \sigma_{33}=0 ; \quad \sigma_{3 i}=\frac{\chi \bar{E}}{2(1+v)}\left(\frac{\partial W}{\partial x_{i}}-\theta_{i}\right) ;  \tag{17}\\
& u_{i}^{\left(x_{3}\right)}=x_{3} \theta_{i} ; W^{\left(x_{3}\right)}=W ; \quad i=1,2,
\end{align*}
$$

where $\theta_{i}$ - normal rotation angles (Fig. 2) $\chi$ - correction factor that takes into account the distribution of shear stresses across the thickness.


Figure 1 : Design scheme


Figure 2 : shows the angle of rotation of the normal In this case, the tensor components of strain and stress take the form:

$$
\varepsilon_{i j}=-\frac{1}{2} x_{3}\left(\frac{\partial \theta_{i}}{\partial x_{j}}+\frac{\partial \theta_{j}}{\partial x_{i}}\right) ;
$$

$$
\begin{align*}
& \varepsilon_{3 i}=\frac{1}{2}\left(\frac{\partial W}{\partial x_{i}}-\theta_{i}\right) \\
& \sigma_{11}=-\frac{E \Gamma_{k}}{1-v^{2}} x_{3}\left(\frac{\partial \theta_{1}}{\partial x_{1}}+v \frac{\partial \theta_{2}}{\partial x_{2}}\right)  \tag{18}\\
& \sigma_{22}=-\frac{E \Gamma_{k}}{1-v^{2}} x_{3}\left(\frac{\partial \theta_{2}}{\partial x_{2}}+v \frac{\partial \theta_{1}}{\partial x_{1}}\right) \\
& \sigma_{12}=-\frac{E \Gamma_{k}}{2(1+v)} x_{3}\left(\frac{\partial \theta_{1}}{\partial x_{2}}+v \frac{\partial \theta_{2}}{\partial x_{1}}\right) \\
& \sigma_{3 i}=\frac{\chi E \Gamma_{k}}{2(1+v)}\left(\frac{\partial W}{\partial x_{i}}-\theta_{i}\right), \quad i, j=1,2
\end{align*}
$$

$\mathrm{N}_{\text {otes }}$

Substitute the expression for the work on virtual displacements, we obtain:

$$
\begin{align*}
\delta A=\int_{-h / s}^{h / 2} \int_{s} & {\left[-\sigma_{i j} \frac{x^{3}}{2}\left(\frac{\partial \delta \theta_{i}}{\partial x_{j}}+\frac{\partial \delta \theta_{j}}{\partial x_{i}}\right)+\sigma_{3 i}\left(\frac{\partial \delta W}{\partial x_{i}}-\delta \theta_{i}\right)+\right.} \\
& \left.+\rho \ddot{W} \delta W+\rho x_{3}^{2} \ddot{\theta}_{i} \delta \theta_{i}\right] d S d x_{3}=0 \tag{19}
\end{align*}
$$

Or by introducing a notation for the corresponding moments:

$$
\begin{align*}
& \bar{M}_{11}=D_{1} \Gamma_{k}\left(\frac{\partial \theta_{1}}{\partial x_{1}}+v \frac{\partial \theta_{2}}{\partial x_{2}}\right)=\Gamma_{k} M_{11} ; \\
& \bar{M}_{22}=D_{1} \Gamma_{k}\left(\frac{\partial \theta_{2}}{\partial x_{2}}+v \frac{\partial \theta_{1}}{\partial x_{1}}\right)=\Gamma_{k} M_{22} ;  \tag{20}\\
& \bar{M}_{12}=D_{2} \Gamma_{k}\left(\frac{\partial \theta_{1}}{\partial x_{2}}+\frac{\partial \theta_{2}}{\partial x_{1}}\right)=\Gamma_{k} M_{12}
\end{align*}
$$

where $D_{2}=\frac{1}{2} D_{1}$

$$
\begin{aligned}
& M_{22}=-D_{1}\left(\frac{\partial \theta_{2}}{\partial x_{2}}+v \frac{\partial \theta_{1}}{\partial x_{1}}\right) \\
& M_{11}=-D_{1}\left(\frac{\partial \theta_{1}}{\partial x_{1}}+v \frac{\partial \theta_{2}}{\partial x_{2}}\right)
\end{aligned}
$$

$$
M_{12}=D_{2}\left(\frac{\partial \theta_{1}}{\partial x_{2}}+\frac{\partial \theta_{2}}{\partial x_{1}}\right)
$$

and integrating over $x_{3}$ we have

$$
\begin{align*}
& \delta A=-\int_{s}\left[-\frac{\partial}{\partial x_{j}}\left(\bar{M}_{i j} \delta \theta_{i}\right)+\frac{\partial}{\partial x_{j}}\left(h \delta_{3 j} \delta W\right)\right] d S+ \\
& +\int_{s}\left(-\frac{\partial \bar{M}_{i j}}{\partial x_{j}} \delta \theta_{i}+\frac{\partial\left(h \bar{\sigma}_{3 j}\right)}{\partial x_{j}} \delta W+h \bar{\sigma}_{3 i} \delta \theta_{i}-\right. \\
& \left.-\rho h \ddot{W} \delta W-\frac{\rho h^{3}}{12} \ddot{\theta}_{i} \delta \theta_{i}\right) d S=0 \tag{21}
\end{align*}
$$

Integrating (21) by parts and equating to zero the coefficients of variation $\delta W$ and $\delta \theta_{i}$ inside the body and on its borders obtain the following system of differential equations

$$
\left\{\begin{array}{l}
-\frac{\partial M_{12}}{\partial x_{2}}-\frac{\partial M_{11}}{\partial x_{1}}+h \sigma_{31}-\frac{\rho h^{3}}{12 \Gamma_{k}} \ddot{\theta}_{1}=0 ;  \tag{22}\\
-\frac{\partial M_{22}}{\partial x_{2}}-\frac{\partial M_{12}}{\partial x_{1}}+h \sigma_{32}-\frac{\rho h^{3}}{12 \Gamma_{k}} \ddot{\theta}_{2}=0 ; \\
\frac{\partial\left(h \sigma_{32}\right)}{\partial x_{2}}+\frac{\partial\left(h \sigma_{31}\right)}{\partial x_{1}}-\frac{\rho h \ddot{W}}{\Gamma_{k}}=0
\end{array}\right.
$$

With natural boundary conditions:

$$
\left\{\begin{array}{l}
M_{12}=0 \\
M_{11}=0 \\
h \sigma_{31}=0, x_{1}=0, l_{1}
\end{array}\right.
$$

The main alternative, which will be the following:

$$
\left\{\begin{array}{l}
\theta_{1}=0 \\
\theta_{2}=0 ; \\
W=0, x_{1}=0, l_{1}
\end{array}\right.
$$

Equation (22) is a differential complex coefficients, it is possible to write in the following form

$$
\left(\begin{array}{l}
-\frac{\partial M_{12}}{\partial x_{2}}-\frac{\partial M_{11}}{\partial x_{1}}+h \tau_{31}-\frac{s h^{3}}{12 \Gamma_{K R}} \theta_{1}^{\prime \prime} \\
\frac{-\partial M_{22}}{\partial x_{2}}-\frac{\partial M_{12}}{\partial x_{1}}+h \tau_{32}-\frac{s h^{3}}{12 \Gamma_{K R}} \theta^{\prime \prime} \\
\frac{\partial\left(h \tau_{32}\right)}{\partial x_{2}}+\frac{\partial\left(h \tau_{31}\right)}{\partial x_{1}}-\frac{s h^{3}}{\Gamma_{K 12}} \ddot{W}
\end{array}\right)+i \Gamma_{K I}\left(\begin{array}{l}
-\frac{\partial M_{12}}{\partial x_{2}}-\frac{\partial M_{11}}{\partial x_{1}}+h \tau_{31} \\
\frac{-\partial M_{22}}{\partial x_{2}}-\frac{\partial M_{12}}{\partial x_{1}}+h \tau_{32} \\
\frac{\partial\left(h \tau_{32}\right)}{\partial x_{2}}+\frac{\partial\left(h \tau_{31}\right)}{\partial x_{1}}
\end{array}\right)=0
$$

The main variables in this system, we assume: $W_{1}, \theta_{1}, \theta_{2}, M_{12}, M_{11}, Q_{1}=h \sigma_{31}$. Out of the equation variables $M_{22}$ and $Q_{2}$.

$$
M_{22}=-\frac{E h^{3}}{12} \frac{\partial \theta_{2}}{\partial x_{2}}+v M_{11} ; Q_{2}=h \sigma_{32}=\frac{\chi E h}{2(1+v)}\left(\frac{\partial W}{\partial x_{2}}-\theta_{2}\right) .
$$

Thus we arrive at the following system of equations:

$$
\left\{\begin{align*}
\frac{\partial W}{\partial x_{1}} & =\theta_{1}+\frac{2(1+v)}{\chi E h} Q_{1} ; \\
\frac{\partial \theta_{2}}{\partial x_{1}} & =-\frac{\partial \theta_{1}}{\partial x_{2}}-\frac{24(1+v)}{E h^{3}} M_{12} ; \\
\frac{\partial \theta_{1}}{\partial x_{1}} & =-v \frac{\partial \theta_{2}}{\partial x_{2}}-\frac{12\left(1-v^{2}\right)}{E h^{2}} M_{12} ;  \tag{23}\\
\frac{\partial M_{11}}{\partial x_{1}} & =-\frac{\partial M_{12}}{\partial x_{2}}+Q_{1}-\frac{p h^{3}}{12 \Gamma_{k}} \ddot{\theta}_{1} ; \\
\frac{\partial M_{22}}{\partial x_{1}} & =-\frac{E h^{3}}{12} \frac{\partial^{2} \theta_{2}}{\partial x_{2}^{2}}-v \frac{\partial M_{11}}{\partial x_{2}}+\frac{\chi E h}{2(1+v)}\left(\frac{\partial W}{\partial x_{2}}-\theta_{2}\right)-\frac{p h^{3}}{12 \Gamma_{k}} \ddot{\theta}_{2} ; \\
\frac{\partial Q_{1}}{\partial x_{1}} & =-\frac{\chi E h}{2(1+v)}\left(\frac{\partial^{2} W}{\partial x_{2}^{2}}-\frac{\partial \theta_{2}}{\partial x_{2}}\right)+\frac{\rho h \ddot{W}}{\Gamma_{k}} .
\end{align*}\right.
$$

or

$$
\begin{align*}
& \frac{\partial y_{1}}{\partial x_{1}}=y_{2}+\frac{y_{4}}{\chi h} ; \quad \frac{\partial y_{2}}{\partial x_{1}}=-v \frac{\partial y_{3}}{\partial x_{2}}-\frac{6(1-v)}{h^{3}} y_{5} ; \\
& \frac{\partial y_{3}}{\partial x_{1}}=-\frac{\partial y_{2}}{\partial x_{2}}-\frac{12}{h^{3}} y_{6} ; \\
& \frac{\partial y_{4}}{\partial x_{1}}=\chi h \frac{\partial}{\partial x_{2}}\left(y_{3}-\frac{\partial y_{1}}{\partial x_{2}}\right)+\frac{h}{\Gamma_{k}} \frac{\partial^{2} y_{2}}{\partial \tilde{t}^{2}} ;  \tag{24}\\
& \frac{\partial y_{5}}{\partial x_{1}}=-\frac{\partial y_{6}}{\partial x_{2}}+y_{4}-\frac{h^{3}}{12 \Gamma_{k}} \frac{\partial^{2} y_{2}}{\partial \tilde{t}^{2}} ; \\
& \frac{\partial y_{6}}{\partial x_{1}}=\frac{\partial}{\partial x_{2}}\left(\frac{(1+v) h^{3}}{6} \cdot \frac{\partial y_{3}}{\partial x_{2}}-v y_{5}\right)+\chi h\left(\frac{\partial y_{1}}{\partial x_{2}}-y_{3}\right)-\frac{h^{3}}{12 \Gamma_{k}} \frac{\partial^{2} y_{3}}{\partial \tilde{t}^{2}} .
\end{align*}
$$

Where

$$
\begin{aligned}
& y_{1}=W ; \quad y_{2}=\theta_{2} ; \quad y_{3}=\theta / v ; \quad y_{4}=\frac{2(1+v)}{E} Q_{1} ; \\
& y_{5}=\frac{4(1+v)}{1-v} M_{12} ; \quad y_{6}=\frac{h\left(1-v^{2}\right)}{E v} M_{12} \\
& M_{22}=-D\left(\frac{\partial \theta_{2}}{\partial x_{2}}+v \frac{\partial \theta_{1}}{\partial x_{1}}\right)+v M_{11}-v M_{11}=
\end{aligned}
$$

$$
\begin{aligned}
& =-D\left(1-v^{2}\right) \frac{\partial \theta_{2}}{\partial x_{2}}+v M_{11}=-\frac{E h^{3}}{12\left(1-v^{2}\right)}\left(1-v^{2}\right) \frac{\partial \theta_{2}}{\partial x_{2}}+v M_{11}= \\
& =-\frac{E h^{3}}{12} \frac{\partial \theta_{2}}{\partial x_{2}}+v M_{11}
\end{aligned}
$$

Finding, as before, the solutions described by a plane harmonic waves propagating along the axis $x_{1}$, we seek a solution of (24) in the form

$$
\left\{\begin{array}{l}
y_{1}=z_{1}\left(x_{1}\right) \cos \left(\kappa x_{2}-\omega t\right) ;  \tag{25}\\
y_{2}=z_{2}\left(x_{1}\right) \cos \left(\kappa x_{2}-\omega t\right) ; \\
y_{3}=z_{3}\left(x_{1}\right) \sin \left(\kappa x_{2}-\omega t\right) ; \\
y_{4}=z_{4}\left(x_{1}\right) \cos \left(\kappa x_{2}-\omega t\right) ; \\
y_{5}=z_{5}\left(x_{1}\right) \cos \left(\kappa x_{2}-\omega t\right) ; \\
y_{6}=z_{6}\left(x_{1}\right) \sin \left(\kappa x_{2}-\omega t\right) .
\end{array}\right.
$$

Substituting relation (25) in the system of differential equations (24) we obtain a system of ordinary differential equations of the first order, solved for the derivative:

$$
\left\{\begin{array}{l}
z_{1}^{\prime}=z_{2}+\frac{z_{n}}{\chi h} ; \\
z_{2}^{\prime}=-v \kappa z_{3}-\frac{6(1-v)}{3} z_{5} ; \\
z_{3}^{\prime}=\kappa z_{2}-\frac{12}{h^{3}} z_{6} ; \\
z_{4}^{\prime}=\chi h \kappa z_{3}+\kappa^{2}\left(\chi h-\frac{h c^{2}}{\Gamma_{n}}\right) z_{1} ;  \tag{26}\\
z_{5}=-\kappa z_{6}+z_{4}+\frac{h^{3}}{12 \Gamma_{\pi}} \omega^{2} z_{2} ; \\
z_{6}^{\prime}=-\chi h \kappa z_{1}-\left[\chi h+\frac{\kappa^{2} h^{3}}{12 \Gamma_{n}}\left(2(1+v)-\frac{c^{2}}{\Gamma_{n}}\right)\right] z_{3}+v \kappa z_{5}
\end{array}\right.
$$

The boundary conditions for this system can be written as follows:
a) free left edge of the plate:

$$
\begin{equation*}
z_{4}=z_{5}=z_{6}=0, \quad x_{1}=0 ; \tag{27}
\end{equation*}
$$

б) free right edge of the plate:

$$
\begin{equation*}
z_{4}=z_{5}=z_{6}=0, \quad x_{1}=I_{1} ; \tag{28,a}
\end{equation*}
$$

в) pinched right edge of the plate:

$$
z_{1}=Z_{2}=z_{3}=0, \quad x_{1}=l_{1}
$$

Thus formulated spectral problem (26-28) in the parameter $\omega$, describing the propagation of flexural waves in a flat edge plate Timoshenko.

## IV. Numerical Analysis of the Dispersion of the Edge Waves in the WedgeShaped Plates

The decision stated above spectral boundary-value problems (14), (15), (16) and (26), (27), (28) was performed by the method of orthogonal sweep Godunov [4]. Numerical implementation of this method was carried out on a computer using software package MAPLE. To test the method and the program was designed version of the album with the boundary conditions can be solved analytically in terms of trigonometric functions.

For resolving the system of equations (14) Kirchhoff-Love plate, these boundary conditions of the form:

$$
\begin{equation*}
X_{1}=0,1 ; \quad z_{2}=z_{4}=0 \tag{29}
\end{equation*}
$$

Here and below we use the dimensionless system of units in which the bandwidth $l$, shear modulus $G$ and bulk density equal to unity.
In this case, the waveform is given by the expression W

$$
\begin{gather*}
z_{1}=z_{o} \cos 2 \pi n x_{1}  \tag{30}\\
z=-(2 \pi n)^{2} z_{1}=A_{2} z_{1} \\
z_{3}=\frac{\left(v K^{2}+(2 \pi n)^{2}\right) h^{3}}{6(1-v)} z=A_{3} z_{1} \\
z_{4}=\left[v k^{2} \frac{\left(v k^{2}+(2 \pi n)^{2}\right) h^{3}}{6(1-v)}+\frac{(1+v)}{6} k^{4}-h\left(\frac{\omega}{g^{*}}\right)^{2}\right] z_{1}=A_{4} z_{1} \\
z_{3}^{1}=-\frac{\left(v k^{2}+(2 \pi n)^{2}\right) h^{3}}{6(1-v)}(2 \pi n) z_{1} \quad z_{2}^{1}=(2 \pi n) z_{2} \\
z_{4}^{1}=-A_{4}(2 \pi n) z_{1}=-(2 \pi n) z_{4}
\end{gather*}
$$

Where $z_{o}$ - arbitrary constant; $c_{n}$ - The real part of the complex frequency; successively substituting the expression (30) into equation (26) we obtain the dispersion equation

$$
\left|\begin{array}{cccc}
2 \pi n & 1 & 0 & 0  \tag{31}\\
v k^{2} & -(2 \pi n) & -\frac{6(1-v)}{n^{3}} & 0 \\
0 & -\frac{n^{3}}{3} k^{2} & -(2 \pi n) & 1 \\
B_{1} & 0 & v k^{2} & 2 \pi n
\end{array}\right|=0
$$

Where ${ }_{B_{1}}=\frac{(1+v) h}{6} k^{4}-n\left(\frac{\omega}{C_{K}+i C_{I}}\right)^{2}$

Similarly, choosing the boundary conditions for the resolution of the system (22) in the form of plates Timoshenko

$$
\begin{equation*}
x=0,1 ; \quad z_{4}=z_{5}=z_{6}=0 \tag{32}
\end{equation*}
$$

Find the expression for the wave form

$$
\begin{array}{ll}
z_{1}=A_{1} \cos 2 \pi n x_{2} ; & z_{4}=A_{4} \sin 2 \pi n x_{2} ; \\
z_{2}=A_{2} \cos 2 \pi n x_{2} ; & z_{5}=A_{5} \sin 2 \pi n x_{2} ;  \tag{33}\\
z_{3}=A_{3} \sin 2 \pi n x_{2} ; & z_{6}=A_{6} \cos 2 \pi n x_{2} .
\end{array}
$$

In (33) permanent $A_{i}(i=1,2,3,4,5,6)$ are determined by solving the system of equations

$$
\left\{\begin{array}{l}
\quad A_{3}+\frac{A_{4}}{\chi h}=0 ; \\
\kappa A_{3}-\frac{12}{h^{3}} A_{5}=0 ; \\
-v \kappa \kappa \kappa_{2}-\frac{6(1-v)}{h^{3}} A_{6}=0 ; \\
\quad \chi h \kappa \kappa_{2}+\kappa^{2}\left(\chi h-h c^{2}\right) A_{1}=0 ;  \tag{34}\\
\quad-\chi h \kappa \kappa_{1}-\left[\frac{(1-v) h^{3}}{6} \kappa^{2}+\chi h-\frac{h^{3}}{12} \omega^{2}\right] A_{2}+v \kappa \kappa_{6}=0 ; \\
\quad-\kappa A_{5}+A_{4}+\frac{h^{3}}{12} \omega^{2} A_{3}=0 .
\end{array}\right.
$$

The system of equations (34) is obtained by substituting (33) in the resolution of the system of differential equations (22). Condition vanishing of the determinant of the system (34) is the dispersion equation boundary value problem (22), (33). The values of the phase velocities found from the above dispersion equations and solving the corresponding test problems (14). (29). (26) and (32) coincide with each other up to the fourth decimal place in the wave number range from 0.1 to 15 for the first two modes $(\mathrm{n}=0.1)$. For the Kirchhoff-Love plates of variable thickness were investigated first five modes with minimum phase velocity of the complexes. Where $C=C_{R}+i C_{I}, C_{R}$ - the phase velocity of wave propagation; $C_{I}$ - speed damping. Figure 3a shows the dispersion curves of the first mode, depending on the thickness varies linearly. Here we assume that the two edges of the plate are free. The straight line I corresponds to a constant thickness $h_{1}=h_{2}=0,1$. In this case, the plate varies as a rod. Curve II - variant $h_{1}=h_{2} / 2=0,05$; curve III - variant $h_{1}=h_{2} / 100=0,001$, curve IV $h_{1}=h_{2} / 1000=0,0001$ and $E_{\min }=6,9 \cdot 10^{6} \kappa / M^{2}, \quad E_{\operatorname{maxs}}=6,9 \cdot 10^{8} \kappa / \mu^{2}, \beta=10^{-4}$. Found that $\kappa>9$ speed damping increase depending on k . For plastics constant thickness $C_{\Gamma}$ on the segment $10^{-4}<C<70$ decreases in a straight line. It can be seen, the dependence of the damping of the wave number starts on the wave number 3-6. With enthusiasm wave number damping factor tends to reduce hand. It can be seen that for a plate of constant thickness, the phase velocity tends to infinity, and for acute wedge plate there is a finite limit as $\kappa \rightarrow \infty$, i.e. the bending edge wave length sufficiently small (compared with the width of the plate) are distributed without dispersion. This fact is evident physical, since the edge of the wedge is no characteristic linear dimension. Land without dispersive waveguide movement begins with a wave of $3-9$, which corresponds to the length of the waves, is less than 1.

It should be noted one fundamental point. Strictly speaking, this study did not consider the case of theoretical $h_{1}=0$ or $\kappa \rightarrow \infty$. All the numerical results obtained by the
simulation of wave processes on a computer that can not operate with an infinitely small and infinitely large numbers. However, the numerical stability can check the result in sufficiently large range of parameters $h_{1}$ or $\kappa$. Despite the lack of theoretical basis, this verification sufficiently suggests that a known controlled precision found the limit value of any quantity at $h_{1} \rightarrow 0$ or $\kappa \rightarrow \infty$. Physically, it is obvious that the parameters $h_{1}$ and $\kappa$ must be coordinated so that the wavelength was substantially greater than the width edge $h_{l}$.

The numerical experiments show that the maximum dimensionless phase velocity (the real part of the complex frequency) during the first mode $K \rightarrow \infty$ largest coincides with the dimensionless thickness $h_{2}$. In dimensional terms, this corresponds to the following changes in the law (the actual number of complex frequency) limit the phase velocity $C_{R 0}$ the angle of the wedge $\varphi_{o}$.

$$
\begin{equation*}
C_{R o}=C_{s} \operatorname{tg} \frac{\varphi_{o}}{2} \tag{35}
\end{equation*}
$$

Coincides with the results of $\left(C_{R o}=C_{o}\right)$ [6]. Numerical experiment also showed that the family of the dispersion curves with different angles at the vertex of the wedge has a certain similarity property, namely: the ratio of the phase velocity to the speed $C_{o}$ (35) does not depend on the angle $\varphi_{o}$. For constant thickness form varies only slightly, while the wedge-shaped plate with increasing K , observed near the localization own form an acute angle. Figure 4 shows the dispersion curves and the second oscillation mode, depending on wavelength in distinguishing values of the thickness of the plate. When $K=0$, the phase velocity is finite. Localization waveform and a limited range of the phase velocity with valid for this mode. Figure 4 b shows the imaginary part of the complex, depending on the speed of the wave hours for different thicknesses. It is seen that the rate of 3-4 imaginary second mode at $K>5$ does not tolerate dispersion. Figure 4 b shows the evolution of the dispersion curves as a function of the wedge angle and thickness $h_{2}$. For small to form close to the line that corresponds to the torsion vibration at large to the observed localization. In contrast to the first mode is available hotspot. Fig. 5 a, b shows the dispersion curves and mode shapes for III and IV of fashion. With integrated small wave numbers phase velocity tends to infinity, and for large - to a finite limit. Also observed localization forms. The number of nodal points two and three, respectively modes (Figure 5).

Figure 6 shows the dependence of the real part of the phase velocities of the first four vibration modes acute wedge plate with different Poisson's ratio. As can be seen from the figure, the maximum phase velocity $C_{o}$ the first mode is virtually independent of Poisson's ratio. In the last phase velocity modes $C_{o}$ increases with increasing v , where the effect of Poisson's ratio is more pronounced at the higher-order modes (real part of complex velocity). Fig. 7 shows the dispersion curves of phase velocities of the four vibration modes for two variants of the legal termination edge of the plate: the free edge (dashed) and fixed (solid). Unlike these options significantly at small wave numbers and virtually absent at large, that is, as one would expect, the maximum phase velocity is independent of the conditions of securing the plate away from the edge of the wedge.

In [7], the distribution of the bending edge waves in the wedge-shaped waveguides in the framework of the linear theory of elasticity. We used the finite element method, based on which the empirical relation for the phase velocities of the normal modes of oscillation depending on the angle of the wedge $\varphi$ :

$$
\begin{equation*}
C_{o}=C_{r} \sin (m \varphi) ; \quad m=1,2, \ldots, \quad m \varphi ; 90^{\circ}, \tag{36}
\end{equation*}
$$

Where: $C_{r}$ - Rayleigh wave speed for a half; $m$ - mode number. It is easy to see that the relations (35) and (36) do not agree with each other at small angles $\varphi$. It is therefore of interest to find out what the limiting phase velocities obtained in the
framework of a more general theory of plates Timoshenko. The spectral problem (2628), which describes the distribution of edge waves plate Timoshenko was solved numerically orthogonal sweep method of Godunov. To control the numerical convergence of the method, the number of points equal to the orthogonalization taken from 10 to 100 . In parallel redundant calculations were carried out in double precision. The result is considered satisfactory if the doubling of the number of points did not change the orthogonalization four significant digits in the phase velocity. Limiting the phase velocity of the first mode for thickness $h_{2}=0$, 2 equal to 0.1945 and is independent of Poisson's ratio. Compared with the same result obtained in the theory of KirchhoffLove, in this case the difference is less than $3 \%$. Figure 8 shows the first three modes indeed part of the complex phase velocity of the plate Timoshenko (b) compared to the corresponding modes of Kirchhoff-Love plate (a) in the Poisson's ratio of 0.25 . In the case of the Kirchhoff-Love plate limit above the phase II and III modes, and with increasing mode number increases contrast. The comparative analysis of the propagation of the edge waves on the basis of these theories plates shows a satisfactory agreement for the first vibration mode. The resulting discrepancy with the results in [5] indicates the need for more detailed research into the general theory of elasticity. Overall, however, conducted a numerical analysis of edge waves in the Kirchhoff-Love plates and Timoshenko suggests that the Kirchhoff-Love hypotheses are justified in the calculation of wave processes in the wedge-shaped plates, including frequencies with a wavelength of the order of the thickness of the plate. This discrepancy with the classical results of the theory of Kirchhoff-Love plates of constant thickness above phenomenon is explained by established localization waveforms with increasing frequency, which occurs only in the plates of variable thickness. At the same time, the relative simplicity of the mathematical apparatus of the theory of Kirchhoff-Love plates, allows us to investigate the dispersion characteristics of the waveguides with a more complex configuration section, which is very difficult to build as part of three-dimensional theory. Consider a plate, whose Thickness varies in accordance

$$
h\left(x_{1}\right)=h_{o} / x_{1} /, \quad-b \leq x_{1} \leq b .
$$

It is clear that such a plate vibrations are reduced to fluctuations in the wedge plate with boundary conditions at $x_{1}=0$, corresponds to the case of symmetry

$$
\begin{equation*}
\varphi=0, \quad Q=0 \tag{37}
\end{equation*}
$$

and of ant symmetry

$$
\begin{equation*}
W=0 \tag{38}
\end{equation*}
$$

$$
M=0
$$

Figure 9.a. and 9.b. (solid lines) shows the dispersion curves of phase velocities of the first three modes in the Kirchhoff-Love plate with a linear variation of thickness.

$$
h\left(x_{1}\right)=h_{0} x_{1}^{p}, \quad o<x_{1} \leq b
$$

where the parameter $p$ taken equal to $1.5 ; 2 ; 2.5 ; 3$ in accordance with designations of curves $1,2,3$ and 4 . For comparison, the dashed lines indicate similar curves discussed above relating to the wedge plate with a thickness $h(1)=h_{o}=0$, 2. Note the qualitative difference in the behavior of solid and dotted lines. When $\mathrm{p}=1$, as mentioned above, the phase velocities approaching asymptoticity nonzero limits, the curve of the first mode increases monotonically. For $\mathrm{p}>1$, the curve of the first mode is not monotonic and has a characteristic maximum in the medium range. Starting with a certain wave number of the phase velocities of all modes decrease monotonically without entering the asymptote nonzero. With increasing $p$ the maximum curve of the first mode is shifted to lower frequencies, and shortwave phase velocities decrease more rapidly. Thus, summarizing the results obtained earlier in the event of a non-linear law of variation of the thickness of the plate, it can be argued that the phase velocity of the first mode in the wedge plate at high frequencies is determined by the rate of change of the thickness in the vicinity of the sharp edge.


Figure 3 : The dispersion curves of the first mode
I. $h_{1}=h_{2}=0,1$; II. $h_{1}=h_{2 / 2}=0,05$; III. $h_{2 / 100}=0,001 ;$ IV. $h_{1}=h_{2 / 1000}=0,001$


Figure 4 a : The dispersion curves of the second mode

1. $h_{2}=0,002, h_{2}=0,2 ; 2 . h_{1}=0,001, h_{2}=0,1 ; 3 . h_{1}=0,0002, h_{2}=0,02$


Figure 4 b : The dependence of the damping rate $\kappa$


Figure 4 V . : Wave form corresponding dispersion curves of the second mode $1 . K=1$; 2. $K=10$


Figure 5a: The dispersion curves of the third mode
$h_{2}=0,002, h_{2}=0,2 ; 2 . h_{1}=0,001, h_{2}=0,1 ; 3 . h_{1}=0,0002, h_{2}=0,02$


Figure 5 b: The dependence of the damping rate $\kappa$

Figure 6 : The dependence of the phase velocity of the first four modal wedge plates with different Poisson's ratios

Figure 7 : The dispersion curves of phase velocities of the four modes for the two variants of the legal termination edge of the plate
-_ free edge,
-------- rigidly fixed



Figure 9a: The dispersion curves of phase velocities at different values P

1. $\mathrm{P}=1,5$;
2. $\mathrm{P}=2$;
3. $\mathrm{P}=2,5 ; \quad$;. $\mathrm{P}=3$


Figure 9.v. : The dispersion curves of phase velocities of the second mode at different values P .

$$
\text { 1. } \mathrm{P}=1,5 ; \quad \text { 2. } \mathrm{P}=2 ; \quad \text { 3. } \mathrm{P}=2,5 ; \quad ; . \mathrm{P}=3
$$

On the basis of these results the following conclusions:

- With increasing wave number of the velocity of propagation is real and the imaginary part of the normal modes in a wedge-shaped (plate) band KirchhoffLove and Timoshenko tend to constant values. At the same time there is the localization movement near the sharp edge of the waveguide.
- For small wedge angles comparison of the results obtained by the KirchhoffLove theory and Timoshenko, shows satisfactory agreement.
- Valid and imaginary parts of the complex phase velocity of the first mode in the wedge plate practically does not depend on the Poisson ratio (change within $0.5 \%)$.
- In the short-range limit value is valid and the imaginary part of the phase velocity of the first mode in the tapered waveguide is determined by the rate of change of the thickness in the vicinity of the sharp edge.
- In wedge-shaped plates with a small angle at the apex of a no dispersive waves propagate with a length not exceeding bandwidth.


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# Loubéré Magic Squares Semigroups and Groups 

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Abstract-This work is a pioneer investigation of semigroups and groups over the Loubéré Magic Squares. By the Loubéré Magic Squares, we understand the magic squares formed by the De La Loubéré Procedure. The set of the Loubéré Magic Squares equipped with the matrix binary operation of addition forms a semigroup if the underlining set so considered is the multi set of natural numbers; and if we consider the multi set of integer numbers as the underlined set of entries of the square, the set of the squares enclosed with the aforementioned operation forms an abelian group. The Loubéré Magic Squares are always recognized with centre piece $C$ and magic sum $M(S)$. We showcase that the set of the centre pieces and the set of the magic sums form respective abelian groups if both are equipped with integer numbers operation of addition. We also explicate that the set of the eigen values of the squares enclosed with the integer addition (operation) forms an abelian group. We reveal that the subelement (a terminology we introduced) Magic Squares of the Loubéré Magic Squares forms a semigroup and the Subelement Magic Squares of the Loubéré Magic Squares Group forms a group, with respect to the matrix binary operation of addition.

Keywords: semigroup, group, centre piece, eigen values, subelement, magic sum.
GJSFR-F Classification : FOR Code : MSC 2010: 16 W 22

Strictly as per the compliance and regulations of :


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# Loubéré Magic Squares Semigroups and Groups 

Babayo A. M. ${ }^{\alpha}$ \& G. U. Garba ${ }^{\sigma}$


#### Abstract

This work is a pioneer investigation of semigroups and groups over the Loubéré Magic Squares. By the Loubéré Magic Squares, we understand the magic squares formed by the De La Louberé Procedure. The set of the Loubéré Magic Squares equipped with the matrix binary operation of addition forms a semigroup if the underlining set so considered is the multi set of natural numbers; and if we consider the multi set of integer numbers as the underlined set of entries of the square, the set of the squares enclosed with the aforementioned operation forms an abelian group. The Loubéré Magic Squares are always recognized with centre piece $C$ and magic sum $M(S)$. We showcase that the set of the centre pieces and the set of the magic sums form respective abelian groups if both are equipped with integer numbers operation of addition. We also explicate that the set of the eigen values of the squares enclosed with the integer addition (operation) forms an abelian group. We reveal that the subelement (a terminology we introduced) Magic Squares of the Loubéré Magic Squares forms a semigroup and the Subelement Magic Squares of the Loubéré Magic Squares Group forms a group, with respect to the matrix binary operation of addition.


Keywords: semigroup, group, centre piece, eigen values, subelement, magic sum.

## I. Introduction

This pioneering work disclosed a new realm of semigroup and group, the Loubéré Magic Squares Semigroup and Group. The set of the Loubéré Magic Squares of the arithmetic sequence of the set of the natural numbers or of its multi set form a semigroup which by analogy we refer to as the Loubéré Magic Squares Semigroup; and the set of the Loubéré Magic Squares of the arithmetic sequence of the set of integer numbers or of the multi set of the integer numbers form a group which by analogy we refer to as the Loubéré Magic Squares Group. The aforementioned semigroup [3] and group [4]are both with respect to the matrix binary operation of addition, thus they are both additive.

The collection of the centre pieces with formula $c_{n}=a_{n}+\left(\frac{m-1}{2}\right) j_{n}$ equipped with the integer addition forms an abelian group and the set of all the magic sums with formula $M\left(S_{n}\right)=\frac{m}{2}\left[2 a_{n}+(m-1) j_{n}\right]$ equipped with the integer numbers binary operation of addition form an abelian group also, where $n=1,2,3, \ldots$ and $a_{n}, j_{n}$ are the corresponding first term and common difference along the main column respectively of $m \times m$ Loubéré Magic Squares.

We also showcase that the set of eigen values of the Loubéré Magic Squares enclosed with integer numbers operation of addition forms an abelian group. This is meaningful for the principal value of the eigen value corresponds to the magic sum [1].

## Definition 1.1.

A magic square $n-1 \times n-1$ formed by removing the border cells of an $n \times n$ Loubéré Magic Square is called the subelement magic square of the $n \times n$ Loubéré Magic Square.
Remarks 1.1.
We have interest in the least subelement which is a subset of $3 \times 3$ Pancolumn Magic Squares. Purposefully, the $3 \times 3$ Loubéré Magic Square has no subelement for it is not a pancolumn. We explicate that the subelement magic squares of the Loubéré Magic Squares Semigroup forms a semi group and the subelement magic squares of the Loubéré Magic Squares Group forms a group with respect to the same underlining set and operation.

## II. Preliminaries

A basic magic square of order $n$ is an arrangement of arithmetic sequence of common difference of 1 from 1 to $n^{2}$ in an $n \times n$ square grid of cells such that every row, column and diagonal add up to the same number, called the magic sum $\mathrm{M}(\mathrm{S})$ expressed as $M(S)=\frac{n^{3}+n}{2}$ and a centre piece C as $C=\frac{M(S)}{n}$.

## a) Loubéré Procedure ( $N E-W-S$ or $N W-E-S$, the cardinal points)

Consider an empty $n \times n$ square of grids (or cells). Start, from the central column or row at a position $\left\lfloor\frac{n}{2}\right\rfloor$ where $\lfloor\overline{1}]$ is the greater natural number less than or equal to, with the number 1. The fundamental movement for filling the square is diagonally up, right (clock wise or NE or SE) or up left (anti clock wise or NW or SW) and one step at a time. If a filled cell (grid) is encountered, then the next consecutive number moves vertically down ward one square instead. Continue in this fashion until when a move would leave the square, it moves due N or E or W or S ( depending on the position of the first term of the sequence) to the last row or first row or first column or last column.

## Definition 2.1

Main Row or Column is the column or row of the Loubéré Magic Squares containing the first term and the last term of the arithmetic sequence in the square.
b) The Proof of the $\left\lfloor\frac{m^{2}}{2}\right\rfloor=a+\left(\frac{m-1}{2}\right) j$ and of the $M(S)=\frac{m}{2}[2 a+(m-1) j]$, where $j=\frac{l-a}{m-1}$ Theorem 2.1.

Let the arithmetic sequence $a, a+d, \ldots, l=a+(n-1) d$ be arranged in an $m \times m$ Loubéré Magic Square. Then the magic sum of the square is expressed as $M(S)=\frac{m}{2}[2 a+(m-1) j]$ and the middle term of the sequence (centre piece of the square) is expressed as $C=a+\left(\frac{m-1}{2}\right) j$ where j denotes the common difference of entries along the main column or row and is given as $j=\frac{l-a}{m-1}$.
Proof.
Consider any arbitrary General Loubéré Magic Square (here we consider $3 \times 3$ ) as follows:

| $c+b$ | $c-b-d$ | $c+d$ |
| :---: | :---: | :---: |
| $c-b+d$ | $c$ | $c+b-d$ |
| $c-d$ | $c+b+d$ | $c-b$ |

Let $a=c-b-d$ and $l=c+b+d$. Then we have (from the square) an arithmetic sequence: $c-b-d, c-b, \ldots, c+b+d$ having the sums $S$ as

$$
\begin{gathered}
S=(c-b-d)+(c-b)+\cdots+(c+b)+(c+b+d) \rightarrow(1) \\
+ \\
S=(c+b+d)+(c+b)+\cdots+(c-b)+(c-b-d) \rightarrow(2) \\
\hline \text { Adding (1)and (2), } 2 \mathrm{~s}=2 \mathrm{c}+2 \mathrm{c}+\cdots \quad \mathrm{n} \text { times }
\end{gathered}
$$

i.e. $2 s=2 n c \Rightarrow c=\frac{s}{n} \ldots$ (3) and $s=\frac{n}{2}(a+l) \ldots$ (4) from the Gaussian High School (Elementary) Method. Since our square is $m \times m$, m number of cells (terms) are on the main column whence $a=c-b-d$. Thus, (3) and (4) become $C=\frac{M(S)}{m} \ldots$ (5) and $M(S)=\frac{m}{2}[a+l] \ldots$ (6) respectively. And, $l=a+(m-1) j \ldots$ (7) where $j$ is along the main column. Substituting (7) in (6), we have: $M(S)=\frac{m}{2}[2 a+(m-1) j] \ldots$ (8). Substituting (8) in (5), we get: $C=a+\left(\frac{m-1}{2}\right) j \ldots$ (9) From (3) and (4), $C=\frac{1}{2}(a+l)=$ $\left(a-\frac{a}{2}\right)+\frac{l}{2}=a+\frac{(l-a)}{2}=a+\frac{l-a}{m-1} \frac{m-1}{2}$, i.e. $C=a+\left(\frac{m-1}{2}\right) \frac{l-a}{m-1} \ldots$ (10). Comparing (9) and (10), we have: $j=\frac{l-a}{m-1} \ldots$ (11).

## Definition 2.2.

A non empty set $S$ equipped with a binary operation * is said to be a Semigroup $\left(S,{ }^{*}\right)$ if it satisfies the following axioms:
i. $\quad a, b \in S \Rightarrow a * b \in S$; and
ii. $a, b, c \in S \Rightarrow a *(b * c)=(a * b) * c$.

If in addition to the 2 axioms above, the following axioms are satisfied; then we call the algebraic structure a group $(G, *)$.
iii. $\exists e \in S \quad \ni a * e=e * a \forall a \in S$; and
iv. $\forall a \in S, \exists a^{-1} \in S \quad \ni a * a^{-1}=a^{-1} * a=e \in S$.

If in addition to the above 4 axioms: I; ii; iii; and iv; the following axiom is satisfied; then we call (G, ${ }^{*}$ ) an abelian group.
v. $\forall a, b \in S, a * b=b * a$

Remark 2.1. The shift in notations from the use of S to G is intentional by the respective specialists.

## iII. The Loubéré Magic Souares Semigroup and Groups

We hereby present that the set of Loubéré Magic Squares $L$ over the set of natural numbers equipped with the matrix binary operation of addition $\oplus$ forms a semigroup, and over the set of integer numbers forms a group-enclosed with the same operation.
a) Definition 3.1.

The square of grid of cells $\left[a_{i j}\right]_{n \times n}$ is said to be Loubéré Magic Square if the following conditions are satisfied.
i. $\quad \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}=k$
ii. $\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\mathrm{k}$
iii. $a_{1,\left\lceil\frac{n}{2}\right.}, a_{\left\lceil\frac{n}{2}\right\rceil\left\lceil\frac{n}{2}\right\rceil}, a_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $a_{\left\lceil\frac{n}{2}, n\right.}, a_{\left.\left.\left\lceil\frac{n}{2}\right\rceil \right\rvert\, \frac{n}{2}\right\rceil}, a_{\left\lceil\frac{n}{2}, 1\right.}$ are on the same main column or row,
where [] is the greater integer less or equal to, T is the transpose (of the square), k is the magic sum (magic product is defined analogously) usually expressed as $k=\frac{n}{2}[2 a+(n-1) j]-$ from the sum of arithmetic sequence, where j is the common difference along the main column or row and a is the first term of the sequence- and $a_{\left.\left.\left\lceil\frac{n}{2}\right\rceil \right\rvert\, \frac{n}{2}\right\rceil}=\frac{k}{n}$.

## b) Theorem 3.2.

$(L, \oplus)$ forms an Infinite Commutative Semigroup if the underlining multi set is of natural numbers and it forms an Infinite Additive Abelian Group if the underlining multi set is of integer numbers.
Proof. Let $\left[a_{i, j}\right]_{n \times n}$ and $\left[b_{i j}\right]_{n \times n} \in L$. Then, by Definition 3.1, $\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}=k$, $\operatorname{trace}\left[a_{i j}\right]_{n \times n}=\operatorname{trace}\left[a_{i j}\right]_{n \times n}^{T}=k$, and $a_{1,\left[\frac{n}{2}\right]}, a_{\left[\frac{n}{2}|,| \frac{n}{2}\right]}, a_{n,\left[\frac{n}{2}\right]}$ are on the same main column or row and $a_{\left[\frac{n}{2}\right], n}, a_{\left[\frac{n}{2}\left\lceil\frac{n}{2}\right]\right.}, a_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row, and $\sum_{i=1}^{n} \sum_{j=1}^{n} b_{i j}=l$, $\operatorname{trace}\left[b_{i j}\right]_{n \times n}=\operatorname{trace}\left[b_{i j}\right]_{n \times n}^{T}=l$, and $b_{1,\left[\frac{n}{2}\right\rceil}, b_{\left.\left[\frac{n}{2}\right\rceil, \left\lvert\, \frac{n}{2}\right.\right]}, b_{n,\left[\frac{n}{2}\right]}$ are on the same main column or row and $b_{\left[\frac{n}{2}\right], n}, b_{\left[\frac{n}{2}\right]\left[\frac{n}{2}\right]}, b_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row.
Then,

$$
\begin{aligned}
\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}+ & \sum_{i=1}^{n} \sum_{j=1}^{n} b_{i j}=k+l=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}+\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}} \\
& \operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}+\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\mathrm{k}+1, \quad \text { and } \quad a_{1,\left[\frac{n}{2}\right]}+
\end{aligned}
$$

$b_{1,\left\lceil\frac{n}{2}\right\rceil} a_{\left.\left.\left\lceil\frac{n}{2}\right\rceil \right\rvert\, \frac{n}{2}\right\rceil}+b_{\left\lceil\frac{n}{2},\left\lceil\frac{n}{2}\right\rceil\right.}, a_{n,\left\lceil\frac{n}{2}\right\rceil}+b_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $a_{\left\lceil\frac{n}{2}\right\rceil, n}+$ $b_{\left\lceil\frac{n}{2}\right], n}, a_{\left\lceil\frac{n}{2}\left\lceil\left[\frac{n}{2}\right\rceil\right.\right.}+b_{\left[\frac{n}{2}\left\lceil\left\lceil\frac{n}{2}\right\rceil\right.\right.}, a_{\left\lceil\frac{n}{2}\right\rceil, 1}+b_{\left\lceil\frac{n}{2}\right\rceil, 1}$ since $a_{1,\left\lceil\frac{n}{2}\right\rceil}, a_{\left.\left.\left\lceil\frac{n}{2}\right\rceil \right\rvert\, \frac{n}{2}\right\rceil}, a_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $a_{\left[\frac{n}{2}\right], n}, a_{\left[\frac{n}{2}\left\lceil\left[\frac{n}{2}\right]\right.\right.}, a_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row, and $b_{1,\left\lceil\frac{n}{2}\right]} b_{\left[\frac{n}{2},\left\lceil\frac{n}{2}\right]\right.}, b_{n,\left\lceil\frac{n}{2}\right]}$ are on the same main column or row and $b_{\left[\frac{n}{2}\right], n}, b_{\left\lceil\frac{n}{2}\left\lceil\left[\frac{n}{2}\right]\right.\right.}, b_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row.

## i. Associativity

Let $\left[a_{i, j}\right]_{n \times n},\left[b_{i, j}\right]_{n \times n}$ and $\left[c_{i, j}\right]_{n \times n} \in L$. Then, by Definition 3.1, $\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}=$ $k$, $\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\mathrm{k}$, and $a_{1, \left\lvert\, \frac{n}{2}\right.,}, a_{\left[\frac{n}{2}|,| \frac{n}{2}\right]}, a_{n,\left[\frac{n}{2}\right]}$ are on the same main column or row and $a_{\left\lceil\frac{n}{2}\right\rceil, n}, a_{\left\lceil\frac{n}{2}\left\lceil\left\lceil\frac{n}{2}\right\rceil\right.\right.}, a_{\left\lceil\frac{n}{2}\right], 1}$ are on the same main column or row, $\sum_{i=1}^{n} \sum_{j=1}^{n} b_{i j}=l$, $\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=1 \quad$ and $b_{1,\left[\frac{n}{2}\right.}, b_{\left[\frac{n}{2}|,| \frac{n}{2}\right]}, b_{\left.n, \frac{n}{2}\right]}$ are on the same main column or row and $b_{\left\lceil\frac{n}{2}\right\rceil, n}, b_{\left\lceil\frac{n}{2}\left\lceil\left\lceil\frac{n}{2}\right\rceil\right.\right.}, b_{\left[\frac{n}{2}\right\rceil, 1}$ are on the same main column or row.
and $\sum_{i=1}^{n} \sum_{j=1}^{n} c_{i j}=m$, $\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\mathrm{m} \quad$ and $c_{1,\left[\frac{n}{2}\right]}, c_{\left[\frac{n}{2}\right],\left[\frac{n}{2}\right]}, c_{n,\left[\frac{n}{2}\right]}$ are on the same main column or row and $c_{\left[\frac{n}{2}, n\right.}, c_{\left[\frac{n}{2}\right]\left[\frac{n}{2}\right]}, c_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row Then,

$$
\begin{aligned}
\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}+\left(\sum_{i=1}^{n}\right. & \left.\sum_{j=1}^{n} b_{i j}+\sum_{i=1}^{n} \sum_{j=1}^{n} c_{i j}\right)=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\left(\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}\right) \\
= & \operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}} \\
= & \left(\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}\right)+\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\left(\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}+\sum_{i=1}^{n} \sum_{j=1}^{n} b_{i j}\right)+\sum_{i=1}^{n} \sum_{j=1}^{n} c_{i j} \\
& \operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\left(\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}\right)=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}+\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=
\end{aligned}
$$

$\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\left(\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\operatorname{trace}\left[\mathrm{c}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}\right)=\mathrm{k}+\mathrm{l}+\mathrm{m}, \quad$ and $a_{\left.1, \left\lvert\, \frac{n}{2}\right.\right]}+\left(b_{\left.1, \left\lvert\, \frac{n}{2}\right.\right\rceil}+\underset{\underset{\sim}{2}}{\frac{n}{2}}\right.$ $c_{\left.1,\left\lceil\frac{n}{2}\right\rceil\right)}, a_{\left\lceil\frac{n}{2} \left\lvert\,\left\lceil\left\lceil\frac{n}{2}\right\rceil\right.\right.\right.}+\left(b_{\left\lceil\frac{n}{2}\right\rceil,\left\lceil\frac{n}{2}\right\rceil}+c_{\left\lceil\frac{n}{2}\right\rceil,\left\lceil\frac{n}{2}\right\rceil}\right), a_{n,\left\lceil\frac{n}{2}\right\rceil}+\left(b_{n,\left\lceil\frac{n}{2}\right\rceil}+c_{n,\left\lceil\frac{n}{2}\right\rceil}\right)$ are on the same main column or row and $a_{\left[\frac{n}{2}\right], n}+\left(b_{\left[\frac{n}{2}\right], n}+c_{\left[\frac{n}{2}\right], n}\right), a_{\left[\frac{n}{2}\right\rceil\left\lceil\frac{n}{2}\right\rceil}+\left(b_{\left[\frac{n}{2}\left\lceil\left[\frac{n}{2}\right]\right.\right.}+c_{\left[\frac{n}{2}\right\rceil\left[\frac{n}{2}\right]}\right), a_{\left[\frac{n}{2}\right], 1}+{ }_{-}^{65}$ $\left(b_{\left[\frac{n}{2}\right], 1}+c_{\left\lceil\frac{n}{2}\right\rceil, 1}\right)$, then $\left(a_{\left\lceil\frac{n}{2}, n\right.}+b_{\left\lceil\frac{n}{2}\right], n}\right)+c_{\left\lceil\frac{n}{2}, n\right.},\left(a_{\left\lceil\frac{n}{2}\left\lceil\left\lceil\frac{n}{2}\right\rceil\right.\right.}+b_{\left\lceil\frac{n}{2}\right\rceil\left\lceil\frac{n}{2}\right\rceil}\right)+c_{\left\lceil\frac{n}{2} \left\lvert\,\left\lceil\frac{n}{2}\right\rceil\right.\right.},\left(a_{\left\lceil\frac{n}{2}\right], 1}+\right.$ $\left.b_{\left[\frac{n}{2}, 1\right.}\right)+c_{\left[\frac{n}{2}, 1\right.}$ and $a_{\left.1, \frac{n}{2}\right]}, a_{\left[\frac{n}{2}\right],\left\lceil\frac{n}{2}\right]}, a_{n,\left[\frac{n}{2}\right]}$ are on the same main column or row and $a_{\left\lceil\frac{n}{2}\right\rceil, n}, a_{\left\lceil\frac{n}{2}\left\lceil\left[\frac{n}{2}\right\rceil\right.\right.}, a_{\left\lceil\frac{n}{2}\right\rceil, 1}$ are on the same main column or row, and $b_{1,\left\lceil\frac{n}{2}\right\rceil}, b_{\left.\left.\left\lceil\frac{n}{2}\right\rceil \right\rvert\, \frac{n}{2}\right\rceil}, b_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $b_{\left[\frac{n}{2}\right], n}, b_{\left[\frac{n}{2}\right]\left[\frac{n}{2}\right]}, b_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row, and $c_{1,\left\lceil\frac{n}{2}\right\rceil}, c_{\left\lceil\frac{n}{2}\right\rceil\left\lceil\frac{n}{2}\right]}, c_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $c_{\left[\frac{n}{2}\right], n}, c_{\left[\frac{n}{2} \left\lvert\,\left\lceil\frac{n}{2}\right]\right.\right.}, c_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row.

## ii. Identity Element

$\exists\left[a_{i j}\right]_{n \times n} \in L$ and is said to be Loubéré Magic Square if the following conditions are satisfied.

$$
\begin{gathered}
\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}=0 \\
\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=0
\end{gathered}
$$

$a_{\left.1, \left\lvert\, \frac{n}{2}\right.\right\rceil}=0, a_{\left\lceil\frac{n}{2}\right\rceil,\left\lceil\frac{n}{2}\right\rceil}=0, a_{n,\left\lceil\frac{n}{2}\right\rceil}=0$ are on the same main column or row and $a_{\left\lceil\frac{n}{2}, n\right.}=$ 0 , $a_{\left\lceil\frac{n}{2} \left\lvert\,\left\lceil\frac{n}{2}\right\rceil\right.\right.}=0, a_{\left\lceil\frac{n}{2}, 1\right.}=0$ are on the same main column or row, whence $a_{i j}=0, \forall \mathrm{i}, \mathrm{j} \Rightarrow$ the identity is $[0]_{n \times n} \in L$
iii. Inverse Element Property

Given $\left[a_{i, j}\right]_{n \times n} \in L \ni \sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}=k$, $\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\mathrm{k}$, and
$a_{\left.1, \left\lvert\, \frac{n}{2}\right.\right]}, a_{\left.\left[\frac{n}{2}\right], \left\lvert\, \frac{n}{2}\right.\right]}, a_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $a_{\left[\frac{n}{2}\right], n}, a_{\left[\frac{n}{2}\left\lceil\left[\frac{n}{2}\right]\right.\right.}, a_{\left[\frac{n}{2}\right], 1}$ are on the same main column or row, there exists $\left[-a_{i, j}\right]_{n \times n} \in L$ such that $\sum_{i=1}^{n} \sum_{j=1}^{n}-a_{i j}=$ $-k$, $\operatorname{trace}\left[-\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[-\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=-\mathrm{k}$, and $-a_{1,\left\lceil\frac{n}{2}\right\rceil},-a_{\left\lceil\frac{n}{2}, \left\lvert\, \frac{n}{2}\right.\right\rceil},-a_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same
main column or row and $-a_{\left[\frac{n}{2}, n\right.},-a_{\left[\frac{n}{2}\right]\left[\frac{n}{2}\right]},-a_{\left[\frac{n}{2}, 1\right.}$ are on the same main column or row. Thus, $\left[a_{i j}\right]_{n \times n}+\left[-a_{i j}\right]_{n \times n}=\left[-a_{i j}\right]_{n \times n}+\left[a_{i j}\right]_{n \times n}=[0]_{\mathrm{n} \times \mathrm{n}}$.
iv. Commutativity
$\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}+\sum_{i=1}^{n} \sum_{j=1}^{n} b_{i j}=k+l=l+k=\sum_{i=1}^{n} \sum_{j=1}^{n} b_{i j}+\sum_{i=1}^{n} \sum_{j=1}^{n} a_{i j}$ and $\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}+\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=k+l=l+k=\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}+\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}=\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}+$ $\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}$
$\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}+\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=\mathrm{k}+\mathrm{l}=\mathrm{l}+\mathrm{k}=\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}+\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}=$ $\operatorname{trace}\left[\mathrm{b}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}+\operatorname{trace}\left[\mathrm{a}_{\mathrm{ij}}\right]_{\mathrm{n} \times \mathrm{n}}^{\mathrm{T}}$, and $a_{1,\left\lceil\frac{n}{2}\right\rceil}+b_{1,\left[\frac{n}{2}\right\rceil}, a_{\left.\left\lceil\frac{n}{2}\right\rceil|,| \frac{n}{2}\right\rceil}+b_{\left\lceil\frac{n}{2}|,| \frac{n}{2}\right\rceil}, a_{n,\left\lceil\frac{n}{2}\right\rceil}+b_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $a_{\left[\frac{n}{2}\right] n}+b_{\left[\frac{n}{2}\right], n}, a_{\left[\frac{n}{2}\left\lceil\left[\frac{n}{2}\right\rceil\right.\right.}+b_{\left[\frac{n}{2}\left\lceil\left[\frac{n}{2}\right]\right.\right.}, a_{\left[\frac{n}{2}\right], 1}+b_{\left[\frac{n}{2}\right], n}$ as well as $b_{1,\left\lceil\frac{n}{2}\right\rceil}+a_{1,\left\lceil\frac{n}{2}\right\rceil}, b_{\left\lceil\frac{n}{2}\right\rceil,\left\lceil\frac{n}{2}\right\rceil}+a_{\left\lceil\frac{n}{2}, \left\lvert\, \frac{n}{2}\right.\right\rceil}, b_{n,\left\lceil\frac{n}{2}\right\rceil}+a_{n,\left\lceil\frac{n}{2}\right\rceil}$ are on the same main column or row and $b_{\left[\frac{n}{2}\right], n}+a_{\left[\frac{n}{2}, n\right.}, b_{\left\lceil\frac{n}{2}\right\rceil\left[\frac{n}{2}\right\rceil}+a_{\left.\left.\left\lceil\frac{n}{2}\right\rceil \right\rvert\, \frac{n}{2}\right]}, b_{\left[\frac{n}{2}, 1\right.}+a_{\left[\frac{n}{2}\right], 1}$.

We can now consider the general multi set. The Loubéré Magic Squares over the multi set of integer numbers, since multi set of natural numbers is its subset, is a semi pandiagonal. By semi pandiagonal, we mean in $n \times n$ square, n elements repeats on every row, column and on a diagonal. Though the sum of the numbers on the rows, the columns and the diagonals add up to the magic sum; yet one diagonal has an $n$ repetition of one element. To change the orientation (from left to right or the reverse) of the pandiagonal of the $3 \times 3$, use the sequence $a, a, a, b, b, b, c, c, c$ rather than $a, b, c, a, b, c, a, b, c$.

We can now show that they form a group as in the above. Consider 3 arbitrary elements of the set of Lefty Semi Pandiagonal Loubéré Magic Squares,

| $b$ | $a$ | $c$ |
| :--- | :--- | :--- |
| $c$ | $b$ | $a$ |
| $a$ | $c$ | $b$ |


| $e$ | $d$ | $f$ |
| :--- | :--- | :--- |
| $f$ | $e$ | $d$ |
| $d$ | $f$ | $e$ |


| $h$ | $g$ | $i$ |
| :--- | :--- | :--- |
| $i$ | $h$ | $g$ |
| $g$ | $i$ | $h$ |

Then,

| $b$ | $a$ | $c$ |
| :--- | :--- | :--- |
| $c$ | $b$ | $a$ |
| $a$ | $c$ | $b$ |


| $e$ | $d$ | $f$ |
| :--- | :--- | :--- |
| $f$ | $e$ | $d$ |
| $d$ | $f$ | $e$ |


| $b+e$ | $a+d$ | $c+f$ |
| :---: | :---: | :---: |
| $c+f$ | $b+e$ | $a+d$ |
| $a+d$ | $c+f$ | $b+e$ |

i. is also a Lefty Semi Pandiagonal $3 \times 3$ Loubéré Magic Squares, hence closure property is satisfied.
ii. Associativity. It is clear (even from inherited property of the underlining set) that

iii. The identity is -as in the above-

| 0 | 0 | 0 |
| :--- | :--- | :--- |
| 0 | 0 | 0 |
| 0 | 0 | 0 |

iv. Let the following be an arbitrary Lefty Semi Pandiagonal Magic Square.

| V | U | W |
| :---: | :---: | :---: |
| W | $V$ | $U$ |
| $U$ | $W$ | $V$ |

Clearly, its inverse is

| $-V$ | $-U$ | $-W$ |
| :--- | :--- | :--- |
| $-W$ | $-V$ | $-U$ |
| $-U$ | $-W$ | $-V$ |

v. Every 2 Loubéré Magic Squares (whether semi pancolumn or not) over multi set of naturals or over multi set of integer numbers commute since natural and integer numbers commutes.

Thus the group and the semigroups of the Loubéré Magic Squares are commutative.

## IV. Centre Pieces and Magic Sums Abelian Groups

## a) Centre Pieces Abelian Group

The set of the centre pieces $c_{1}, c_{2}, c_{3}, \ldots$ of $m \times m$ Loubéré Magic Squares equipped with the integer number binary operation of addition forms an infinite abelian group. Given the centre pieces $c_{1}, c_{2}, c_{3}, \ldots$ of $m \times m$ Loubéré Magic Squares with corresponding formula

$$
c_{1}=a_{1}+\left(\frac{m-1}{2}\right) j_{1}, c_{2}=a_{2}+\left(\frac{m-1}{2}\right) j_{2}, c_{3}=a_{3}+\left(\frac{m-1}{2}\right) j_{3}, \ldots \text {; then }
$$

i. $\quad c_{1}+c_{2}=\left(a_{1}+a_{2}\right)+\left(\frac{m-1}{2}\right)\left(j_{1}+j_{2}\right)$ is the centre piece of the $m \times m$ Loubéré Magic Square with first term $a_{1}+a_{2}$ and common difference along the main column $j_{1}+j_{2}$. Hence, the set is closed.
ii. This is an inherited property of the set of integer numbers:

$$
c_{1}+\left(c_{2}+c_{3}\right)=\left(a_{1}+a_{2}+a_{3}\right)+\left(\frac{m-1}{2}\right)\left(j_{1}+j_{2}+j_{3}\right)=\left(c_{1}+c_{2}\right)+c_{3}
$$

iii. The identity element is the zero centre piece e.g.

| $C$ | $-D$ | $A$ |
| :---: | :---: | :---: |
| $-B$ | 0 | $B$ |
| $-A$ | $D$ | $-C$ |

iv. Given an arbitrary centre piece $c_{n}=a_{n}+\left(\frac{m-1}{2}\right) j_{n}$ of the $m \times m$ Loubéré Magic Square, there exists another centre piece $c_{-n}$ of another $m \times m$ Loubéré Magic Square having first term as $-a_{n}$ and common difference along the main column or row as $-j_{n}$, thus its formula is $c_{-n}=-a_{n}+\left(\frac{m-1}{2}\right)\left(-j_{n}\right)$ such that $c_{n}+c_{-n}=c_{-n}+c_{n}=$ $\left(a_{n}-a_{n}\right)+\left(\frac{m-1}{2}\right)\left[j_{n}-j_{n}\right]=0=c_{i}$, the identity centre piece.
v. Clearly $c_{1}+c_{2}=a_{1}+a_{2}+\left(\frac{m-1}{2}\right)\left(j_{1}+j_{2}\right)=a_{2}+a_{1}+\left(\frac{m-1}{2}\right)\left(j_{2}+j_{1}\right)=c_{2}+c_{1}$

The set equipped with the operation is an abelian group.

## b) Magic Sum Abelian Groups

The set of the magic sums $M\left(s_{1}\right), M\left(s_{2}\right), M\left(s_{3}\right), \ldots$ of $m \times m$ Loubéré Magic Squares equipped with the integer binary operation of addition form an infinite abelian group. Given the magic sums $M\left(s_{1}\right), M\left(s_{2}\right), M\left(s_{3}\right)$, ..of $m \times m$ Loubéré Magic Squares with corresponding formula

$$
M\left(s_{1}\right)=\frac{m}{2}\left[2 a_{1}+(m-1) j_{1}, M\left(s_{2}\right)=\frac{m}{2}\left[2 a_{2}+(m-1) j_{2}\right], M\left(s_{3}\right)=\frac{m}{2}\left[2 a_{3}+(m-1) j_{3}, \ldots ;\right.\right.
$$

then (as in the above):
i. $M\left(s_{1}\right)+M\left(s_{2}\right)=M\left(s_{?}\right)$ where $M\left(s_{?}\right)$ is a magic sum of another $\mathrm{m} \times \mathrm{m}$ Loubéré Magic Square with first term $a_{1}+a_{2}$ and common difference along the main column as $j_{1}+j_{2}$.
The axioms: ii, iii, iv andv follow, by analogy to the centre piece abelian group, immediately.

## V. Eigen Values Abelian Group

The Eigen values computation in the magic squares is what is zealotly prophesized that magic squares are special type of matrices, hence the definition of the magic squares, we do not love to like such a sudden conclusion if loving to liking forces choosing the definitions in terms of just the square grids (or cells).

We want to show through concrete examples that the set of Eigen Values of the Loubéré Magic Squares with the usual integer numbers binary operation of addition forms a group. Consider the following arbitrary two $3 \times 3$ Loubéré Magic Squares which we let


1


We compute the eigen values for $a$ as follows: The corresponding matrix of $a$ is $(a)=\left(\begin{array}{ccc}4 & -3 & 2 \\ -1 & 1 & 3 \\ 0 & 5 & -2\end{array}\right)$, its eigen vector is $|a-\lambda I|=\left|\begin{array}{ccc}4-\lambda & -3 & 2 \\ -1 & 1-\lambda & 3 \\ 0 & 5 & -2-\lambda\end{array}\right|=0$, i. e. $\lambda^{3}-3 \lambda^{2}-24 \lambda-72=(\lambda-3)\left(\lambda^{2}-24\right)=0$ having characteristic equation as $\lambda_{a_{1}}=3, \lambda_{a_{2}}=4.9$ and $\lambda_{a_{3}}=-4.9$.

We compute the eigen values for $b$ as follows: The corresponding matrix of $b$ is $(b)=\left(\begin{array}{ccc}2 & -5 & 0 \\ -3 & -1 & 1 \\ -2 & 3 & -4\end{array}\right)$, $|b-\lambda I|=\left|\begin{array}{ccc}2-\lambda & -5 & 0 \\ -3 & -1-\lambda & 1 \\ -2 & 3 & -4-\lambda\end{array}\right|=0$ i.e. $\lambda^{3}+3 \lambda^{2}-24 \lambda-72=(\lambda+3)\left(\lambda^{2}-24\right)=0$ with eigen values $\lambda_{b_{1}}=-3, \lambda_{b_{2}}=4.9$ and $\lambda_{b_{3}}=-4.9$.

We compute the eigen values for c as follows: The corresponding matrix of c is (c) $=\left(\begin{array}{ccc}6 & -8 & 2 \\ -4 & 0 & 4 \\ -2 & 8 & -6\end{array}\right)$, its characteristic equation is $|c-\lambda I|=\left|\begin{array}{ccc}6-\lambda & -8 & 2 \\ -4 & -\lambda & 4 \\ -2 & 8 & -6-\lambda\end{array}\right|=$ 0 , i.e. $\lambda^{3}-96 \lambda=0$ with corresponding eigen values $\lambda_{c_{1}}=0, \lambda_{c_{2}}=9.8$ and $\lambda_{c_{3}}=-9.8$.
We now conclude this session by showing that the set of eigen values satisfies The Properties of a Group as follows:
Closure Property. Consider any 3 arbitrary Loubéré Magic Squares a, b, c; such that $a+b=c$; then from the example above, the corresponding eigen values of $a$; $\lambda_{a_{1}}, \lambda_{a_{2}} \lambda_{a_{3}}$; the corresponding eigen values of $\mathrm{b} ; \lambda_{b_{1}}, \lambda_{b_{2}}, \lambda_{b_{3}}$; are such that $\lambda_{a_{1}}+\lambda_{b_{1}}=$ $\lambda_{c_{1}}, \lambda_{a_{2}}+\lambda_{b_{2}}=\lambda_{c_{2}}$, and $\lambda_{a_{3}}+\lambda_{b_{3}}=\lambda_{c_{3}}$ where $\lambda_{c_{1}}, \lambda_{c_{2}}, \lambda_{c_{3}}$ are the corresponding eigen values of c .

Associativity Property. Since Loubéré Magic Squares are a semigroup (which is easy to observe), the eigen values are associative.
Identity Element Property. The eigen value 0 is the identity element that corresponds to the sum of the Loubéré Magic Squares of opposite eigen values as in the above.
Inverse Elements Property. For any arbitrary eigen value $\lambda_{m}$ corresponding to a Loubéré Magic Square $m$, there exist a $-\lambda_{m}$ eigen value corresponding to another Loubéré Magic Square such that $\lambda_{m}+\left(-\lambda_{m}\right)$ gives the identity element which is formed as a result of matrix addition of the aforementioned Loubéré Magic Squares.
Commutativity. Consider any 2 arbitrary Loubéré Magic Squares a, b, ; such that $a+b=b+a$; then from the example above, the corresponding eigen values of a; $\lambda_{a_{1}}, \lambda_{a_{2}} \lambda_{a_{3}}$; the corresponding eigen values of $\mathrm{b} ; \lambda_{b_{1}}, \lambda_{b_{2}}, \lambda_{b_{3}}$; are such that $\lambda_{a_{1}}+\lambda_{b_{1}}=$ $\lambda_{b_{1}}+\lambda_{a_{1}}, \lambda_{a_{2}}+\lambda_{b_{2}}=\lambda_{b_{2}}+\lambda_{a_{2}}$, and $\lambda_{a_{3}}+\lambda_{b_{3}}=\lambda_{b_{3}}+\lambda_{a_{3}}$.

The idea of eigen values computation of a magic square is conceived from the work of [1].

## VI. The Subelement Magic Souares Semigroup and Group

The set of least subelement of Loubéré Magic Squares is a subset of pancolumn $3 \times 3$ Magic Squares. By convention, the $3 \times 3$ Loubéré Magic Square(since not pancolumn) is not a self subelement. The sum of two arbitrary subelements of $m \times m$ Loubéré Magic Squares is a subelement of $m \times m$ Loubéré Magic Square, hence closure property is exhibited. Associativity, Identity, Inverse and Commutativity Properties are inherited from the super elements, the Loubéré Magic Squares. Both the binary and the unary operations of the super elements and of the subelements are equal.

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## Note :

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## Author Guidelines:

1. General,
2. Ethical Guidelines,
3. Submission of Manuscripts,
4. Manuscript's Category,
5. Structure and Format of Manuscript,
6. After Acceptance.

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- As always, give awareness to spelling, simplicity and correctness of sentences and phrases.


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This part is supposed to be the easiest to carve if you have good skills. A sound written Procedures segment allows a capable scientist to replacement your results. Present precise information about your supplies. The suppliers and clarity of reagents can be helpful bits of information. Present methods in sequential order but linked methodologies can be grouped as a segment. Be concise when relating the protocols. Attempt for the least amount of information that would permit another capable scientist to spare your outcome but be cautious that vital information is integrated. The use of subheadings is suggested and ought to be synchronized with the results section. When a technique is used that has been well described in another object, mention the specific item describing a way but draw the basic principle while stating the situation. The purpose is to text all particular resources and broad procedures, so that another person may use some or all of the methods in one more study or referee the scientific value of your work. It is not to be a step by step report of the whole thing you did, nor is a methods section a set of orders.

## Materials:

- Explain materials individually only if the study is so complex that it saves liberty this way.
- Embrace particular materials, and any tools or provisions that are not frequently found in laboratories.
- Do not take in frequently found.
- If use of a definite type of tools.
- Materials may be reported in a part section or else they may be recognized along with your measures.


## Methods:

- Report the method (not particulars of each process that engaged the same methodology)
- Describe the method entirely
- To be succinct, present methods under headings dedicated to specific dealings or groups of measures
- Simplify - details how procedures were completed not how they were exclusively performed on a particular day.
- If well known procedures were used, account the procedure by name, possibly with reference, and that's all.

Approach:

- It is embarrassed or not possible to use vigorous voice when documenting methods with no using first person, which would focus the reviewer's interest on the researcher rather than the job. As a result when script up the methods most authors use third person passive voice.
- Use standard style in this and in every other part of the paper - avoid familiar lists, and use full sentences.

What to keep away from

- Resources and methods are not a set of information.
- Skip all descriptive information and surroundings - save it for the argument.
- Leave out information that is immaterial to a third party.


## Results:

The principle of a results segment is to present and demonstrate your conclusion. Create this part a entirely objective details of the outcome, and save all understanding for the discussion.

The page length of this segment is set by the sum and types of data to be reported. Carry on to be to the point, by means of statistics and tables, if suitable, to present consequences most efficiently.You must obviously differentiate material that would usually be incorporated in a study editorial from any unprocessed data or additional appendix matter that would not be available. In fact, such matter should not be submitted at all except requested by the instructor.

- Sum up your conclusion in text and demonstrate them, if suitable, with figures and tables.
- In manuscript, explain each of your consequences, point the reader to remarks that are most appropriate.
- Present a background, such as by describing the question that was addressed by creation an exacting study.
- Explain results of control experiments and comprise remarks that are not accessible in a prescribed figure or table, if appropriate.
- Examine your data, then prepare the analyzed (transformed) data in the form of a figure (graph), table, or in manuscript form. What to stay away from
- Do not discuss or infer your outcome, report surroundings information, or try to explain anything.
- Not at all, take in raw data or intermediate calculations in a research manuscript.
- Do not present the similar data more than once.
- Manuscript should complement any figures or tables, not duplicate the identical information.
- Never confuse figures with tables - there is a difference.

Approach

- As forever, use past tense when you submit to your results, and put the whole thing in a reasonable order.
- Put figures and tables, appropriately numbered, in order at the end of the report
- If you desire, you may place your figures and tables properly within the text of your results part.

Figures and tables

- If you put figures and tables at the end of the details, make certain that they are visibly distinguished from any attach appendix materials, such as raw facts
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- In spite of position, each table must be titled, numbered one after the other and complete with heading
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- Recommendations for detailed papers will offer supplementary suggestions.

Approach:

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